

Finding the Generalized Solitary Wave Solutions within the (G'/G) -Expansion Method

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Abstract: In this study, the solitary wave solutions for third order equal-width wave-Burgers (EW-Burgers) equation, the second order Bratu and sinh-Bratu type equations will be discussed. The EW-Burgers equation models the propagation of nonlinear and dispersive waves with certain dissipative effects and furthermore the Bratu type problem appears a simplification of the solid fuel ignition model in thermal combustion theory. Our methodology, is investigated by using (G'/G) -expansion method. The obtained results can be extended to the other models.

Keywords: (G'/G) -expansion method, Nonlinear evolution equations, *EW*- Burgers equation, Bratu and sinh-Bratu type equations.

1 Introduction

It is well known that most of the phenomena that appear in physics can be described by partial differential equations Ebaid (2007), Bekir and Boz (2008), Yan (1996), Parkes (2010) and Ramos (2006). In this paper, the following problems will be investigated:

a. Canonical EW-Burgers equation as a special case of the generalized regularized long-wave (GRLW) equation which can be written as following

$$\frac{\partial}{\partial x} \left(\frac{\partial^2 u}{\partial t^2} \right) + \alpha \frac{\partial u}{\partial x} + \gamma \frac{\partial}{\partial t} \left(\frac{\partial u^p}{\partial x} \right) = \delta \frac{\partial^2 u}{\partial t^2} + \beta \frac{\partial}{\partial x} \left(\frac{\partial u}{\partial t} \right), (x, t) \subseteq \mathbb{R}^2, \quad (1)$$

where α , γ , δ and β are given non-negative real constants and $p \geq |t| + 2$, is a constant.

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The equal width wave equation is a model partial differential equation for the simulation of one-dimensional wave propagation in nonlinear media with dispersion processes. This equation is an important mathematical model arising in many different physical contexts to describe many phenomena which are simultaneously involved in nonlinearity, dissipation, dispersion, and instability, especially at the description of turbulence processes Ebaid (2007).

b. The nonlinear second order Bratu and sinh-Bratu type equations in the following form

$$\frac{d^2u}{dx^2} + \alpha \exp(u) + \beta \exp(-u) = 0, \quad \alpha, \beta \in \mathbb{R}. \quad (2)$$

The Bratu type problem can be used as the model of thermal reaction process, chemical reaction theory, radiative heat transfer and nanotechnology to the expansion of universe Wazwaz (2005), Abbasbandy, Hashemi, and Liu (2011) and Boyd (2011). The solving procedure of this method, by the help of Maple, Matlab, or any Mathematical package, is of utter simplicity.

Recently, several direct methods such as Exp-function method Bekir and Boz (2008) and Yan (1996), sine-cosine method Parkes (2010), tanh-coth method Wang, Zhou, and Li (1996), the homogeneous balance method Wang and Li (2005b), F-expansion method Wang and Li (2005a) and also wavelet methods Ray and Gupta (2014); Gupta and Ray (2014) have been proposed to obtain exact solutions of nonlinear partial differential equations. Using these methods many exact solutions, including the solitary wave solutions, shock wave solutions and periodic wave solutions are obtained for some kinds of nonlinear evolution equations.

The application of (G'/G) -expansion method to obtain more explicit traveling wave solutions to many nonlinear differential equations has been developed by many researchers Wang, Zhang, and Li (2008), Zayed and Gepreel (2009) Aslan and Öziş (2009) and Bekir (2008). The (G'/G) -expansion method is based on the assumption that the travelling wave solutions can be expressed by a polynomial in (G'/G) . It has been shown that this method is straightforward, concise, basic and effective. Eq. (1) and Eq. (2) for different values of $\alpha, \beta, \delta, \gamma$ and p are presented in **Table 1** and **Table 2** respectively. Finally, the paper is organized as follows. In the next section, the basic (G'/G) -expansion method is introduced. Application of G'/G -expansion method to our equations is presented in Section 3. Section 4 ends this work with a brief conclusion.

2 The basic (G'/G) -expansion method

We suppose that the given nonlinear partial differential $u(x, t)$ to be in the form

$$P(u, u_x, u_t, u_{xx}, u_{xt}, u_{tt}, \dots) = 0, \quad (3)$$

Table 1: Eq. (1) for different values of $\alpha, \beta, \delta, \gamma$ and p

Equation	Condition
First order linear wave	$\beta = \delta = \gamma = 0, \alpha \neq 0$
First order nonlinear wave	$\alpha = \beta = \delta = 0, \gamma \neq 0$
One dimensional heat transfer	$\alpha = \beta = \gamma = 0, \delta \neq 0$
On dimensional linear advection \ddot{I} -diffusion	$\beta = \gamma = 0, \alpha, \delta \neq 0$
One dimensional nonlinear advection \ddot{I} -diffusion	$\alpha = \beta = 0, \gamma, \delta \neq 0$
One dimensional nonlinear Burgers	$\alpha = \beta = 0, \gamma, \delta \neq 0, p = 2$
EW	$\alpha = \delta = 0, \gamma = 1, \beta \neq 0, p = 2$
RLW	$\alpha = 1, \delta = 0, \gamma, \beta \neq 0, p = 2$
EW-Burgers	$\alpha = 0, \gamma = 1, \delta, \beta \neq 0, p = 2$

Table 2: Eq. (2) for different values of $\alpha, \beta, \delta, \gamma$ and p

Equation	Condition
Bratu type	$\alpha \neq 0, \beta = 0$ or $\alpha = 0, \beta \neq 0$
Sinh-Bratu type	$\alpha = -\beta \neq 0$

where P is a multivariate polynomial in its arguments. In the following, it is explained the essential steps for implementing (G'/G) -expansion method.

Step 1. Taking the change of variable $\xi = x - wt$, gives $u(x, t) = U(\xi)$, where w is a constant parameter to be determined later. Substituting $\xi = x - wt$ into the Eq. (3) yields an ODE for $U(\xi)$ of the form

$$Q(U, U', -wU', U'', -wU'', w^2U'', \dots) = 0, \tag{4}$$

where $U^{(j)} = \frac{d^j U(\xi)}{d\xi^j}$. So, if possible, integrate Eq. (4), term by term one or more times. This introduces one or more constants of integration.

Step 2. Introduce the approach

$$U(\xi) = \sum_{j=0}^N \left(\frac{G'}{G}\right)^j, \tag{5}$$

where $G = G(\xi)$ satisfies the differential equation

$$G'' + \lambda G' + \mu G = 0, \tag{6}$$

here N is a positive integer (to be determined). The b_j , $j = 0, \dots, N$ and λ and μ are real constants with $b_N \neq 0$, and the prime denotes derivative with respect to ξ . (G'/G) satisfies the differential equation

$$\frac{d(\frac{G'}{G})}{d\xi} = -[\mu + \lambda(G'/G) + (G'/G)^2], \quad (7)$$

and so,

$$\frac{d}{d\xi} = -[\mu + \lambda(G'/G) + (G'/G)^2] \left(\frac{d}{(G'/G)} \right), \quad (8)$$

and

$$\begin{aligned} \frac{d^2}{d\xi^2} &= (\lambda + 2(\frac{G'}{G}))[\mu + \lambda(\frac{G'}{G}) + (\frac{G'}{G})^2] \left(\frac{d}{d(G'/G)} \right) \\ &+ [\mu + \lambda(\frac{G'}{G}) + (\frac{G'}{G})^2]^2 \left(\frac{d^2}{d(G'/G)^2} \right). \end{aligned} \quad (9)$$

Substituting Eq. (5) and (8) and (9) into the ODE from step 1, yields an algebraic equation in powers of the (G'/G) . Then, the positive integer N is determined by the balance of linear and nonlinear terms of the highest order in the resulting algebraic equation.

Step 3. With N being determined, the coefficients of each power of (G'/G) in the algebraic equation from Step 2 put equal to zero. This yields a system of algebraic equations involving b_j , $j = 0, \dots, N, w$ and the integration constants. Finally, the general solution of Eq. (6) is to be substituted into Eq. (5).

3 Application of G'/G -expansion method to EW-Burgers equation

To look for travelling wave solutions of Eq. (1), we use the wave transformation $\xi = x - wt$ and change Eq. (2) into the form of an ODE

$$-wU' + 2UU' - \delta U'' + w\beta U''' = 0. \quad (10)$$

Integrating it with respect to ξ and setting the constant of integration to zero, we obtain

$$-wU + U^2 - \delta U' + w\beta U'' = 0. \quad (11)$$

Now, we make an approach Eq. (6) for the solution of Eq. (10). Balancing the terms U^2 and U'' in Eq. (11), then we get $2N = N + 2$ which yields the leading term order $N = 2$. Therefore, we can write the solution of Eq. (11) in the form

$$U(\xi) = b_0 + b_1\left(\frac{G'}{G}\right) + b_2\left(\frac{G'}{G}\right)^2. \tag{12}$$

Substituting Eq. (12) into Eq. (11), collecting the coefficients of $\left(\frac{G'}{G}\right)^j, j = 0, \dots, 4$, and set it to zero we obtain the system of algebraic equations for b_0, b_1, b_2 and w . Then, solving the system by Mathematica 7., we obtain the following answers

$$\begin{aligned} b_0 &= \frac{6\delta\mu}{\lambda[\beta(\lambda^2 - 4\mu) - 1]}, & b_1 &= \frac{6\delta}{[\beta(\lambda^2 - 4\mu) - 1]}, \\ b_2 &= \frac{6\delta}{\lambda[\beta(\lambda^2 - 4\mu) - 1]}, & w &= -\frac{\delta(\lambda^2 - 4\mu)}{\lambda[\beta(\lambda^2 - 4\mu) - 1]}, \end{aligned} \tag{13}$$

$$\begin{aligned} b_0 &= -\frac{\delta(2\mu + \lambda^2)}{\lambda[\beta(\lambda^2 - 4\mu) + 1]}, & b_1 &= -\frac{6\delta}{[\beta(\lambda^2 - 4\mu) + 1]}, \\ b_2 &= -\frac{6\delta}{\lambda[\beta(\lambda^2 - 4\mu) + 1]}, & w &= -\frac{\delta(\lambda^2 - 4\mu)}{\lambda[\beta(\lambda^2 - 4\mu) + 1]}, \end{aligned} \tag{14}$$

where λ and μ are arbitrary constants. Substituting Eqs. (13) and (14) into Eq. (12) yields

$$U^\pm(\xi) = \begin{cases} \frac{6\delta}{\lambda[\beta(\lambda^2 - 4\mu) - 1]} \left[\mu + \lambda\left(\frac{G'}{G}\right) + \left(\frac{G'}{G}\right)^2 \right], \\ -\frac{6\delta}{\lambda[\beta(\lambda^2 - 4\mu) + 1]} \left[\lambda^2 + 2\mu + \lambda\left(\frac{G'}{G}\right) + \left(\frac{G'}{G}\right)^2 \right], \end{cases} \tag{15}$$

where $\xi = x - \left(-\frac{\delta(\lambda^2 - 4\mu)}{\lambda[\beta(\lambda^2 - 4\mu) \mp 1]}\right)t$.

Substituting the general solutions of Eq. (6) into Eq. (15) we have three types of travelling wave solutions of the EW-Burger equation as follows

1. when $\lambda^2 - 4\mu > 0$, we have

$$U_1^\pm(\xi) = \begin{cases} \frac{3\delta(\lambda^2 - 4\mu)}{2\lambda[\beta(\lambda^2 - 4\mu) - 1]} \left[\left(\frac{c_1 \sinh \frac{1}{2}\sqrt{\lambda^2 - 4\mu}\xi + c_2 \cosh \frac{1}{2}\sqrt{\lambda^2 - 4\mu}\xi}{c_1 \cosh \frac{1}{2}\sqrt{\lambda^2 - 4\mu}\xi + c_2 \sinh \frac{1}{2}\sqrt{\lambda^2 - 4\mu}\xi} \right)^2 - 1 \right], \\ -\frac{30\delta\mu}{\lambda[\beta(\lambda^2 - 4\mu) + 1]} - \frac{3\delta(\lambda^2 - 4\mu)}{2\lambda[\beta(\lambda^2 - 4\mu) - 1]} \left[\left(\frac{c_1 \sinh \frac{1}{2}\sqrt{\lambda^2 - 4\mu}\xi + c_2 \cosh \frac{1}{2}\sqrt{\lambda^2 - 4\mu}\xi}{c_1 \cosh \frac{1}{2}\sqrt{\lambda^2 - 4\mu}\xi + c_2 \sinh \frac{1}{2}\sqrt{\lambda^2 - 4\mu}\xi} \right)^2 + 3 \right], \end{cases} \tag{16}$$

2. when $\lambda^2 - 4\mu < 0$, we have

$$U_2^\pm(\xi) = \begin{cases} \frac{3\delta(4\mu - \lambda^2)}{2\lambda[\beta(4\mu - \lambda^2) + 1]} \left[\left(\frac{-d_1 \sin \frac{1}{2} \sqrt{4\mu - \lambda^2} \xi + d_2 \cos \frac{1}{2} \sqrt{4\mu - \lambda^2} \xi}{d_1 \cos \frac{1}{2} \sqrt{4\mu - \lambda^2} \xi + d_2 \sin \frac{1}{2} \sqrt{4\mu - \lambda^2} \xi} \right)^2 - 1 \right], \\ -\frac{30\delta\mu}{\lambda[1 - \beta(4\mu - \lambda^2)]} - \frac{3\delta(4\mu - \lambda^2)}{2\lambda[\beta(4\mu - \lambda^2) - 1]} \left[\left(\frac{-d_1 \sin \frac{1}{2} \sqrt{4\mu - \lambda^2} \xi + d_2 \cos \frac{1}{2} \sqrt{4\mu - \lambda^2} \xi}{d_1 \cos \frac{1}{2} \sqrt{4\mu - \lambda^2} \xi + d_2 \sin \frac{1}{2} \sqrt{4\mu - \lambda^2} \xi} \right)^2 + 3 \right], \end{cases} \quad (17)$$

3. when $\lambda^2 - 4\mu = 0$, we have

$$U_3^\pm(\xi) = \begin{cases} -\frac{6\delta}{\lambda} \left[\frac{1}{(\xi - x_0)^2} \right], \\ -\frac{6\delta}{\lambda} \left[\mu + \lambda^2 + \frac{1}{(\xi - x_0)^2} \right], \end{cases} \quad (18)$$

where c_1, c_2 in Eq. (16), d_1, d_2 in Eq. (17) and x_0 in Eq. (18) are arbitrary real constants.

4 Application of G'/G -expansion method to Bratu and sinh-Bratu type equations

4.1 The Bratu type equations

We first consider the Bratu type equations

$$\frac{d^2u(x)}{dx^2} + \alpha \exp(u(x)) = 0, \quad \alpha > 0. \quad (19)$$

To look for solutions of Eq. (28), we use the wave transformation $\xi = kx + x_0$ and change Eq. (28) into the form of an ODE

$$k^2 U'' + \alpha \exp(U(\xi)) = 0. \quad (20)$$

Using the transformation $U(\xi) = \ln v(\xi)$, will change Eq. (29) into the ODE in the form

$$k^2 (v'' v - (v')^2) + \alpha v^3 = 0, \quad (21)$$

where $v' = \frac{dv}{d\xi}$ and $v'' = \frac{d^2v}{d\xi^2}$. Now, we make an Eq. (6) for the solution of Eq. (17). Balancing the terms v^3 and vv'' in Eq. (17), then we get $3N = 2N + 2$ which yields the leading term order $N = 2$. Therefore, we can write the solution of v in the form

$$v(\xi) = b_0 + b_1 \left(\frac{G'}{G} \right) + b_2 \left(\frac{G'}{G} \right)^2, \quad (22)$$

substituting Eqs. (8), (9) and (13) into Eq. (12), collecting the coefficients of $(\frac{G'}{G})^j, j = 0, \dots, 4$, and set it to zero we obtain a system of algebraic equations for b_0, b_1, b_2 and k , that solving this system by Mathematica 7. gives

$$b_0 = -\frac{2\mu k^2}{\alpha}, b_1 = -\frac{2\lambda k^2}{\alpha}, b_2 = -\frac{2k^2}{\alpha}, k = k, \tag{23}$$

where λ and μ are arbitrary constants. Substituting Eq. (31) into Eq. (12) yields

$$v(\xi) = -\frac{2k^2}{\alpha}[\mu + \lambda(\frac{G'}{G}) + (\frac{G'}{G})^2], \tag{24}$$

where $\xi = kx + x_0$. Substituting the general solutions of Eq. (5) into Eq. (15) we have three types of solutions of the Bratu type equations as follows

1. When $\lambda^2 - 4\mu > 0$, we have

$$v(\xi) = -\frac{k^2(\lambda^2 - 4\mu)}{2\alpha} \left[\left(\frac{c_1 \sinh \frac{1}{2} \sqrt{\lambda^2 - 4\mu} \xi + c_2 \cosh \frac{1}{2} \sqrt{\lambda^2 - 4\mu} \xi}{c_1 \cosh \frac{1}{2} \sqrt{\lambda^2 - 4\mu} \xi + c_2 \sinh \frac{1}{2} \sqrt{\lambda^2 - 4\mu} \xi} \right)^2 - 1 \right], \tag{25}$$

2. When $\lambda^2 - 4\mu < 0$, we have

$$v(\xi) = \frac{k^2(4\mu - \lambda^2)}{2\alpha} \left[\left(\frac{c_1 \sin \frac{1}{2} \sqrt{4\mu - \lambda^2} \xi + c_2 \cos \frac{1}{2} \sqrt{4\mu - \lambda^2} \xi}{c_1 \cos \frac{1}{2} \sqrt{4\mu - \lambda^2} \xi + c_2 \sin \frac{1}{2} \sqrt{4\mu - \lambda^2} \xi} \right)^2 - 1 \right], \tag{26}$$

3. When $\lambda^2 - 4\mu = 0$, we have

$$v(\xi) = -\frac{2K^2}{\alpha} \left(\frac{1}{(\xi - x_0)^2} \right), \tag{27}$$

where $\xi = kx + x_0, c_1$ and c_2 in Eqs. (16) and (17), and $x_0 = -\frac{c_1}{c_2}$ in Eq. (18) are arbitrary constants. In particular, if we choose $c_2 \neq 0, c_1^2 < c_2^2$ then the solution Eq. (16) give the solution

$$v_1(\xi) = -\frac{k^2(\lambda^2 - 4\mu)}{2\alpha} \operatorname{sech}^2 \left(\frac{1}{2} \sqrt{\lambda^2 - 4\mu} \xi + \xi_0 \right), \tag{28}$$

where $\lambda^2 - 4\mu > 0$, $\xi_0 = \tanh^{-1} \frac{c_1}{c_2}$ and solution Eq. (17) give the following solution

$$v_2(\xi) = -\frac{k^2(\lambda^2 - 4\mu)}{2\alpha} \operatorname{csch}^2\left(\frac{1}{2}\sqrt{\lambda^2 - 4\mu}\xi + \xi_1\right), \tag{29}$$

where $\lambda^2 - 4\mu < 0$, $\xi_0 = \tanh^{-1} \frac{c_1}{c_2}$ and recall that $u(x) = U(\xi) = \ln(v(\xi))$, hence we obtain the following solutions of the Bratu type equations from Eqs. (28) and (29) as follows, respectively

$$u_1(x) = \ln\left[\frac{k^2(\lambda^2 - 4\mu)}{2\alpha} \operatorname{sech}^2\left(\frac{1}{2}\sqrt{\lambda^2 - 4\mu}\xi + \xi_0\right)\right], \tag{30}$$

where $\lambda^2 - 4\mu > 0$ and $\xi_0 = \tanh^{-1} \frac{c_1}{c_2}$,

$$u_2(x) = \ln\left[-\frac{k^2(\lambda^2 - 4\mu)}{2\alpha} \operatorname{csch}^2\left(\frac{1}{2}\sqrt{\lambda^2 - 4\mu}\xi + \xi_0\right)\right], \tag{31}$$

where $\lambda^2 - 4\mu < 0$, $\xi_1 = \tanh^{-1} \frac{c_1}{c_2}$, and another solution from Eq. (28) is as follows:

$$u_3(x) = \ln\left[-\frac{2}{\alpha} \left(\frac{1}{x^2}\right)\right]. \tag{32}$$

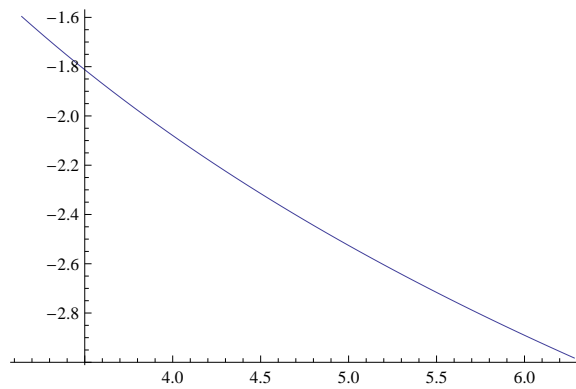


Figure 1: Approximate solution of Eq. (20) where $\alpha = -1$ and $x \in [1, 2]$.

4.2 The sinh-Bratu type equations

We secondly consider sinh-Bratu type equations

$$\frac{d^2u(x)}{dx^2} + \alpha \sinh(u(x)) = 0. \quad \alpha > 0. \tag{33}$$

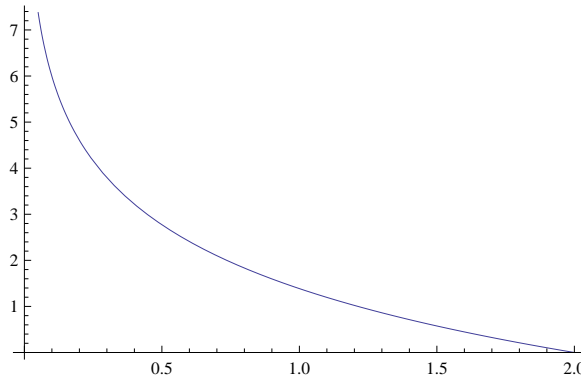


Figure 2: Approximate solution of Eq. (20) where $\alpha = -1.5$ and $x \in [1, 2]$.

In order to apply the (G'/G) method, we use the transformation $\xi = kx + x_0$ and change Eq. (33) into the form

$$k^2U'' + \alpha \sinh(U) = 0. \tag{34}$$

And then we use the transformation $U(\xi) = \ln v(\xi)$, so that

$$\sinh(U(\xi)) = \frac{v(\xi) - v^{-1}(\xi)}{2}, \cosh(U(\xi)) = \frac{v(\xi) + v^{-1}(\xi)}{2}, \tag{35}$$

this transformation will change Eq. (34) into the ODE in the form

$$2k^2(v''v - (v')^2) + \alpha(v^3 - v) = 0, \tag{36}$$

where $v' = \frac{dv}{d\xi}$, $v'' = \frac{d^2v}{d\xi^2}$. Now, we make an Eq. (6) for the solution of Eq. (36). By balancing the terms v^3 and $v v''$ in Eq. (36), then we get $3N = 2N + 2$ which yields the leading term order $N = 2$. Therefore, we can write the solution of v in the form

$$v(\xi) = b_0 + b_1\left(\frac{G'}{G}\right) + b_2\left(\frac{G'}{G}\right)^2, \tag{37}$$

substituting Eqs. (8), (9) and (37) into Eq. (36), collecting the coefficients of $(\frac{G'}{G})^j$, $j = 0, \dots, 4$, and set it to zero we obtain a system of algebraic equations for b_0, b_1, b_2 and k . By solving the resulted system with the help of Mathematica 7. we have the following sets of solutions

$$b_0 = \pm \frac{\lambda^2}{\lambda^2 - 4\mu}, b_1 = \pm \frac{4\lambda}{\lambda^2 - 4\mu}, b_2 = \pm \frac{4}{\lambda^2 - 4\mu}, k = \pm \sqrt{\frac{\alpha}{|\lambda^2 - 4\mu|}}, \tag{38}$$

where λ, μ and $\lambda^2 \neq 4\mu$ are arbitrary constants. Substituting Eq. (38) into Eq. (37) yields

$$v(\xi) = \pm \frac{1}{\lambda^2 - 4\mu} \lambda^2 + 4\lambda \left(\frac{G'}{G} \right) + 4 \left(\frac{G'}{G} \right)^2, \tag{39}$$

where $\xi = kx + x_0$. By substituting the general solutions of Eq. (5) into Eq. (39) we have three types of solutions of the sinh-Bratu type equations as follows

1. When $\lambda^2 - 4\mu > 0$, we have

$$v_1^\pm(\xi) = \pm \left[\left(\frac{c_1 \sinh \frac{1}{2} \sqrt{\lambda^2 - 4\mu} \xi + c_2 \cosh \frac{1}{2} \sqrt{\lambda^2 - 4\mu} \xi}{c_1 \cosh \frac{1}{2} \sqrt{\lambda^2 - 4\mu} \xi + c_2 \sinh \frac{1}{2} \sqrt{\lambda^2 - 4\mu} \xi} \right)^2 \right], \tag{40}$$

where $\xi = \pm \sqrt{\mp \frac{\alpha}{\lambda^2 - 4\mu}} x + x_0$, c_1 and c_2 are arbitrary constants.

2. When $\lambda^2 - 4\mu < 0$, we have

$$v_2^\mp(\xi) = \mp \left[\left(\frac{-c_1 \sin \frac{1}{2} \sqrt{4\mu - \lambda^2} \xi + c_2 \cos \frac{1}{2} \sqrt{4\mu - \lambda^2} \xi}{c_1 \cos \frac{1}{2} \sqrt{4\mu - \lambda^2} \xi + c_2 \sin \frac{1}{2} \sqrt{4\mu - \lambda^2} \xi} \right)^2 \right], \tag{41}$$

where $\xi = \pm \sqrt{\pm \frac{\alpha}{\lambda^2 - 4\mu}} x + x_0$, c_1 and c_2 are arbitrary constants.

In particular, if we choose $c_2 \neq 0, c_1^2 < c_2^2$ then the solutions Eqs. (40) and (41) give the following solutions

$$v_1^\pm(\xi) = \pm \tanh^2 \left(\frac{1}{2} \sqrt{\lambda^2 - 4\mu} \xi + \xi_0 \right), \tag{42}$$

where $\lambda^2 - 4\mu > 0, \xi_0 = \tanh^{-1} \left(\frac{c_1}{c_2} \right)$ and

$$v_2^\mp(\xi) = \mp \coth^2 \left(\frac{1}{2} \sqrt{4\mu - \lambda^2} \xi + \xi_1 \right), \tag{43}$$

where $\lambda^2 - 4\mu < 0, \xi_1 = \tan^{-1} \left(\frac{c_1}{c_2} \right)$ and recall that $u(x) = U(\xi) = \ln(v(\xi))$, hence we obtain the some solutions of the sinh-Bratu type equations from Eqs. (42) and (43) as follows, respectively

$$u_1^\pm(x) = \text{Arccos} \left[\pm \frac{1}{2} \left(\tanh^2 \left(\frac{1}{2} \sqrt{\lambda^2 - 4\mu} \xi + \xi_0 \right) \right) \right] + \left(\coth^2 \left(\frac{1}{2} \sqrt{\lambda^2 - 4\mu} \xi + \xi_0 \right) \right), \tag{44}$$

where $\lambda^2 - 4\mu > 0, \xi_0 = \tanh^{-1}(\frac{c_1}{c_2}),$

$$u_2^\mp(x) = \text{Arccos}[\mp \frac{1}{2}(\tan^2(\frac{1}{2}\sqrt{4\mu - \lambda^2}\xi + \xi_1))] + (\cot^2(\frac{1}{2}\sqrt{4\mu - \lambda^2}\xi + \xi_1)), \tag{45}$$

where $\lambda^2 - 4\mu < 0, \xi_1 = \tan^{-1}(\frac{c_1}{c_2}).$

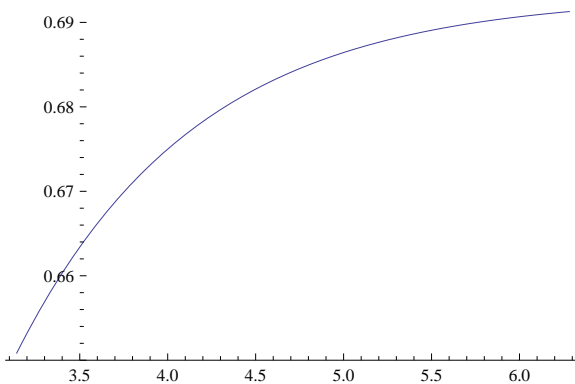


Figure 3: Approximate solution of Eq. (33) where $\mu = 0.125, \lambda = 0.7, k = 1, x_0 = 0, \xi_0 = 1$ and $x \in [\pi, 2\pi]$

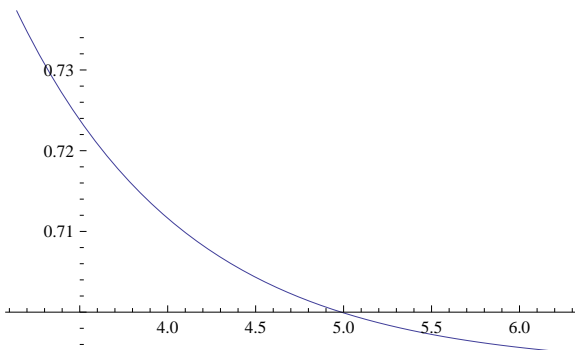


Figure 4: Approximate solution of Eq. (33) where $\mu = 0.125, \lambda = 0.7, k = 1, x_0 = 0, \xi_1 = 1$ and $x \in [\pi, 2\pi].$

Remark : The computations associated in Figures [1-4] were performed by using Mathematica 7.

5 Conclusions

In this paper, we apply the (G'/G) -expansion method to third order EW-Burgers equation and Bratu and sinh-Bratu type equations. Our results show that the (G'/G) -expansion method is entirely efficient and well suited for finding exact solutions of these equation. The advantage of this method over other methods is that we can obtain the exact solution by using a simple computer program.

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