

Hierarchical Adaptive Cross Approximation GMRES Technique for Solution of Acoustic Problems Using the Boundary Element Method

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Abstract: In this paper a new Rapid Acoustic Boundary Element Method (RABEM) is presented using a Hierarchical GMRES solver for 3D acoustic problems. The Adaptive Cross Approximation is used to generate both the system matrix and the right hand side vector. The ACA is also used to evaluate the potential and the particle velocity values at selected internal points. Two different GMRES solution strategies (without preconditioner and with a block diagonal preconditioner) are developed and tested for low and high frequency problems. Implementation of different boundary conditions (i.e. Dirichlet, Neumann and mixed Robin) is also described. The applications presented include the problem of noise acting on a row of aircraft seats and problem of engine noise emanating from the Falcon aircraft. For the first example, the new RABEM is shown to be faster than the Fast Multipole Methods. The tests demonstrated that the new solver can achieve CPU times of almost $O(N)$ for low frequency and $O(N \log N)$ for high frequency problems.

1 Introduction

The Boundary Element Method (BEM) is one of the most general and efficient numerical technique for solving acoustic problems [Wrobel and Aliabadi (2002)]. The boundary element discretisation of the surface of the problem leads to a non-symmetric and fully populated system matrix. For a standard BEM formulation both the memory storage and the setting up of the system matrices are of $O(N^2)$, where N denotes the degree of freedom. Moreover, direct solvers require $O(N^3)$ operations while iterative solvers $O(kN^2)$, where k is the number of iterations.

To overcome the difficulties related to storage and solution time a number of techniques have been proposed which include using block-based solvers [Crotty (1982); Rigby and Aliabadi (1995)], lumping techniques [Kane and Kumar (1990)] and

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iterative solvers [Mansur, Araújo and Malaghini (1992)]. The Adaptive Cross Approximation is an effective technique for solving non-symmetric and fully populated matrices and decreases the CPU time significantly (see Bebendorf and Rjasanow (2003)) and has been applied to the Helmholtz equation by, for example, Von Estorff et al. (2005). The solution of linear system of equations is accelerated by calculating only few entries of the original matrix. The basic idea is to divide the whole matrix into two rank (*low* and *full* rank) blocks based on (ACA) size and distance between a group of collocation points and a group of boundary elements. The ACA algorithm has been applied to the low rank blocks achieving approximately $O(N)$ for both storage and matrix-vector multiplication [Benedetti, Aliabadi and Daví (2007)]. An alternative approach to the fast BEM solver is the Fast Multipole Method (FMM) (see for example Aoki, Amaya, Urago and Nakayama (2004); Chew, Song, Cui, Velamparambil, Hastriter and Hu (2004); He, Lim and Lim (2008); Wang and Yao (2005, 2008)). Although FMM techniques are efficient for the fast solution of boundary element problems, their main disadvantage is that the knowledge of the kernel expansion is required in order to carry out the integration; all the terms of the series needed to reach a given accuracy must be computed in advance and then integrated, which can lead to a significant modification of the integration procedures in standard BEM codes. From an algebraic point of view however, the integration of a degenerate kernel, i.e. of a kernel expanded in series, over a cluster of elements corresponds to the approximation of the corresponding matrix block by a low rank block. Comparison between the FMM and ACA method were made by Wang, Hall, Yu and Yao (2008) for three-dimensional Laplace problems. In their study a fully pivoted ACA was shown to require considerably more computational effort than FMM. However, it must be noted that the full pivoted approach is well known to be much slower than partially pivoted approach developed in our paper. The main reason behind this is that fully pivoted approach requires the knowledge of the full matrix whereas the partially pivoted approach would only require generation of individual matrix entries. Indeed, in our paper the partially pivoted ACA approach for acoustic problems is shown to be superior to the FMM for an example considered.

Developing iterative solvers for non symmetric linear systems has been widely investigated. One of the most popular is the generalized minimal residual method (GMRES) proposed by Saad and Schultz (1986) and further developed by other authors [Leung and Walker (1997); Merkel, Bulgakov, Bialecki and Kuhn (1998); Amini and Maines (1998)].

In this paper a new hierarchical adaptive cross approximation technique coupled with GMRES is presented for 3D boundary element solution of Helmholtz problems. Particular attention is paid to implementation of different types of boundary

conditions (i.e. Dirichlet, Neumann and mixed Robin) into the ACA solution algorithm. The ACA and the H -Matrix format are also employed for the post processing steps required to evaluate the values of the potential and particle velocity at the internal points. The constant elements are utilized to discretize the problem. Initially two simple benchmark problems of the pulsating sphere and of the scattering of a plane wave from a rigid sphere are investigated. Next, large scale applications of involving a row of three seats representing the seats in an aircraft cabin and noise emanating from engines of the Dassault Falcon aircraft are presented. The tests carried out show that the new assembly and solution technique can achieve CPU times of almost $O(N)$ for low frequency and $O(N \log N)$ for high frequency problems.

2 The Boundary Element Method

In this section the three dimensional Boundary Element Method (BEM) for sound propagation problems is briefly reviewed. The boundary integral formulation for sound propagation can be written as

$$\begin{aligned} C(\mathbf{x}')u^j(\mathbf{x}') + \int_{\Gamma} q^*(\mathbf{x}', \mathbf{x})u(\mathbf{x})d\Gamma(\mathbf{x}) \\ = \int_{\Gamma} u^*(\mathbf{x}', \mathbf{x})q(\mathbf{x})d\Gamma(\mathbf{x}) + \int_{\Omega} u^*(\mathbf{x}', \mathbf{X}^s)\frac{1}{c^2}b(\mathbf{X}^s)d\Omega(\mathbf{X}^s) \end{aligned} \quad (1)$$

where $u(\mathbf{x})$ and $q(\mathbf{x})$ denote the potential and flux, respectively on the boundary Γ . The terms $u^*(\mathbf{x}', \mathbf{x})$ and $q^*(\mathbf{x}', \mathbf{x})$ are the potential and the flux fundamental solutions, respectively, and $C(\mathbf{x}')$ is the free term which its value depends on the position of \mathbf{x}' on the boundary (see Wrobel and Aliabadi (2002)). The last term refers to the presence of sources within the domain Ω with strength $b(\mathbf{X}^s)/c^2$ where c is the sound velocity. The BE formulation automatically satisfies the Sommerfeld radiation condition [Wrobel and Aliabadi (2002)].

The boundary integral equation (1) is discretised with constant elements. The discretised form of (1) can be written in a matrix form as

$$\mathbf{H}\mathbf{U} = \mathbf{G}\mathbf{Q} + \mathbf{P}(\mathbf{X}^s) \quad (2)$$

where \mathbf{H} and \mathbf{G} are $N \times N$ coefficient matrices corresponding to integrals of the product of the Jacobian of transformation with flux and potential fundamental solutions, respectively, and \mathbf{U} and \mathbf{Q} are the $N \times 1$ potential and flux vectors, respectively. Finally, the last integral in equation (1) produces the $N \times 1$ vector \mathbf{P} , created by NP sources within the domain Ω .

The implementation of the boundary conditions for acoustics wave equations has received much attention in the last decades. In particular Clayton and Engquist

(1977) and more recently Higdon (1992) included the absorbing boundary conditions for acoustic simulations in their computing numerical models. The possible presence of absorbent material with a certain value of impedance has also been introduced by a quantity that relates directly the potential and the flux, here called *artificial admittance* (ζ) which is defined as follows

$$\zeta(\omega) = \frac{q}{u} \quad (3)$$

Equation (3) creates the impedance boundary conditions and can be easily accounted for in (2). After substitution of the prescribed boundary conditions, the resulting system of algebraic equations may be written as

$$\mathbf{A}\mathbf{Y} = \mathbf{F} \quad (4)$$

where \mathbf{Y} is the vector containing the unknown boundary potentials and fluxes, \mathbf{A} is a coefficient matrix which is non-symmetric and densely populated, and \mathbf{F} is obtained by multiplying the prescribed boundary conditions by the corresponding columns of the \mathbf{G} and \mathbf{H} matrices.

Similarly, for internal potential values the resulting system of equations can be written as

$$\mathbf{U}(\mathbf{X}) = -\bar{\mathbf{H}}\mathbf{U} + \bar{\mathbf{G}}\mathbf{Q} + \mathbf{P}(\mathbf{X}^s) \quad (5)$$

where $\bar{\mathbf{H}}$ and $\bar{\mathbf{G}}$ are similar to \mathbf{H} and \mathbf{G} , but evaluated at internal points (NI), and for the internal particle velocity

$$\mathbf{U}'(\mathbf{X}) = -\bar{\mathbf{H}}'\mathbf{U} + \bar{\mathbf{G}}'\mathbf{Q} + \mathbf{P}'(\mathbf{X}^s) \quad (6)$$

where superscript “ $'$ ” denotes derivative with respect to Cartesian coordinates (x , y , z); $\bar{\mathbf{H}}'$ and $\bar{\mathbf{G}}'$ contain the integrals of the derivatives of the fundamental solutions and are $NI \times N$ matrices calculated by substituting the fundamental solutions $u^*(\mathbf{x}', \mathbf{x})$ and $q^*(\mathbf{x}', \mathbf{x})$ and their derivatives (see Dominguez (1993)) and $\mathbf{P}'(\mathbf{X}^s)$ is a $NI \times 1$ vector containing the contribution to the particle velocity generated by the sources within the domain Ω .

3 Hierarchical BEM for Acoustics

The use of hierarchical matrices [Hackbush (1999); Hackbush and Khoromskij (2000); Borm, Grasedyck and Hackbush (2003)] for the representation of BEM system of equations, in conjunction with Krylov subspace methods [Leung and Walker (1997); Merkel, Bulgakov, Bialecki and Kuhn (1998)], is extended to BEM acoustic problems with different boundary conditions. The technique speeds up the computation, whilst maintaining the required accuracy and saving on the memory storage.

3.1 Matrix Assembly using the Collocation Method

The hierarchical representation of a boundary element matrix can be regarded as a subdivision of the matrix itself into a collection of blocks, some of which, called *low rank blocks*, allow a special compressed representation, while others, called *full rank blocks* are represented in their entirety. The subdivision and the subsequent classification is achieved starting from the boundary element mesh and is based on the grouping of the nodes and elements into clusters of *close* nodes and elements. A block populated by integrating over a cluster of elements whose distance, suitably defined, from the cluster of collocation nodes is above a certain threshold is called *admissible* and it can be represented in the low rank format. The remaining blocks are generated and stored in their entirety.

The process leading to the subdivision in sub-blocks and to their further classification is based on a preliminary *hierarchical partition* of the matrix index set aimed at grouping subsets of indices corresponding to contiguous nodes and elements, on the basis of some computationally efficient geometrical criterion. The process starts from the complete set of indices $I = \{1, 2, \dots, n\}$ where n denotes the number of collocation points. This initial set constitutes the *root* of the tree. Each cluster in the tree, called *tree node* (not to be confused with geometrical discretisation nodes) is split into two subsets, called *sons*, on the basis of some selected criterion. For this study the longest extended dimension of the whole geometry (step 1 of the figure 1) is detected first and its central point calculated. This point divides the mesh into two tree nodes (step 2). The geometries of the mesh of each of these blocks is then divided into two other tree nodes by following the same procedure (step 3). Such an iterative procedure is repeated until each tree node is constituted by a smaller number of elements than the cardinality.

The common tree node from which two sets originate is called the *parent*. The tree nodes that cannot be further split are the *leaves* of the tree. Usually a tree node cannot be further split when it contains a number of indices equal to or less than a minimum number n_{min} , called *cardinality* of the tree, previously fixed value. This partition is stored in a binary tree of index subsets, or *cluster tree*, that constitutes the basis for the subsequent construction of the hierarchical *block* subdivision, that will be stored in a quaternary *block tree*.

The block tree is built recursively starting from the complete index $I \times I$ (both rows and columns) of the collocation matrix and the previously found cluster tree. The objective of this process is to split hierarchically the matrix into sub-blocks and to classify the leaves of the tree into admissible (low-rank) or non admissible (full-rank) blocks. The classification is based on a geometrical criterion that assesses the separation of the clusters of boundary elements associated to the considered block.

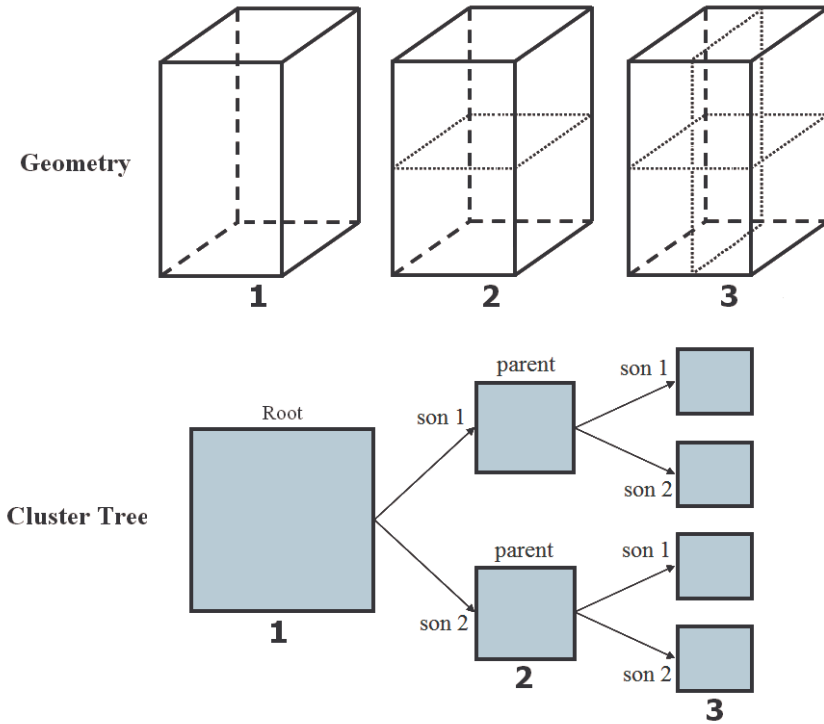


Figure 1: A schematic of the first two iterations of the cluster tree creation: 1) the whole geometry is divided 2) into two parts, each of which is then subdivided again 3) into two parts.

Such a criterion relates directly to the boundary mesh.

Let Ω_{x_0} denote the cluster of elements containing the discretisation nodes corresponding to the row indices of the considered block and Ω_x the set of elements over which the integration is carried out to compute the coefficient corresponding to the column indices. The admissibility condition can be written as

$$\min(\text{diam } \Omega_{x_0}, \text{diam } \Omega_x) \leq \eta \cdot \text{dist}(\Omega_{x_0}, \Omega_x) \tag{7}$$

where $\eta > 0$ is a parameter influencing the number of admissible blocks on one hand and the convergence speed of the adaptive approximation of low rank blocks on the other hand [Borm, Grasedyck and Hackbush (2003)].

Since the actual diameters and the distance between two clusters are generally time consuming to be exactly computed, the condition is usually assessed with respect to bounding boxes parallel to the reference axes [Giebermann (2001); Grasedyck

(2005)]. In this case Ω_{x_0} and Ω_x in the equation (7) are replaced by the boxes Q_{x_0} and Q_x . The bounding box clustering technique adopted in the present work is generally used for its simplicity, although it produces non-optimal partitions that can be improved by suitable procedures, as will be illustrated next. Other clustering techniques able to produce better initial partitions have been proposed in the literature. The construction using the *principal component analysis* [Bebendorf (2006)] significantly improves the quality of the initial partition.

The algorithm is graphically illustrated in the figure 2. Starting from the root (the entire matrix), each block is subdivided into four sub-blocks until either the admissibility condition is satisfied or the block is sufficiently small that it cannot be further subdivided. The clear grey boxes represent low rank blocks while the dark grey boxes are the full rank ones.

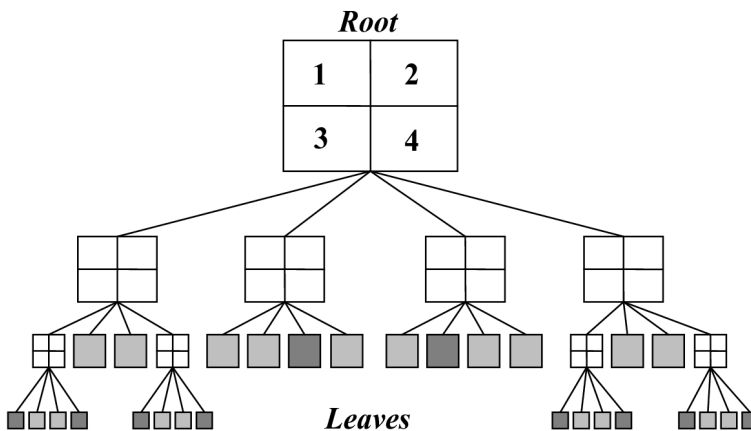


Figure 2: A schematic presentation of the first three iterations that form the block tree.

An admissible block can be represented in low rank format. Such representation constitutes an approximation of the discrete integral operator based, from the analytical point of view, on a suitable expansion of the kernel of the continuous integral operator [Bebendorf and Rjasanow (2003); Tyrtshnikov (1996); Goreinov, Tyrtshnikov and Zamarashkin (1997); Bebendorf (2000)]. This expansion, and consequently the existence of low rank *approximants*, is based on the *asymptotic smoothness* of the kernel functions [Bebendorf and Grzhibovskis (2006)], i.e. on the fact that the kernels $u^*(\mathbf{x}', \mathbf{x})$ and $q^*(\mathbf{x}', \mathbf{x})$ are singular only when $\mathbf{x}' = \mathbf{x}$. Asymptotic smoothness represents a sufficient condition for the existence of low rank approximates. For more detailed information about asymptotic

smoothness the readers are referred to the works [Bebendorf and Rjasanow (2003); Bebendorf (2000)].

Let C be an $m \times n$ admissible block. It admits the low rank representation

$$C \simeq C_k = A \cdot B^T = \sum_{i=1}^k a_i \cdot b_i^T \tag{8}$$

where A is of order $m \times k$ and B is of order $n \times k$, with k being the *rank* of the new representation. The approximating block C_k satisfies the relation $\|C - C_k\|_F \leq \varepsilon \|C\|_F$, where $\|\cdot\|_F$ represents the *Frobenius norm* and ε is the set accuracy. Sometimes it is useful to represent the matrix using the alternative sum representation, where a_i and b_i are the i -th columns of A and B , respectively. The approximate representation allows storage savings with respect to the full rank representation and speeds up the matrix-vector product [Grasedyck and Hackbusch (2003)].

The low rank blocks are built by computing and storing only *some* of the entries of the original blocks. Such entries allow for computation of the columns a_i and b_i of the representation [8] through suitable algorithms, known as adaptive cross approximation (ACA) [Bebendorf and Rjasanow (2003); Bebendorf (2000)]. The ACA algorithms allow to reach the selected collocation matrix accuracy ε_c adaptively. The stopping criterion is based on the assessment of the convergence of the approximating block in terms of the Frobenius norm [Bebendorf and Rjasanow (2003); Bebendorf (2000)]. Once the k -th couple (a_k, b_k) has been computed the Frobenius norm of the approximation can be computed by the following recursive formula

$$\|C_k\|_F^2 = \|C_{k-1}\|_F^2 + 2 \sum_{i=1}^{k-1} (a_k^T a_i)(b_i^T b_k) + \|a_k\|_F^2 \|b_k\|_F^2 \tag{9}$$

where a_k and b_k represent the column and row computed at the k -th iteration. A suitable stopping criterion can be expressed as

$$\|a_k\|_F \|b_k\|_F \leq \varepsilon \|C_k\|_F \tag{10}$$

that prescribes to stop the iteration when the inequality is satisfied for a required preset accuracy ε .

3.2 Boundary Conditions and Right Hand Side Setting

By applying the ACA, the actual setting of the final system for mixed boundary condition problems requires some additional considerations. In this work different types of boundary conditions are considered in such a way that rigid, soft and absorbing surfaces can all be studied in order to simulate a real situation and to

perform an eventual parametric analysis. An initial routine selects the matrix that is calculated first based on the boundary conditions that are predominant for that block. If the flux (or potential) is predominant then the matrix \mathbf{H} (or \mathbf{G}) is calculated first, otherwise, in the case the boundary condition are in terms of impedance, the ACA algorithm is applied to both the matrices.

In order to speed up significantly the system solution of equation (4) the “*BlockNode*” pointer suggested by Benedetti, Aliabadi and Daví (2007) is used here. The contributions of each block to the solving matrix \mathbf{A} is stored in the *BlockNode* pointer and calculated at each step as well as the contribution of the right hand side vector that is, instead, stored in the global vector \mathbf{F} to save memory.

The ACA algorithm is applied to one or both of the matrices \mathbf{G} and \mathbf{H} depending upon the boundary conditions that are predominant for each block matrix. Particular care must be taken when not pure Dirichlet, Neumann or mixed Robin conditions are applied. There may be four different cases:

- a) BCs mainly in terms of flux (or potential) except for the k^{th} value expressed in terms of potential (or flux).

As an example, one can consider that the boundary conditions are expressed in terms of flux (the first block matrix calculated with ACA is \mathbf{H}) and at the node k is expressed in terms of potential. A routine replaces each k^{th} value of each row previously calculated with zero, adds the opposite in sign of the k^{th} column of the \mathbf{G} block matrix and also adds a zero row with 1 at the k^{th} element (see figure 3).

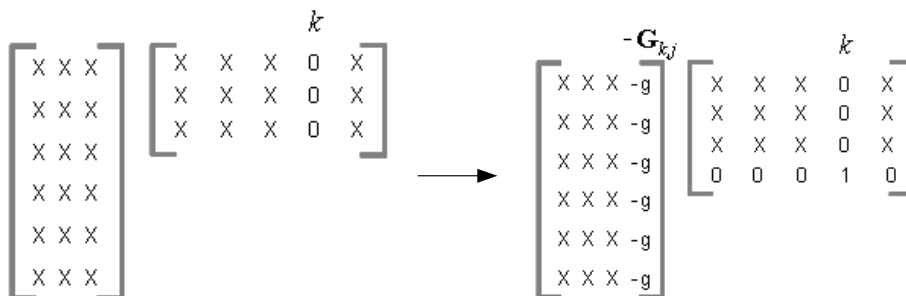


Figure 3: Substitution of the k^{th} values of each row with zero, addition of the k^{th} column of the \mathbf{G} block matrix and of a zero row with 1 at the k^{th} element.

- b) BCs mainly in terms of flux except for the k^{th} value expressed in terms of impedance.

Here simply add the k^{th} column of the \mathbf{G} block matrix not yet calculated multiplied by the artificial admittance ζ opposite in sign, and adds a zero row with 1 at the k^{th} element (see figure 4).

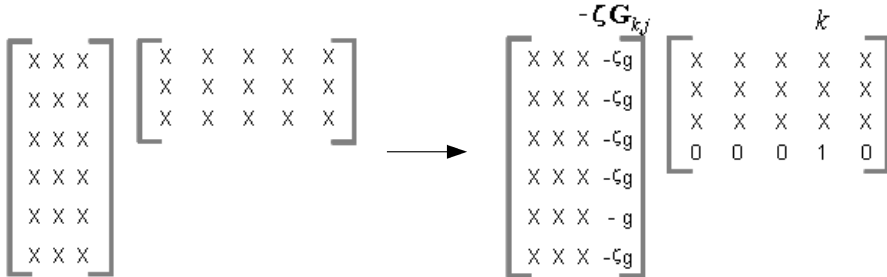


Figure 4: Addition of the k^{th} column of the \mathbf{G} block matrix multiplied by the artificial admittance ζ opposite in sign and of a zero row with 1 at the k^{th} element.

- c) BCs mainly in terms of potential except for the k^{th} value expressed in terms of impedance.

In this case multiply the k^{th} row of the \mathbf{G} block matrix (previously calculated) by the artificial admittance ζ opposite in sign. Then another routine adds the k^{th} column of the \mathbf{H} block matrix not yet calculated and a zero row with 1 at the k^{th} element (see figure 5).

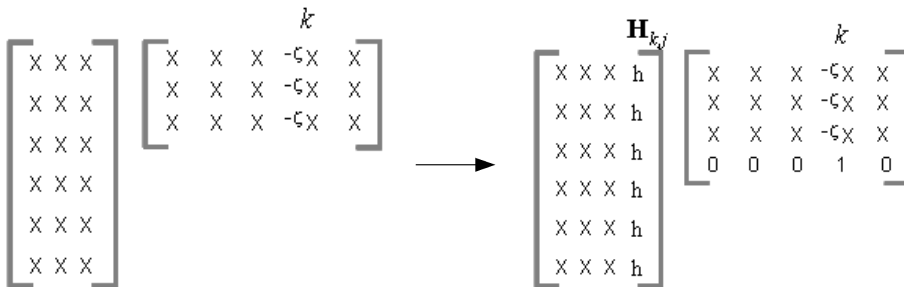


Figure 5: Multiplication of the k^{th} row of the \mathbf{G} block matrix by the artificial admittance ζ opposite in sign, addition of the k^{th} column of the \mathbf{H} block matrix and of a zero row with 1 at the k^{th} element.

- d) BCs mainly in terms of impedance except for the k^{th} value expressed in terms of potential (or flux).

In this case the ACA algorithm acts earlier on the \mathbf{H} block matrix. If the ACA stop criterion is reached and the k^{th} value is in terms of potential, the k^{th} value of each row calculated is set to zero. The ACA algorithm is finally applied to the \mathbf{G} block matrix whose columns are multiplied $-\zeta$, and the k^{th} elements of its rows are divided by ζ (see figure 6).

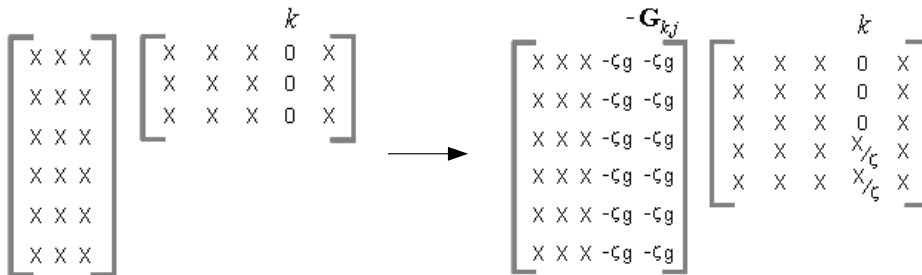


Figure 6: Substitution of the k^{th} values of each row with zero, addition of the columns of the \mathbf{G} block matrix multiplied by the artificial admittance opposite in sign and of its rows where the k^{th} elements are divided by ζ .

Regarding the setting up of the right hand side vector some additional considerations are required as explained next. Recalling the ACA algorithm, the routine that calculates the i^{th} row of one of the block matrices \mathbf{G} or \mathbf{H} also calculates the i^{th} row of the other block matrix. Thus, the right hand side contribution of that row for the block matrix analyzed is directly calculated.

Now, there are two main cases to analyze.

- BCs mainly in terms of flux (or potential).

Here, one may need to apply the ACA on the matrix \mathbf{G} (or \mathbf{H}). In fact, if all the boundary conditions are zero, there is no need to calculate the contribution of the block matrix to the final right hand side and another block matrix can be analyzed. Moreover, owing to the fact that the number of entries needed for the ACA is, in most the cases, equal for both the block matrix \mathbf{G} and \mathbf{H} , may be convenient to calculate the contribution to the right hand side with a standard procedure (see appendix).

- BCs mainly in terms of admittance.

Once the ACA as been applied to the block matrix \mathbf{H} , the possible contribution to the right hand side vector \mathbf{F} , due to the presence of the k^{th} value of the boundary conditions expressed in terms of potential, is easily calculated by multiplying each column of the ACA to the k^{th} value of each row with opposite in sign. Finally, if the ACA algorithm applied to the block matrix \mathbf{G} is also successful reached, the contribution to the right hand side of the eventual presence of the boundary condition expressed in terms of flux is clearly calculated.

3.3 System Solution

Once the hierarchical representation of the collocation matrix has been established, the solution of the system can be computed through iterative solvers with or without preconditioners. When the condition number is high it slows down the convergence rate, as is often the case when dealing with BEM systems, a preconditioner can be computed taking full advantage of the representation in hierarchical format. For $Ay = F$ system a left preconditioner is an easily invertible matrix M such that the condition number of the system $M^{-1}Ay = M^{-1}F$ results lower than the original one, improving thus the convergence rate of the iterative solver.

In the present work a GMRES iterative solver [Saad and Schultz (1986)] without and with preconditioner are used for the solution of the system of equations. The method, proposed by Saad and Schultz (1986) and further developed by many other author, is a Krylov based iterative method for the numerical solution of a system of linear equations, as in equation (4). The Arnoldi iteration is used to form a basis for the Krylov subspace.

In all the on-line complex GMRES routines available for FORTRAN, the generation of the whole matrix \mathbf{A} is needed. In the present work the whole matrix is never calculated, resulting in a reduction in memory requirements. The routine adopted is an $O(N)$ procedure.

Internal Points

The ACA algorithm can also be used to speed up the calculation of potential and particle velocity at internal points from the boundary values of potential and flux. To achieve this, a cluster tree for the internal points is generated following the procedure seen at the beginning of this section. Thus the hierarchical tree is generated by considering the columns and the rows of the block matrices $\overline{\mathbf{H}}$, $\overline{\mathbf{G}}$, $\overline{\mathbf{H}}'$ and $\overline{\mathbf{G}}'$ (see equations (5) and (6)). The ACA procedure is applied four times to all the matrices. The name of such a pointer, the values of the parameters η and *cardinality* remain unchanged and hence the same routines as for the boundary points can be applied.

4 Numerical Results

In order to demonstrate the efficiency of the proposed method a series of numerical tests have been performed. Comparisons with the standard BEM with a direct solver [Dominguez (1993)], with LMS Virtual.Lab [LMS International (2008)] and with Fast Multipole Method (FMM) for acoustics [FastBEM Acoustics (R)] shows that the formulation presented here can significantly reduce the system solution time. All of the numerical simulations were executed on a Xeon(TM) CPU 3.00GHz processor with 2GB of RAM.

In all the simulations performed the value of the parameter $\eta=10$ and the cardinality is set to 22. These values are found to give the best performances. Moreover, the value of the cardinality is quite restricted. In fact, as this value decreases, the stopping criterion of equation (10) is not successful for many low rank blocks, rather, it increases the CPU time due to bigger full block dimensions. The value of the parameter η can be modified between 1 and 1000 without loss of accuracy and resulting in a 5% CPU time acceleration. Furthermore, the optimum value of this parameter depends upon the geometry and the elements of the mesh.

Concerning the storage memory requirements of the matrix \mathbf{A} , figure 7 shows the block-wise structure of the collocation matrix as generated by the ACA algorithm for a pulsating sphere constituted with 6100 elements for four different values of η (1, 2, 4 and 10). The tone of grey is proportional to the ratio between the memory required for the low rank representation and the memory required for a standard format. Hence black blocks stand for the full rank block matrix, while almost white blocks are those for which the ACA compression works better.

RABEM vs. SBEM vs. LMS

The problem of the sound radiated by a pulsating sphere with radius, wave number k and uniform radial velocity all equal to unity is investigated. The acoustic wave velocity and the medium density are set equal to unity. The three-dimensional scalar wave propagation problems using constant boundary elements [Dominguez (1993)], will be referred to as “SBEM”, and is used for comparison purposes. Table 1 compares the two codes speed up ratio (defined as CPU time of SBEM divided by CPU time of RABEM) for setting up the system solution, finding the solution and for the total time for eight different mesh accuracies. The percentage error in all the simulations is less than 0.02 %. Comparison are also made with the commercial code LMS Virtual.Lab as shown in table 1.

The comparison with SBEM demonstrates that the RABEM formulation is effective in the population of the linear system when more than 2000 degrees of freedom mesh is studied. However, the GMRES routine is effective even considering a

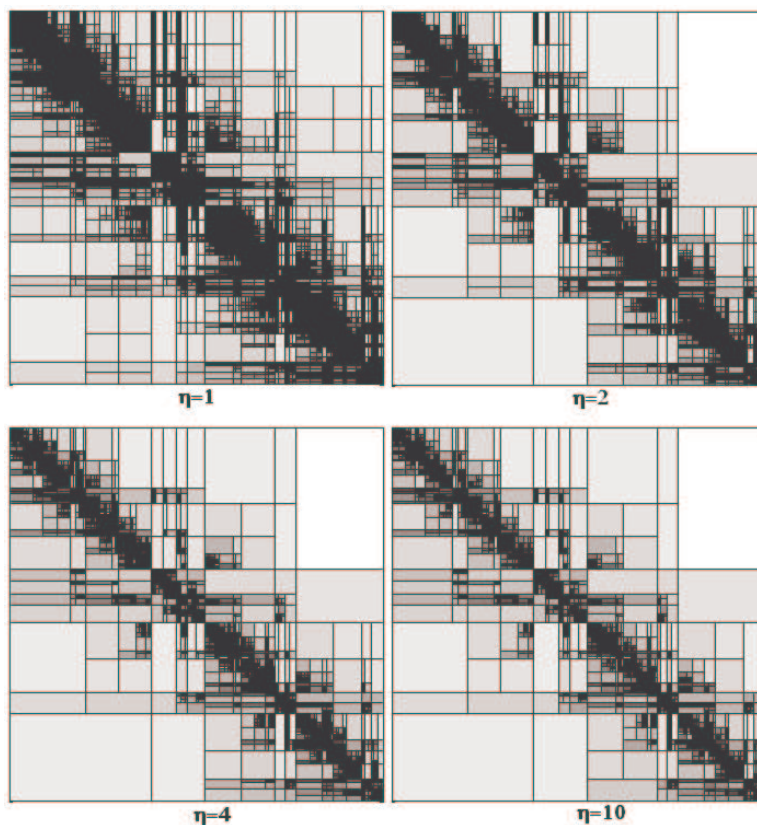


Figure 7: Block-wise representation of the ACA generated matrix for $\eta = 1$, $\eta = 2$, $\eta = 4$ and $\eta = 10$.

limited number of elements.

Comparison of RABEM and LMS for the rigid sphere (same medium characteristic as seen in the first example scattered by a plane waves with unit amplitude and zero phase at the origin) is shown in figure 8.

It must be pointed out that no attempt has been made to optimize the general BEM routines of integration etc. in the RABEM and they remain the same as in the standard BEM (SBEM).

Internal Points To demonstrate the speed up that can be achieve when the ACA is applied at selected internal points, the CPU time of the RABEM and standard code is compared for the 3240 dof pulsating sphere and for four sets of internal

Table 1: Speed up ratio between the RABEM, SBEM and LMS.

dof	SBEM/RABEM			LMS/RABEM
	Setting up	System Solution	Solver	Total
646		0.8	107.9	2.2
1126		1.0	180.6	3.6
2112		1.5	611.6	7.8
3240		2.0	1669.9	18.3
4764		2.6	44098.7	44.6
6100		3.2	5562.1	55.4
8488		4.1	13165.7	132
10340		5.0	18513.2	175.7
20232		-	-	-
28194		-	-	-

Scattering of a Plane wave from a rigid Sphere
LMS/RABEM

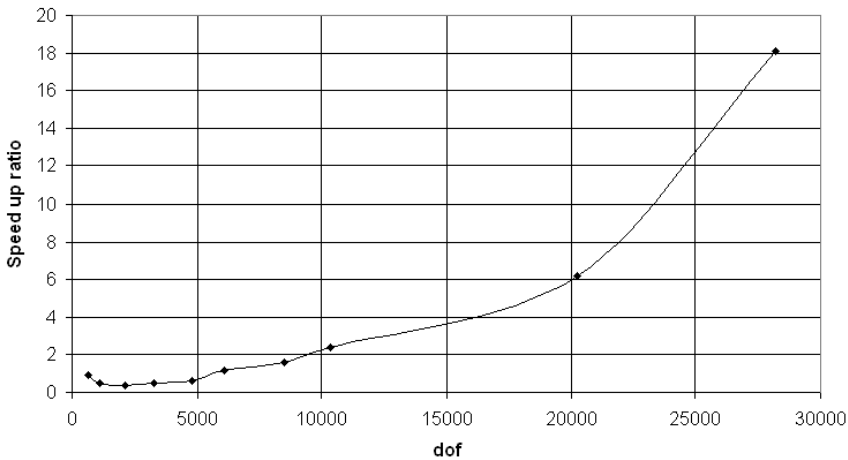


Figure 8: Speed up ratio for a scattered rigid sphere by a plane wave - LMS/RABEM.

points: 546, 2184, 4368 and 10374. Once again the speed up ratio is defined as the ratio of CPU times obtained by SBEM over RABEM. Figure 9 shows the resulting acceleration during the procedure for calculating the internal point potential. The RABEM CPU time required depends upon the location of such internal points as

well as the dof of the boundary mesh, where as the standard procedure is invariant to the location, while the rank of the ACA decreases as the points are further away.

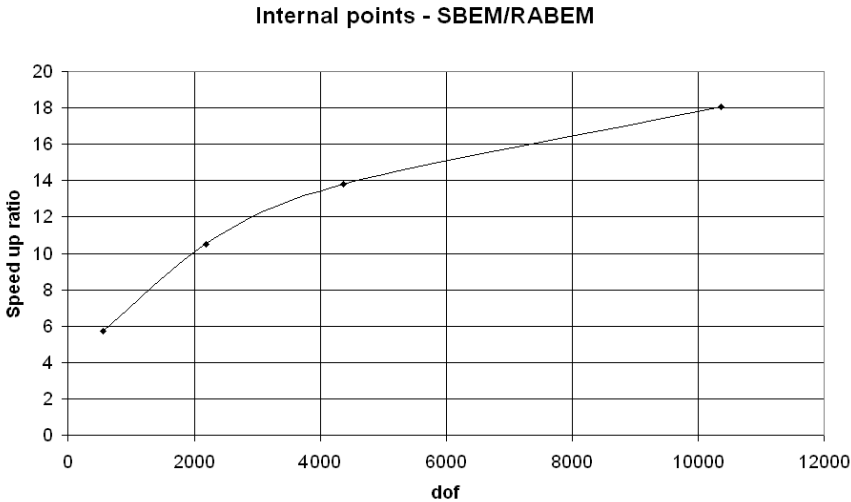


Figure 9: Speed up ratios relates to four different groups of internal points for a pulsating sphere discretised with 3240 constant triangular elements.

Frequency Variation

The variation of the computational time with respect to varying frequency is studied in the problem of pulsating sphere. The example tested in the first above (with the same boundary conditions) is studies for 8 angular frequencies (1, 2, 4, 8, 12, 16, 20, 30 Hz/rad). The mesh has 10340 constant triangular elements. Figure 10 shows the ratios obtained by dividing the CPU times for different angular frequencies with the CPU time for the angular frequency of 1 Hz/rad. As evident the CPU time increases almost linearly.

A Row of Seat in an Aircraft Cabin

Another example to which the RABEM has been applied is a row of three seats of an aircraft cabin as shown in figure 11. The surface of the seats have been divided into 27,284 constant triangular elements. The maximum extensions along the *x*, *y* and *z* axis are 694, 1563 and 960 millimeters, respectively. The sizes of the smallest and the largest triangles are 3.58E-03 and 2.38E-02 millimeter, respectively, so the maximum frequency that can be applied is around 1.3-1.5 kHz. In the simulation performed all the surfaces has been set as hard and a monopole

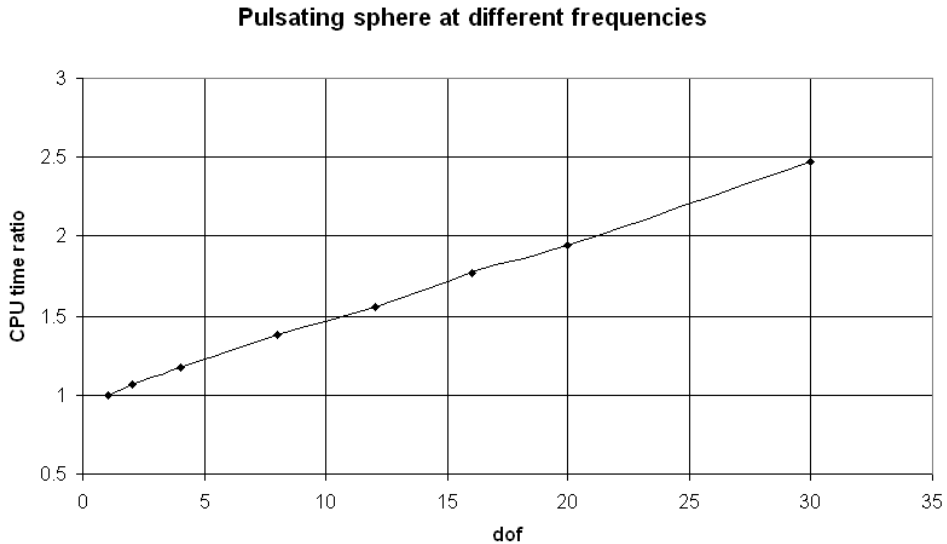


Figure 10: CPU time ratios for different angular frequencies for a pulsating sphere meshed with 10340 triangular elements.

source with unit complex potential amplitude has been inserted at the center of the left seat with a distance of 7.4 and 21.8 cm from the cushion and the back of the seat, respectively. The speed up ratio related to the Fast Multipole Method FMM and for the RABEM without preconditioner (RABEM) and with a block diagonal preconditioner (RABEMpdb) are displayed in the figure 12 for nine different frequencies (100, 200, 400, 600, 800, 900, 1000, 1100 and 1200 Hz). As evident RABEM is faster than the FMM code [FastBEM Acoustics (R)] for higher frequencies. In particular for frequencies lower than 400Hz FastBEM is slightly faster, but as the frequency increases the speed up ratio grows up almost linearly up to 2.6 for 1200Hz. In this example the RABEMpdb is never faster than RABEM.

The same geometry has been simulated by Virtual.Lab LMS for a frequency of 100Hz. The RABEM code solved it 4.2 times faster than LMS.

Finally, in figure 13 the solution in terms of real part of the potential for the 100Hz frequency is displayed.

Dassault Falcon Airplane

The simulations of a 47998 elements mesh (see figure 14) of a model representing the Dassault Falcon airplane is presented. The total length of the aircraft is 18.5m

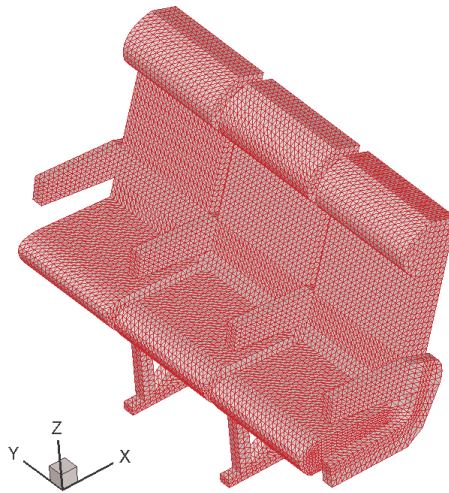


Figure 11: Surface of an airplane seat row meshed with 27,284 triangular elements.

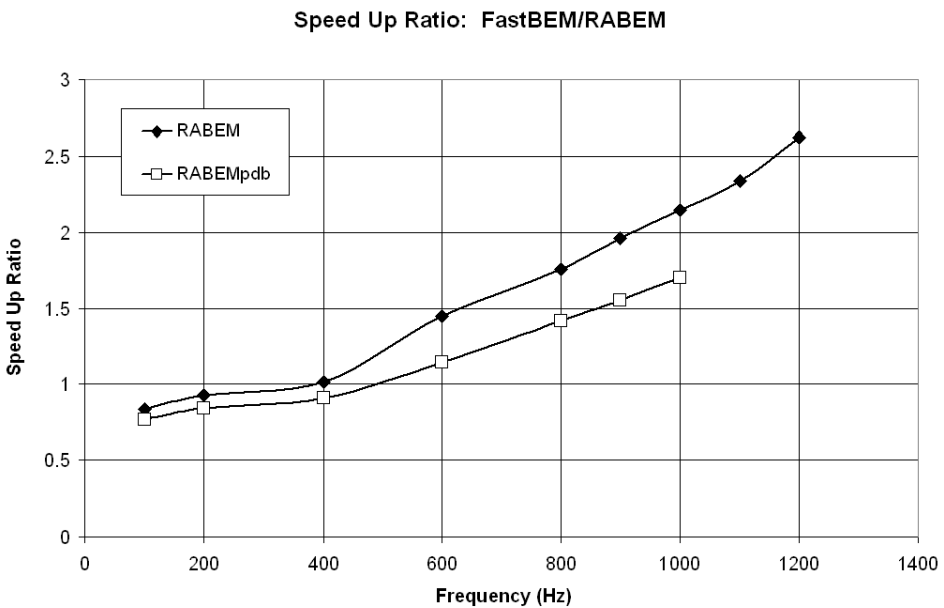


Figure 12: Comparison between the unpreconditioned, and block diagonal preconditioned GMRES for an airplane seat row meshed with 27,284 triangular elements.

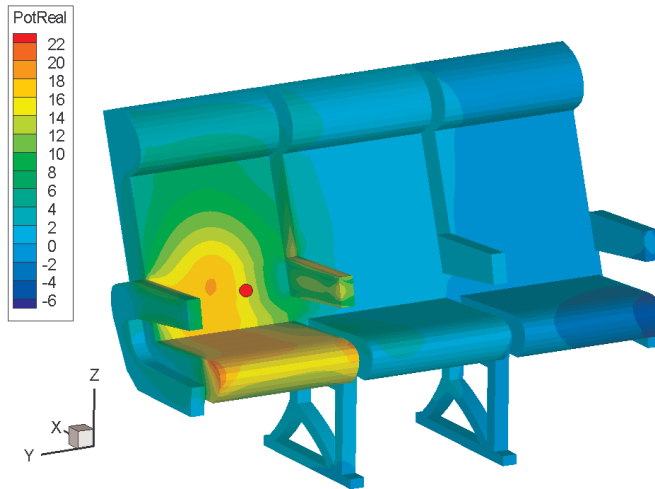


Figure 13: Real part of the potential for the aircraft seats when a monopole acts at 100Hz.

and the wing extension is 22.46m. The highest frequency applied was 180 Hz. The sizes of the smallest and the largest triangles are $1.02E-02$ and 0.23millimeter, respectively, so the maximum frequency that can be applied is around 185 Hz. In the simulation performed all the surfaces has been set as hard and two monopole sources with unit complex potential amplitude have been inserted just in front of the compressors of two engines.

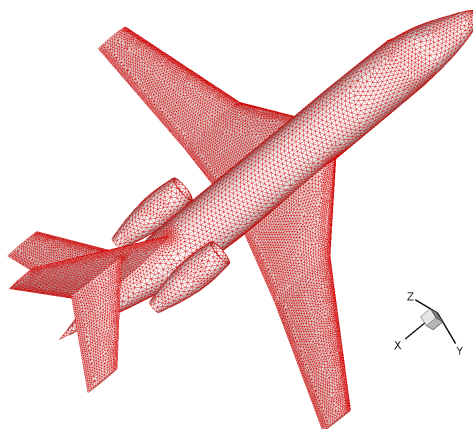


Figure 14: 47,998 triangular elements Dassault Falcon mesh.

Figure 15 shows the sound pressure level for 25Hz.

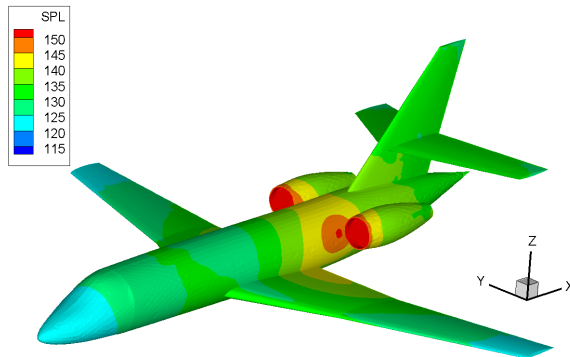


Figure 15: Sound pressure level at 25Hz for a Falcon geometry.

5 Conclusions

A new Hierarchical ACA GMRES BEM fast solver (RABEM) for 3D acoustic numerical simulations has been presented. The new approach is shown to result in significant savings in both storage and solution time. The proposed method is shown to speed up the solution time as much as 18 times the commercial code LMS. The solutions time are shown to be sensitive to the frequency with higher frequencies requiring more solution times. The tests demonstrated that the new solver can achieve CPU times of almost $O(N)$ for low frequency and $O(N \log N)$ for high frequency problems.

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Appendix

In this appendix a possible procedure to optimize the calculation of the right hand side vector is described.

In order to calculate the minimum number of terms four different alternatives which depend upon the boundary conditions need to be analyzed.

1. The values of the boundary conditions are all zero;
2. Application of ACA algorithm a second time;
3. Population of the missing columns;
4. Population of the missing rows.

The first case occurs when all the boundary conditions are zero. Thus there is no need to calculate the contribution of the block matrix to the final right hand side and another block matrix can be analyzed. The last three options are to optimize the number of the elements of the matrices that still need to be calculated.

Owing to the fact that the number of entries needed for ACA is, in most the cases, equal for both the block matrix \mathbf{G} and \mathbf{H} , the number of elements that must be calculated is

$$N_{ACA}^{II} = N_{ACA}^I \times m + N_{ACA}^I \times n \quad (11)$$

where N_{ACA}^I is the number of rows (or columns) calculated in the first application of the ACA.

Moreover, in the case that the non-zero boundary conditions are few, say N_{zBC} , only selected columns of the original block matrix are needed and the total number of elements to be calculated is

$$N_{dir1} = N_{zBC} \times m \quad (12)$$

The final option corresponds to calculating the remaining contribution of the right hand side block matrix directly and depends on the following numbers

$$N_{dir2} = n \times (m - N_{ACA}^I) \quad (13)$$

In conclusion, if the first alternative is not verified, the lower number between N_{ACA}^{II} , N_{dir1} and N_{dir2} drives the choice for calculating the right hand side contribution of each block.