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# Numerical Examination of Free Convection Flow of Casson Ternary Hybrid Nanofluid across Magnetized Stretching Sheet Impacted by Newtonian Heating

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Received: 27 July 2023 Accepted: 28 September 2023 Published: 30 November 2023

## ABSTRACT

In current study, the influence of magnetic field (MHD) on heat transfer of natural convection boundary layer flow in Casson ternary hybrid nanofluid past a stretching sheet is studied using numerical simulation. The Newtonian heating boundary conditions that depend on the temperature and velocity terms are taken into this investigation. The particular dimensional governing equations, for the studied problem, are converted to the system of partial differential equations utilizing adequate similarity transformation. Consequently, the system of equations is numerically solved using well-known Keller box numerical techniques. The obtained numerical results are in excellent approval with previous literature results. The existence of types for composite ternary nanoparticles, which as alumina  $Al_2O_3$ , copper oxide  $CuO$ , and silver  $Ag$  or copper  $Cu$ , as well as water  $H_2O$  which is considered an essential fluid, are examined. The physical features, such as temperature, velocity, local skin friction coefficient, and local Nusselt number are explored with the variation of MHD Casson ternary hybrid nanofluids parameters and displayed graphically and tabularly in the present study. The numerical outcomes exhibit that when the conjugate parameter values rise, the temperature and local Nusselt number increase, however they decrease as the Casson parameter is reduced.

## KEYWORDS

Casson ternary hybrid nanofluid fluid; magnetic field; stretching sheet; newtonian heating

## Nomenclature

$P_y$	Defined the yield stress
$B_0$	Magnetic field
$C_f$	Skin friction coefficient
$\tau_{ij}$	Rheological property



$Re$	Reynold number
$g$	Gravity vector
$\chi$	Volume fraction parameter
$M$	Magnetic parameter
$Nu$	Nusselt number
$Pr$	Prandtl number
$\sigma$	Electrical conductivity
$x, y$	Dimensional variables along velocity component $u, v$
$\pi_c, \mu_\beta$	Critical value of product based, the plastic dynamic viscosity
$\theta$	Temperature
$\psi$	Stream function
$e$	Positive constant
$Hs$	Conjugate parameter
$hs$	Heat transfer parameter
$U$	Component of velocity
$v$	Component of velocity
$\nu_f$	Kinematic viscosity
$k$	Thermal conductivity
$\beta$	Casson parameter
$T_w$	Wall temperature
$\mu$	Dynamic viscosity
$\rho$	Density
$T_\infty$	Ambient temperature
$\rho c_p$	Heat capacity
$\pi$	Component of the deformation rate

### Subscript

$f$	Based fluid
$nf$	Nanofluid
$hnf$	Hybrid nanoliquid
$thnf$	Ternary hybrid nanoliquid

## 1 Introduction

In fluid dynamics, the models that relate shear stress to shear rate with a linear relationship are insufficient to describe the behavior of complicated fluids; however they are particularly helpful in this field. Models that are common in use to describe the rheological behavior of such fluids are based on Casson model, which is firstly presented in 1959 by Casson [1]. There are two crucial parameters included in the model: The Casson yield stress, which represents the minimal stress needed to onset the flow, and Casson viscosity, which reflects the fluid's resistance to deformation once it has started flow. In recent years, this model has received attention due to its numerous applications. For example, simulation of fluid flows of drilling fluids and the behavior of inks and paints in printing sectors. In the literature, different numerical studies have investigated the flow behavior of the Casson fluid, highlighting its potential applications in various fields. The study by Hazarika et al. [2]

investigated numerically the behavior of a steady MHD micropolar Casson fluid and its heat transfer characteristics as it flowed over a spherical object. They reported that the micropolar factor increases the temperature, while Casson factor decreases it. These results could be used in applications of tissue engineering and bioengineering. A recent numerical study presented by Gireesha et al. [3] analyzed the impact of chemical reactions, a magnetic field, and a space-dependent energy source on the thermal behavior of Casson liquid moving across a curved stretch sheet. Their findings indicated that the increase of the Casson factor is resulted in substantial flow profile variations. Moreover, the study reported that temperature is an increasing function in terms of thermal Biot number, while concentration is an increasing function in terms of concentration Biot number. A study by Nagaraja et al. [4] carried out the flow and melting energy transport across a stretching sheet for a magneto dusty Casson fluid, focusing on the Cattaneo-Christov heat flux model. They found that, elevating the Casson parameter raises the Nusselt number and reduces the drag force. More comprehensive studies can be found in [5,6].

Nanoparticles have emerged as promising materials for enhancing heat transport owing to their distinctive characteristics, which include a high surface area to volume ratio, tunable size and shape, and relatively high thermal conductivity. When dispersed in a fluid, nanoparticles can increase the thermal conductivity of the original liquid, leading to improve heat transport performance. The first known study to introduce the use of nanoparticles for enhancing energy transport was conducted by Choi et al. in 1995 [7]. They proposed the concept of nanofluids, which are fluids containing nanoparticles. In their seminal paper, Choi and Eastman demonstrated that the inclusion of nanosolids in an original fluid could dramatically improve its thermal conductivity, thereby supporting its thermal performance. Their work sparked a lot of interest in the scientific community, and studies continued, confirming that the use of these unique nanoparticles in energy transport applications has the potential to improve heat efficiency and reduce costs in various industries [8–10]. Their work still inspires researchers to explore the potential of nanoparticles and optimizing them to enhance the process of energy transport. Hybrid nanoparticles have emerged as a promising approach for enhancing the heat transfer of fluids in various engineering applications. By combining two or more types of nanosolids in a single liquid, a hybrid nanoliquid is obtained with unique thermal and rheological properties. Indeed, this will enhance thermal conductivity and stability and reduce agglomeration compared to mono-nanofluids. Additionally, hybrid nanoparticles can alter rheological properties, such as viscosity and density, affecting fluid flow patterns and heat transfer rates [11–16]. However, for further enhancement of the performance of nanoliquids, researchers have developed the third generation of nanoliquids, known as ternary hybrid nanoliquid. It consists of three different kinds of nanosolids dispersed in a fluid and offer superior energy transport properties compared to their mono and hybrid nanoparticle counterparts [17–20]. The thermal capabilities of these ternary hybrid nanofluids have been scrutinized through a multitude of numerical investigations. Manjunatha et al. [21] explored the convective energy transport characteristics of ternary nanofluids passing over a stretching sheet. Their research revealed that the thermal conductivity of the ternary nanofluid surpasses that of the hybrid nanofluid, highlighting the enhanced heat transfer capabilities of the ternary nanofluid. In their numerical study, Algehyne et al. [22] delved the behavior of ternary hybrid nanoliquid flowing over a stretched permeable surface. By employing a pseudoplastic model and the Darcy-Forchheimer relation, they uncovered significant enhancements in velocity, energy transport rate, and thermal conductivity, varying as much as 11.73%, 32%, and 61%, depending on the type of nanocomposite used. Guedri et al. [23] carried out a computational analysis to examine the flow of a tri-hybrid

nanofluid across a stretching sheet subjected to Joule heating and viscous dissipation. They observed that the magnetic factor and electric field had a considerable effect on fluid flow, and these factors supported the Bejan number and reduced the production of entropy in the system. Related research can be seen in references [24–33].

Numerous industrial and engineering heat-transmitting operations critically rely on fluid movement over a stretching sheet. In the polymer industry, film casting involves stretching a polymer sheet while simultaneously cooling it with fluid flow to prevent overheating and maintain its desired properties. Similarly, in the glass industry, fluid flow is utilized to deftly control the cooling rate of stretched molten glass sheets, ensuring they meet precise thickness and smoothness standards. Wire drawing also benefits from fluid flow to cool wires during stretching, which improves their surface finish for use in electrical and construction applications. On the other hand, fluid flow is employed to regulate the drying rate of stretched paper sheets, resulting in desirable properties such as strength, thickness, and smoothness which is an essential process in paper production and manufacturing. By ensuring that stretching sheets meet stringent industry standards, fluid flow plays an indispensable role in elevating the quality of industrial processes. On the other hand, the interaction between magnetic fields and conductive fluids, such as plasma and liquid metals, in the context of heat transfer through fluids, or what is known as magnetic hydrodynamics (MHD), can have a significant impact on fluid flow patterns and heat transfer rates. In the presence of a magnetic field, the motion of a conductive fluid can be affected by the Lorentz force, which arises from the interaction between the magnetic field and the current in the fluid. This force can cause movement of fluid in different ways than it would in the absence of a magnetic field. MHD has numerous applications in heat transfer through fluids, ranging from the design of fusion reactors, where its effects can play an important role in determining the efficiency of energy transfer between the plasma and the reactor walls, to the cooling of nuclear reactors by improving the rate of heat transfer between the coolant and the reactor walls. This is important to prevent overheating and maintain safe operating conditions. MHD can also be used in many industrial processes, such as metal casting and crystal growth, to improve heat transfer efficiency and control the flow of a conductive liquid. In the presence of these enormous applications of fluid flow through a stretched plate under the influence of a magnetic field, wide numerical studies have tackled this problem in an attempt to improve the thermal performance of these fluids as well as gain a deep understanding of the factors affecting that performance. Sohail et al. [34] examined the role of tri-hybrid nanosolids in a mixture of pseudo-plastic liquids over a 2D porous stretched surface. Fatima et al. [35] conducted a study to examine the convective heat transport in hydromagnetic  $H_2O$  nanoliquid,  $H_2O$  hybrid nanoliquid, and  $H_2O$  tri hybrid nanoliquid flow while subjecting it to the influence of Cattaneo-Christov heat flux and nonlinear thermal radiation over a stretching sheet [36]. The impact of Hall current and thermal radiation on the movement of magnetohydrodynamic liquid carrying tiny particles across a nonlinear stretching surface was investigated by Ali et al. [37].

The investigation focuses on the numerical analysis of Casson fluid behavior, considering the non-Newtonian characteristics inherent in many real-world fluids. This aspect is of paramount importance as it allows for a more accurate representation of the fluid flow and heat transfer phenomena encountered in practical engineering applications. The study incorporates a ternary hybrid nanofluid composition, comprising a mixture of three distinct nanoparticles. This aspect introduces a new dimension to the research as it explores the synergistic effects of multiple nanoparticles on the convective heat transfer characteristics of the fluid. The examination of such complex nanofluid compositions holds significant potential for advancing the understanding and optimization of heat transfer processes in various industrial domains. Moreover, the investigation accounts for the influence of a magnetic field on the fluid flow, thus incorporating magnetohydrodynamic effects. The inclusion

of the magnetic field as a governing parameter is crucial, as it allows for the analysis of the coupled interactions between fluid flow, heat transfer, and the magnetic field. Understanding these interactions is vital for the design and optimization of magnetohydrodynamically-driven systems, with applications in areas such as energy conversion and thermal management. Lastly, the study investigates the impact of Newtonian heating, which introduces an additional heat source to the system. By considering this aspect, the research contributes to a deeper understanding of convective heat transfer phenomena in the presence of Newtonian heating, which has practical implications for a wide range of industrial processes. The outcomes of this research endeavor are anticipated to enhance the fundamental knowledge base and pave the way for advancements in the design and optimization of heat transfer systems in various engineering applications. More specifically, the current investigation can be summarized via the following questions:

1. How do the governing equations for Casson tri-hybrid nanoliquid, under magnetic field, models produce in the presence of unsteady free convection flow over a stretching sheet?
2. How can a mathematical formulation for the problem of MHD Casson tri-hybrid nanoliquid be introduced over a stretching sheet?
3. How do the MHD Casson tri-hybrid nanoliquid models compare with the previously investigated natural heat and mass transfer flow problems?
4. How do analysis of the numerical results that can be gained from the influences of MHD Casson tri-hybrid parameters on the heat transfer and physical quantities?
5. How does the heat transfer demeanor of the employed nanoparticles that are suspended in the based fluid under the impact of considered parameters?

## 2 Mathematical Formulation

In this study, a two dimensional natural convection boundary layer flow is investigated. The physical model under consideration is a stretching sheet with an incompressible magnetohydrodynamic Casson ternary hybrid nanofluid with a Newtonian heating on the boundaries. Two types of Casson ternary hybrid nanofluid are taken into account. The first type is ternary hybrid nanofluid compositions  $\text{Al}_2\text{O}_3\text{-CuO-Ag}$  and  $\text{Al}_2\text{O}_3\text{-CuO-Cu}$  nanoparticles, and the second type is a hybrid nanofluid composition  $\text{Al}_2\text{O}_3\text{-CuO}$ . Magnetic field  $\beta_0^2$  intensity is employed, and it varies the orientation of the ternary hybrid nanofluid flow. Additionally, the ternary hybrid nanofluids acquire a characteristic through the Casson environment. The Cartesian frame is chosen in such a way that the fluid flow in the  $x$ -direction, which is gauged over in the horizontal direction of the stretching sheet. Also, the  $y$ -axis is the distance orthogonal on the considered stretching sheet, as displayed in Fig. 1. The form of Newtonian heating is applied to the wall. In the energy equation, the viscous dissipation effect was assumed to be neglected. The influence of magnetic, electrical conductivity, Casson, and nanoparticle volume fraction are included in the momentum equation.

The rheological property of Casson fluid is given as follows (see Casson [1]):

$$\tau_{ij} = \begin{cases} 2 \left( \mu_\beta + \frac{P_y}{\sqrt{2\pi}} \right) e_{ij, \pi > \pi_i} \\ 2 \left( \mu_\beta + \frac{P_y}{\sqrt{2\pi_c}} \right) e_{ij, \pi < \pi_i} \end{cases}, \quad (1)$$

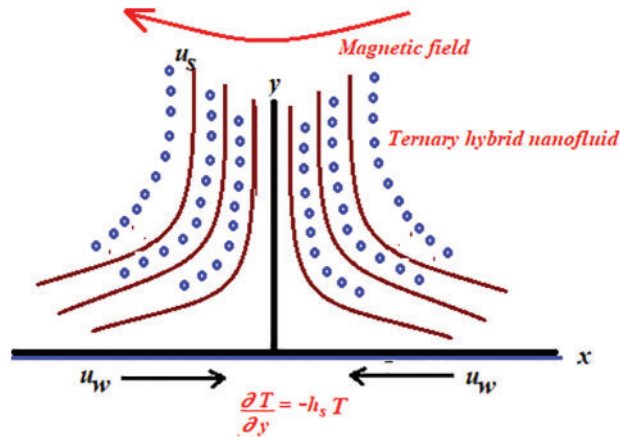
Here, the symbol  $P_y$  is defined the yield stress of the fluid.  $\pi_c$ ,  $\mu_\beta$  are called critical value of this product based and the plastic dynamic viscosity of the non-Newtonian model, respectively.  $\pi = e_{ij}e_{ij}$ ,  $e_{ij}$  is the component of the deformation rate.

In the state of the Casson fluid, case  $\pi > \pi_c$ , we get:

$$\mu = \mu_\beta + P_y \sqrt{2\pi}, \quad (2)$$

where  $P_y = \mu_\beta \sqrt{2\pi} / \beta$ . Then we get:

$$\frac{\mu}{\rho} = \frac{\mu_\beta}{\rho} + \left[ 1 + \frac{1}{\beta} \right], \quad (3)$$



**Figure 1:** Physical model and coordinate system

Upon all the above regarding, the ternary hybrid nanofluids model, regarding magnetic, and Casson impacts, the governing equations, namely (continuity, momentum, and thermal) can be established as [38–42]:

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0, \quad (4)$$

$$u \frac{\partial u}{\partial x} + v \frac{\partial v}{\partial x} = \nu_{THnf} \left( 1 + \frac{1}{\beta} \right) \frac{\partial^2 u}{\partial y^2} - \frac{\sigma_{THnf}}{\rho_{THnf}} B_0^2 u, \quad (5)$$

$$u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \frac{K_{THnf}}{(\rho C_p)_{THnf}} \frac{\partial^2 T}{\partial y^2}, \quad (6)$$

according to Newtonian heating (NH) boundary conditions, according to Newtonian heating (NH) boundary conditions, can be written as [42]:

$$\begin{aligned} u = u_w = ax, v = 0, T = -hs \frac{\partial T(x)}{\partial y}, \text{ at } y = 0 \\ u \rightarrow 0, T \rightarrow T_\infty \text{ as } y \rightarrow \infty \end{aligned} \quad (7)$$

$T$  is referred to the temperature;  $u$  and  $v$  are symbolized as the velocity components which depend on the  $x$  and  $y$  directions.  $\beta_0^2$  is the magnetic field strength,  $hs$  and  $\rho_s$  are heat transfer parameter,

thermal expansion coefficient, and density of nanoparticles, respectively.  $\rho_f, \nu_f, \beta$  are density, kinematic viscosity, and Casson parameter, respectively.  $\alpha_{THnf}, \rho_{THnf}, (\rho c_p)_{HTnf}, \mu_{THnf}, k_{THnf}$  and  $\sigma_{THnf}$  are thermal diffusivity, density, heat capacity, viscosity, thermal conductivity, and electrical conductivity of ternary hybrid nanofluid, respectively, and they are given in [Table 1](#), where  $\chi_1, \chi_2,$  and  $\chi_3$  are defined the nanoparticle volume fraction parameter refer to  $Al_2O_3, CuO,$  and  $Ag$  or  $Cu,$  respectively.  $\mu_f, k_f$  and  $\sigma_{nf}$  are viscosity, thermal and electrical conductivity, respectively of base fluid.

**Table 1:** Thermo-physical properties of Tri-hybrid nanofluid [43]

$(\mu)_{HTnf} = \frac{\mu_f}{(1 - \chi_1)^{2.5}(1 - \chi_2)^{2.5}(1 - \chi_3)^{2.5}},$
$(\rho)_{HTnf} = (1 - \chi_1) [(1 - \chi_2) [(1 - \chi_3) \rho_f + \chi_3 \rho_3] + \chi_2 \rho_2] + \chi_1 \rho_1,$
$(\rho c_p)_{HTnf} = (1 - \chi_1) [(1 - \chi_2) [(1 - \chi_3) (\rho c_p)_f + \chi_3 (\rho c_p)_3] + \chi_2 (\rho c_p)_2] + \chi_1 (\rho c_p)_1,$
$\frac{k_{Hnf}}{k_{Hnf}} = \frac{(k_1 + 2k_{hnf}) - 2\chi_1(k_{hnf} - k_1)}{(k_1 + 2k_{hnf}) + \chi_1(k_{hnf} - k_1)},$
$\frac{k_{Hnf}}{k_{Hnf}} = \frac{(k_2 + 2k_{nf}) - 2\chi_2(k_{nf} - k_2)}{(k_2 + 2k_{nf}) + \chi_2(k_{nf} - k_2)},$
$\frac{k_{nf}}{k_f} = \frac{(k_3 + 2k_f) - 2\chi_3(k_f - k_3)}{(k_3 + 2k_f) + \chi_3(k_f - k_3)},$
$\sigma_{HTnf} = \left[ 1 + \frac{3((\sigma_1/\sigma_{Hnf}) - 1) \chi_1}{((\sigma_1/\sigma_{Hnf}) + 2) - \chi_1((\sigma_1/\sigma_{Hnf}) - 1)} \right] \sigma_{Hnf},$
$\sigma_{Hnf} = \left[ 1 + \frac{3((\sigma_2/\sigma_{nf}) - 1) \chi_1}{((\sigma_2/\sigma_{nf}) + 2) - \chi_2((\sigma_2/\sigma_{nf}) - 1)} \right] \sigma_{nf},$
$\sigma_{nf} = \left[ 1 + \frac{3((\sigma_3/\sigma_f) - 1) \chi_3}{((\sigma_3/\sigma_f) + 2) - \chi_3((\sigma_3/\sigma_f) - 1)} \right] \sigma_f,$

Now, consider the following similarity transformations:

$$v = -\frac{\partial \psi}{\partial x}, u = \frac{\partial \psi}{\partial y},$$

$$\psi = (ev)^{\frac{1}{2}}x, f(\eta)\eta = (ev)^{\frac{1}{2}}y, \theta(\eta) = \frac{T - T_\infty}{T_w - T_\infty}, \tag{8}$$

where  $\psi$  represents the stream function.

Substitute [Eqs. \(7\)](#) and [\(8\)](#) and the resultant properties in [Table 1](#) into [Eqs. \(4\)–\(6\)](#), we get the following partial differential equations:

$$\frac{\rho_f}{\rho_{thnf}} \left[ \frac{1}{(1 - \chi_1)^{2.5} (1 - \chi_2)^{2.5} (1 - \chi_3)^{2.5}} \right] \left( 1 + \frac{1}{\beta} \right) f'''(\eta) + f(\eta)f''(\eta) - (f'(\eta))^2 - \frac{\rho_f}{\rho_{thnf}} \frac{\sigma_{thnf}}{\sigma_f} Mf'(\eta) = 0, \tag{9}$$

$$\frac{1}{Pr} \left( \frac{k_{thnf}/k_f}{(1-\chi_1) \left[ (1-\chi_2) \left[ (1-\chi_3) + \chi_3 \frac{(\rho c_p)_3}{(\rho c_p)_f} \right] + \chi_2 \frac{(\rho c_p)_2}{(\rho c_p)_f} \right] + \chi_1 \frac{(\rho c_p)_1}{(\rho c_p)_f}} \right) (\theta''(\eta)) + f(\eta)\theta'(\eta) = 0, \quad (10)$$

and the boundary conditions

$$\begin{aligned} f(0) = 0, f'(0) = 1, \theta'(0) = -Hs - \theta(0), \text{ at } \eta = 0, \\ f'(\eta) \rightarrow 0, \theta(\eta) \rightarrow 0 \text{ at } \eta \rightarrow \infty, \end{aligned} \quad (11)$$

where  $Pr = \frac{v_f}{k_f}$ ,  $M = \frac{\sigma_f B_0^2 e^2}{\rho_f v_{thnf}}$ ,  $Hs = \left(\frac{v_f}{a}\right)^{1/2} h_s$  are the Prandtl number, magnetic parameter, the conjugate parameter, respectively. It is worth noting that Newtonian heating is ruled by a non-dimensional conjugate parameter, which has a minimum value of 0 (insulated wall) to  $\infty$  (wall temperature remains constant). Furthermore, the position with Newtonian heating occurs in what is usually termed conjugate convection flow, where the heat is provided to the convective fluid via a bounding surface with a finite heat capacity. This composition happens in numerous necessary engineering machines, for example in heat exchangers where the conduction in the convection significantly impacts solid tube wall in the fluid flowing over it.

Due to engineering purposes, the local skin friction coefficient  $C_f$  and the Nusselt number  $Nu$  of the current model can be expressed as

$$\begin{aligned} C_f &= \frac{\tau_w}{\rho u_w^2}, \\ Nu &= \frac{xq_w}{k(T_w - T_\infty)}, \end{aligned} \quad (12)$$

where  $\tau_w$  and  $q_w$  correspond to shear stress and heat flux on the plane of the wall, and are given by

$$q_w = k_{THnf} \frac{\partial T}{\partial y}, \tau_w = \mu_{THnf} \left( \frac{\partial u}{\partial y} \right) \text{ at } y = 0, \quad (13)$$

As a result, the local skin friction coefficient and Nusselt number become

$$\begin{aligned} C_f &= Re^{-\frac{1}{2}} \left[ \frac{1}{(1-\chi_1)^{2.5} (1-\chi_2)^{2.5} (1-\chi_3)^{2.5}} \right] \left( 1 + \frac{1}{\beta} \right) f''(0), \\ Nu &= Re^{\frac{1}{2}} \frac{k_{THnf}}{k_f} (\theta'(0)), \end{aligned} \quad (14)$$

where  $Re = \frac{ax^2}{v_f}$  is the local Reynolds number.



### 3 Numerical Method and Accuracy

The generated equations obtained in previous section, are solved numerically using Keller-box scheme. This numerical method relies on steps and methods as follows: Reducing equations with high order partial derivative to the first order using finite difference method. Then the resultant equations are centered using centering method, which gives the difference equations. In third step, the gained difference equations are linearized using Newton’s method, consequently, the linear system is solved by the block tridiagonal elimination technique. Lastly, the resulting system is coded via MATLAB software and we compare results with relevant literature. An excellent agreement was found between the present results and the literature. The comparison is presented in Table 2. Table 3 displays the thermos-physical properties of utilized nanoparticles and base fluid.

**Table 2:** Comparison of present results with published literature ( $M = 0, \beta \rightarrow \infty, \alpha = 0,$  and  $hs = 1$ ), with different values of Pr

Pr	Salleh et al. [38]		Present	
	$\theta(0)$	$\theta'(0)$	$\theta(0)$	$\theta'(0)$
3	6.02577	7.02577	6.02581	7.02579
5	1.76594	2.76594	1.76598	2.76597
7	1.13511	2.13511	1.13532	2.13515
10	0.76531	1.76531	0.76536	1.76537
100	0.16115	1.16115	0.16121	1.16122

**Table 3:** Thermo-physical characteristics of the original liquid and considered nanomaterials [44]

Thermo-physical feature	Ag	Cu	Al <sub>2</sub> O <sub>3</sub>	Water
$C_p$ (J/kg K)	235	385	765	4179
$\rho$ (kg/m <sup>3</sup> )	10500	8933	3970	997.1
$K$ (W/m K)	429	401	40	0.613
$\sigma$ (S/m)	$6.3 \times 10^7$	$5.96 \times 10^7$	$3.5 \times 10^7$	$5.5 \times 10^{-6}$
Pr	–	–	–	6.2

Briefly, the mathematical formulations and the Keller box method numerical method can be summarized as a flow chart of research methodology, Fig. 2 as follows [45]:

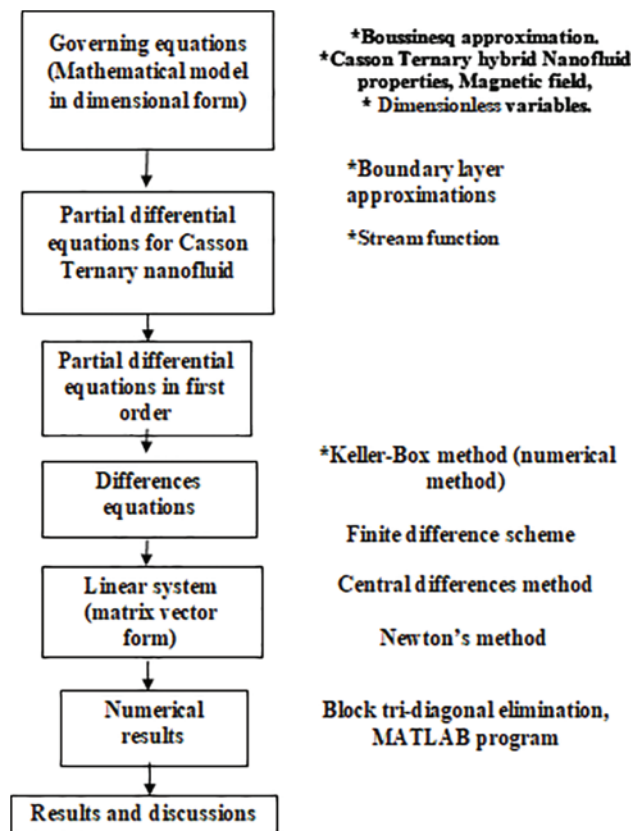
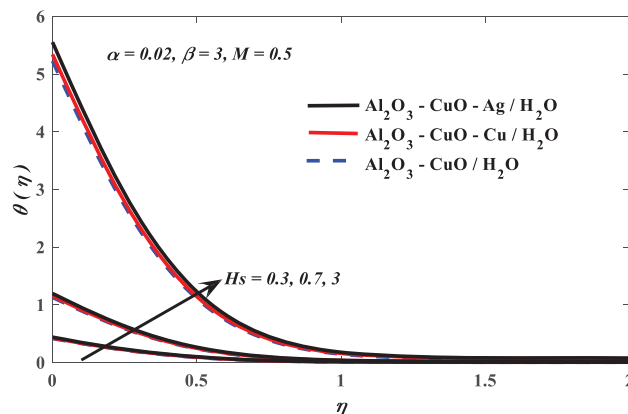


Figure 2: Flow chart of research methodology

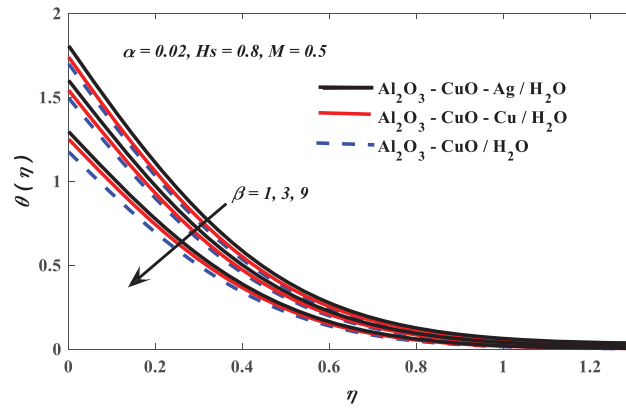
#### 4 Results and Discussions

Graphical impersonations of the influences of different important parameters on physical collections, regarding heat transfer, are offered in Figs. 3–9. The conduct of Casson ternary hybrid nanofluid outcoming from the effect of these parameters is also debated and analyzed. The impacts of Casson  $\beta$ , nanoparticle volume fraction  $\alpha$ , which is defined as  $\alpha = \chi_1 = \chi_2 = \chi_3$ , conjugate  $Hs$ , and magnetic  $M$  parameters have been taken into this investigation, in which their domains are  $0.3 \leq Hs \leq 3$ ,  $1 \leq \beta \leq 9$ ,  $0.01 \leq \alpha \leq 0.03$ , and  $0.1 \leq M \leq 4$ . The conjugate parameter impresses on temperature profiles of ternary hybrid nanofluids while stabilizing the values of the nanoparticle volume fraction  $\alpha = 0.02$ , Casson parameter  $\beta = 3$ , and  $M = 0.5$  as constant values, as clarified in Fig. 3. As the conjugate values are raised, the ternary hybrid nanofluid temperature is boosted. This can be attributed to the improved thermal contact and enhanced thermal conductivity between the stretching sheet and the ternary hybrid nanofluid. As a result, the heat generated by the Newtonian heating source is more efficiently transferred to the ternary hybrid nanofluid, leading to an increase in its temperature. Besides that, when the impacts of the conjugate parameter were employed on the velocity profiles, it was found that the velocity profiles did not respond. This is due to Eq. (5), where the parameter affects only the temperature terms, and does not affect the velocity profiles. Therefore, their figures were not included. From Fig. 4, it can be noted that the temperature profiles of the ternary hybrid nanofluid are getting down by the increase in the Casson parameter values. The repression in the nanofluid

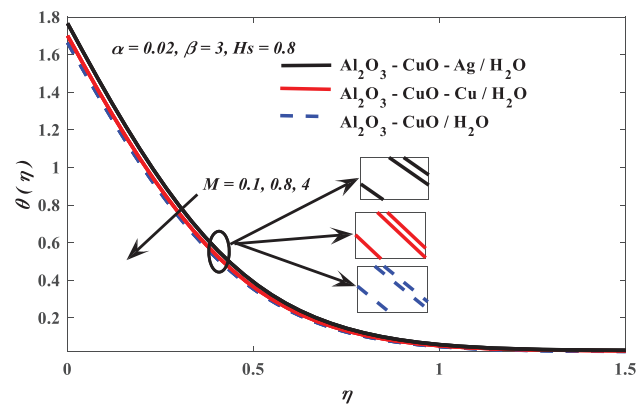
temperature occurs because the increment in the Casson parameter progresses the viscosity of the nanofluids, and this performs a rise in the nanofluid's opposition to temperature, that way decreasing the temperature profiles. Fig. 5 explains the magnitude to which magnetic parameter condensation impacts the temperature profiles of nanofluids while fixing the values of other parameters. It is noted that temperature profile amounts turn down as raise the values of the magnetic parameters. Decreasing occurs due to the limitation in the flow of the fluid resulting from an increment of the magnetic field intensity. Also, Fig. 6 describes the changes in temperature profiles for different values of the nanoparticle volume fraction parameter at  $H_s = 0.8$ ,  $\beta = 3$ , and  $M = 0.5$ . Clearly, growing the values of this parameter supports the increase in the temperature profile. As the volume fraction of nanoparticle is raised, the number of nanoparticles per unit volume increases. This leads to an increase in the total surface area available for heat transfer within the original fluid. The tri-hybrid nanoparticles have high thermal conductivity, and as they are embedded in the original fluid, they enhance the overall thermal conductivity of the original fluid. This increased thermal conductivity facilitates a more efficient transfer of heat from the stretching sheet to the tri-hybrid nanofluid, raising its temperature. Now, we direct our alertness to the investigation of graphical outcomes that supply the velocity profile variations. As both  $\beta$  and  $\alpha$  increase, the fluid velocity increases, and the curves become more betterment pointing to the raise of velocity profiles, as shown in Figs. 7 and 8. In the opposite behavior, Fig. 9, the velocity profiles decrease when the magnetic parameter increases. A magnetic field's presence can induce magneto-hydrodynamic (MHD) effects on the fluid flow. The Lorentz force is one of the primary impacts of MHD. The Lorentz force acts perpendicular to both the magnetic field direction and the fluid velocity. An augmentation in the magnetic parameter corresponds to a heightened intensity of the applied magnetic field, leading to the Lorentz force becoming dominant and tending to oppose the fluid motion. As visible in all presented figures, the comparisons between the compositions (ternary hybrid nanofluids, hybrid nanofluids) are considered. Interestingly, the ternary hybrid nanofluid, formed of water with  $\text{Al}_2\text{O}_3$ -CuO-(Ag or Cu) had higher temperature and velocity profiles compared with the hybrid nanofluid containing water and  $\text{Al}_2\text{O}_3$ -CuO under the Casson, magnetic, nanoparticle volume fraction parameters effects. Additionally, two ternary hybrid nanofluids are compared as physical comparisons. Where it is found that the  $\text{Al}_2\text{O}_3$ -CuO-Ag/water has higher temperature profiles than  $\text{Al}_2\text{O}_3$ -CuO-Cu/water, but the opposite case happens, the velocity profiles for  $\text{Al}_2\text{O}_3$ -CuO-Cu/water are higher than  $\text{Al}_2\text{O}_3$ -CuO-Ag/water, as shown in the all figures included.



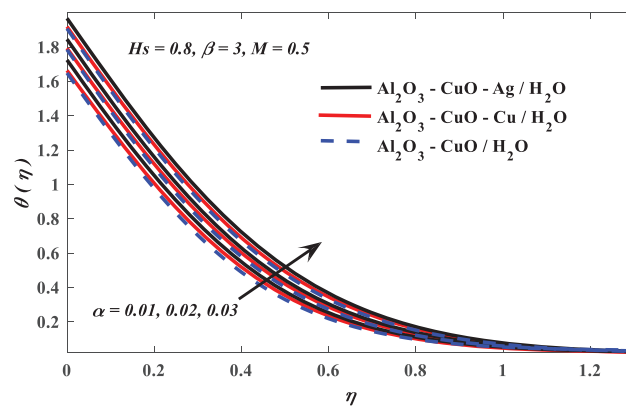
**Figure 3:** Conjugate parameter vs. temperature profiles



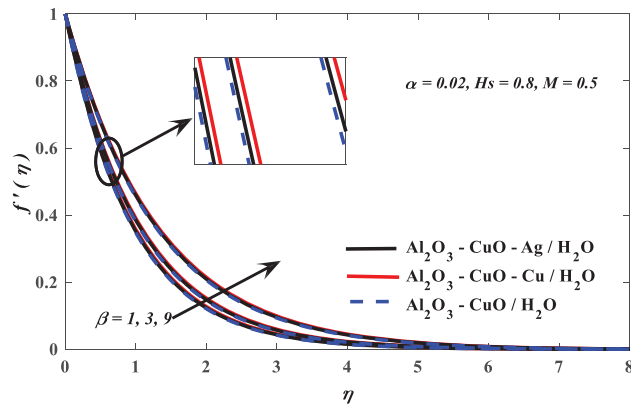
**Figure 4:** Casson parameter vs. temperature profiles



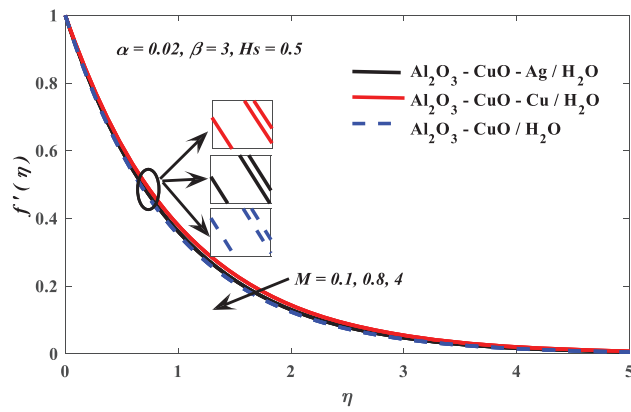
**Figure 5:** Magnetic parameter vs. temperature profiles



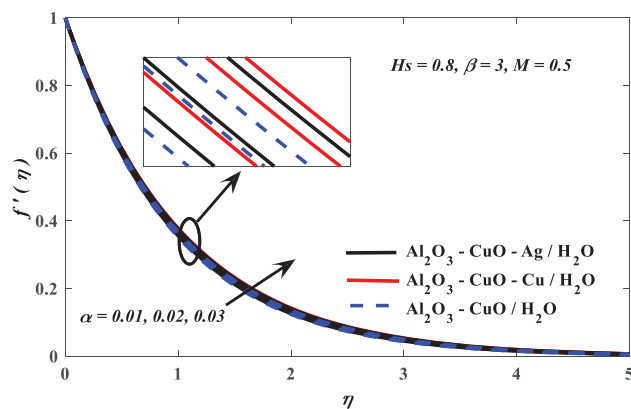
**Figure 6:** Nanoparticle volume fraction parameter vs. temperature profiles



**Figure 7:** Casson fraction parameter vs. velocity profiles



**Figure 8:** Magnetic parameter vs. velocity profiles



**Figure 9:** Nanoparticle volume fraction parameter vs. temperature profiles

Table 4 displays the influences of  $H_s$ ,  $\alpha$ ,  $M$ , and  $\beta$  on the local skin friction coefficient and local Nusselt number, respectively. It is observed that the  $Nu$  distribution is changed at the variation of all utilized parameters values, in this study. Moreover, the local Nusselt number decreases with

the increasing values of both  $Hs$  and  $\beta$ . On the contrary, when the other parameters,  $\alpha$ ,  $M$ , are increased,  $Nu$  values are decreased. Besides that, it is noticed that the  $C_f$  distribution is unaffected by the conjugate parameter  $Hs$ , because there is no effect of this parameter in equation number 20. On the other hand,  $C_f$  is influenced by the other parameters,  $\alpha$ ,  $M$ , and  $\beta$ . We get that the growth in the parameters ( $\alpha$ ,  $M$ , and  $\beta$ ), leads to the rate of local skin friction being reduced. Likewise, it is observed that the ternary hybrid nanofluids have the highest  $Nu$  values than the hybrid nanofluids, with different values of employed parameters,  $Hs$ ,  $\alpha$ ,  $M$ , and  $\beta$ . But it is noticed that the ternary hybrid nanofluids have the lowest  $C_f$  values than the hybrid nanofluids, at the move of these parameters. Additionally,  $Al_2O_3$ -CuO-Ag/water has a lower  $C_f$  and higher  $Nu$  than  $Al_2O_3$ -CuO-Cu/water.

**Table 4** Outcomes of  $Nu$  and  $C_f$  with several values  $Hs$ ,  $\alpha$ ,  $M$ , and  $\beta$

$Hs$	$Al_2O_3$ -CuO-Ag/Water			$Al_2O_3$ -CuO-Cu/Water			$Al_2O_3$ -CuO Water		
	$\alpha$	$M$	$\beta$	$C_f$	$Nu$	$C_f$	$Nu$	$C_f$	$Nu$
0.3	0.02	0.5	3	-4.5209	3.8860	-4.4612	3.9292	-4.2348	3.5583
0.7				-4.5209	2.1665	-4.4612	2.1832	-4.2348	1.9990
3				-4.5209	1.3864	-4.4612	1.3901	-4.2348	1.2981
1	0.01	0.5	3	-3.9981	1.6823	-3.9662	1.6882	-3.8534	1.6212
	0.02			-4.5209	1.8462	-4.4612	1.8576	-4.2348	1.7107
	0.03			-5.0422	2.0349	-4.9570	2.0520	-4.6132	1.8102
1	0.02	0.1	3	-4.5204	1.8462	-4.4607	1.8576	-4.2343	1.7107
		0.8		-4.5213	1.8463	-4.4616	1.8577	-4.2352	1.7108
		4		-4.5255	1.8464	-4.4657	1.8578	-4.2390	1.7109
1	0.02	0.5	1	-1.8461	2.0377	-1.8218	2.0489	-1.7293	1.8789
			3	-4.5209	1.8462	-4.4612	1.8576	-4.2348	1.7107
			9	-12.3806	1.7641	-12.2169	1.7750	-11.5971	1.6384

## 5 Conclusions

This investigation aimed to supply the research gap, which includes the impacts of ternary hybrid nanofluids with the magnetic fields, and subject to Newtonian heating boundary conditions, in the presence of natural convection. The current study developed a mathematical model (governing equations) and used mathematical techniques to convert these equations to partial differential equations (PDEs). Moreover, these PDEs have been successfully solved by the Keller box method, programmed the resulting linear system, and discussed the resulting numerical outcomes. Consequently, we acquired new results which were compared with previous literature. So, it can contribute this consideration to establishing coming studies. Relying on that, this study has outlined the following conclusions:

1. Temperature profiles drop with increased values of the Casson and magnetic parameters, whereas it is boosted when the increment of values of nanoparticle volume fraction and conjugate parameters.
2. Velocity profile escalated with the Casson and nanoparticle volume fraction parameters, but it diminishes due to the growing values of the magnetic parameter.
3. The decrement of the local skin friction coefficient was evaluated with increased Casson, magnetic, and nanoparticle volume parameters, but did not affect with conjugate parameter.

4. The decrement of the local skin friction coefficient was evaluated with increased Casson, magnetic, and nanoparticle volume parameters, but did not affect with conjugate parameter. On the hand, the values of the Nusselt number increased by increasing the nanoparticle volume fraction and magnetic parameters but decreased due to an increase in conjugate and Casson parameters.
5. Local Nusselt number creates better values for ternary hybrid nanofluids compared with the hybrid nanofluid. Likewise, temperature profiles for  $\text{Al}_2\text{O}_3\text{-CuO-Ag/Water}$  (ternary hybrid nanofluids) are higher values than  $\text{Al}_2\text{O}_3\text{-CuO-Cu/Water}$ .
6. Based on this study, we can take into consideration some future studies. The current problem can be developed in future research using more physical influences, such as thermal radiation, chemical reactions, viscous dissipation, and joule heating impacts, as well as it can also develop to comprise ternary hybrid nanofluids with viscous dissipation and Joule heating impacts. On the other hand, it can utilize other boundary conditions; such as convective boundary conditions, constant heat flux, etc.

**Acknowledgement:** The authors happily thank the Center for Graduate Studies Management, Al-Hussein Bin Talal University, Ma'an, Jordan, for the financial support through Voting No. (93/2023) for this research.

**Funding Statement:** This research was funded by the Center for Graduate Studies Management, Al-Hussein Bin Talal University through Voting No. (93/2023).

**Author Contributions:** The authors confirm their contribution to the paper as follows: construct the mathematical formulations, discussions, investigation: Mohammed Z. Swalmeh, Wejdan Mesa'adeen; write literature, review, editing: Firas A. Alwawi; methodology, software: A. A. Altawallbeh; resourcing, writing original draft: Feras M. Al Faqih, Ahmad M. Awajan. All authors reviewed the results and approved the final version of the manuscript.

**Availability of Data and Materials:** Not applicable.

**Conflicts of Interest:** The authors declare that they have no conflicts of interest to report regarding the present study.

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