

Stress distribution due to inclined line loads in fibrous polymer composites

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Summary

Uni-Directional Fiber-Reinforced Plastic (UD-FRP) laminates have been modeled as a quasi-homogeneous monoclinic half-space subjected to an inclined load at the surface. Complete closed form stress-fields derived previously [1] based on Lekhtinskii's formulation [2] in conjunction with the principle of analytical continuation [3] were used to study the stress-response in relation to a few parameters identified, namely, fiber-orientation, load inclination angle and spatial coordinates with respect to line load application position in the half-space domain.

Introduction

Due to the extensive use of fibrous composite as load-bearing components in recent times, there has been a renewed interest in the anisotropic half-space problem as an idealized descriptive model for analyzing UD-FRP laminates [4, 5]. Anisotropic half-space problem with surface load acting is a classical elasticity problem that has received widespread treatment by researchers in the last several decades [6]. Complete closed-form stress-fields for a line-load of fixed width acting symmetrical about the origin at the surface of an anisotropic (monoclinic) elastic half-space have been derived in [1]. The present work deals with understanding the stress-fields by conducting a parametrical analysis vis-à-vis fiber-orientation, load inclination angle and spatial coordinates with respect to the line-load application position for a given width of line load in the half-space domain. The stress-fields thus generated are believed to throw light on the edge-trimming process of UD-FRP laminates.

Inclined load problem

UD-FRP composite laminate is modeled as a continuous quasi-homogeneous linear elastic monoclinic half-space with symmetry existing about z , $X_3 = 0$ axis whose stiffness properties have been calculated from standard tensor transformation relations. A Cartesian coordinate system is defined such that the angle between problem X -axis and the material X_1 -axis in the anti-clockwise direction represents the fiber angle of the UD-FRP as shown in Figure 1. Based on the need to simplify the mathematical analysis, the following three assumptions have been made to derive the stress field equations [1]:- a) Generalized plane strain deformation; b) UD-FRP laminate is a continuous quasi-homogeneous linear elastic monoclinic half space; and c) Quasi-static loading scenario is considered.

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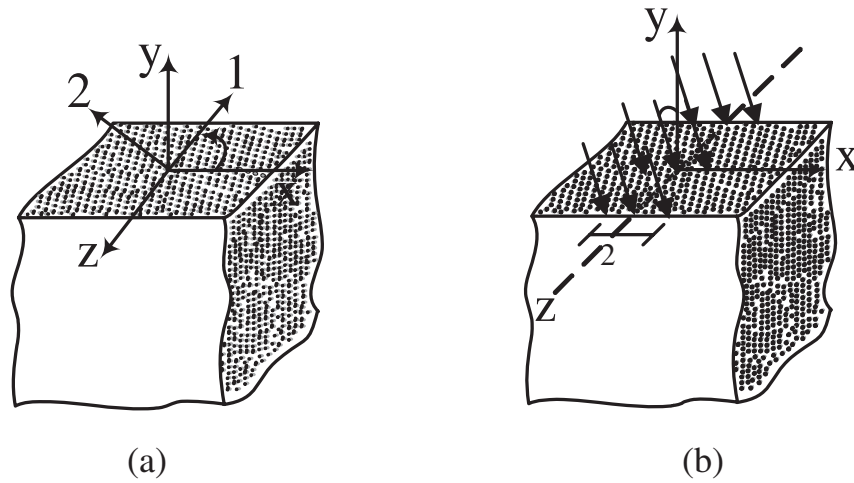


Figure 1: UD-FRP composite representation as a quasi-homogeneous continuous anisotropic elastic half-space; (a) Cartesian coordinate definition; and (b) line load with uniform magnitude along z – axis.

Stress-fields derivation

From Cauchy-Reinmann relations for analytic functions, for a function $f(z = x + iy)$ analytic in the lower half-plane $y \leq 0$, continuous up to the boundary with $f(\infty) = 0$ with ξ , a point on the boundary (abscissa):

$$\frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{f(\xi)}{\xi - z} d\xi = -f(z), \quad \frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{\overline{f(\xi)}}{\xi - z} d\xi = 0 \quad (1)$$

where \bar{f} denotes the complex conjugate of f . This property has been used in the development of the complex potentials which are critical in the determination of stress fields.

For a half-space problem with normal and tangential components of the external tractions on the boundary $y = 0$ are denoted by $N(\xi)$ and $T(\xi)$ respectively. The boundary conditions when $y = 0$ take the form $\sigma_{yy} = N(\xi)$, $\tau_{xy} = T(\xi)$, $\tau_{yz} = 0$. From Muskhelishvili's approach of analytic continuation (extension of Equation (1)), the complex stress functions for a monoclinic material are given by Lekhnit-

skii [2]: -

$$\begin{aligned}\Phi'_1(z_1) &= \frac{1}{2\pi i(\mu_1 - \mu_2)} \int_{-\infty}^{\infty} \frac{N(\xi)\mu_2 + T(\xi)}{\xi - z_1} d\xi \\ \Phi'_2(z_2) &= -\frac{1}{2\pi i(\mu_1 - \mu_2)} \int_{-\infty}^{\infty} \frac{N(\xi)\mu_1 + T(\xi)}{\xi - z_2} d\xi \\ \Phi'_3(z_3) &= 0; \quad z_1 = x + \mu_1 y; \quad z_2 = x + \mu_2 y\end{aligned}\quad (2)$$

where μ_1, μ_2 are the eigenvalues, derived from the stiffness matrix of the half-space [1]. The corresponding stresses are thus known and computed as: -

$$\begin{aligned}\sigma_{xx} &= 2Re [\mu_1^2 \Phi'_1(z_1) + \mu_2^2 \Phi'_2(z_2)] \\ \sigma_{yy} &= 2Re [\Phi'_1(z_1) + \Phi'_2(z_2)] \\ \sigma_{zz} &= -\frac{1}{a_{33}} (a_{13}\sigma_{xx} + a_{23}\sigma_{yy} + a_{63}\tau_{xy}) \\ \tau_{xy} &= -2Re [\mu_1 \Phi'_1(z_1) + \mu_2 \Phi'_2(z_2)] \\ \tau_{xz} &= \tau_{yz} = 0\end{aligned}\quad (3)$$

For a line load of width 2ε (x -direction) of constant magnitude N per unit length symmetric about the z -axis acting at an inclination of α with respect to the y -axis in counter-clockwise direction (Figure 1 (b)), the normal component $N(\xi) = \frac{-N \cos \alpha}{2\varepsilon} (u(\xi + \varepsilon) - u(\xi - \varepsilon))$, and the transverse component $T(\xi) = \frac{-N \sin \alpha}{2\varepsilon} (u(\xi + \varepsilon) - u(\xi - \varepsilon))$, where $u(\xi - \varepsilon)$ is the Heaviside function indicating the location of the line load compressive stresses are taken to be negative. Performing the integrations of (2), the complex stress functions are represented by

$$\begin{aligned}\Phi'_1(z_1) &= \frac{N(\mu_2 \cos \alpha + \sin \alpha)}{4\varepsilon\pi i(\mu_1 - \mu_2)} Ln \left(\frac{z_1 - \varepsilon}{z_1 + \varepsilon} \right), \\ \Phi'_2(z_2) &= \frac{-N(\mu_1 \cos \alpha + \sin \alpha)}{4\varepsilon\pi i(\mu_1 - \mu_2)} Ln \left(\frac{z_2 - \varepsilon}{z_2 + \varepsilon} \right) \\ \Phi'_3(z_3) &= 0\end{aligned}\quad (4)$$

which reduce to Flamant's solution for an isotropic half space as demonstrated in [1]. The complete closed-form stress-field expressions are thus obtained by substituting Equation (4) in (3). It should be noted that there exists a mathematical singularity at $(\pm\varepsilon, 0)$ due to the assumed discontinuous geometry of the line load. This is circumvented by considering stress-fields sufficiently away from the singularity zone.

Stress-Analysis Example

UD-FRP Graphite/Epoxy (Gr/Ep) laminate with $E_{11} = 160.65$ GPa, $E_{22} = 9.62$ GPa, $E_{33} = 9.62$ GPa, $G_{12} = 6.32$ GPa, $G_{13} = 6.32$ GPa, $G_{23} = 3.58$ GPa, $\nu_{12} = 0.2965$, $\nu_{13} = 0.2965$ and $\nu_{23} = 0.342$ has been used in the present investigation. Based on the stress fields derived, a few variables have been identified to perform a parametric analysis, namely (1) Fiber-orientation θ of the UD-FRP laminate; (2) magnitude of load N acting on the UD-FRP laminate; (3) the angle α that the load makes with the vertical axis; and (4) magnitude of width 2ε of the line load. Fiber orientation θ has been identified as the most significant parameter. To this end, the variation of stress-fields with respect to fiber-orientation has been discussed in detail. For the sake of conceptual understanding, the magnitude of N has been set to unity to facilitate comparisons among the various cases. Additionally, the stresses are evaluated at locations in $x - y$ space representative of non-integer multiples of half-width of load profile ε . C.W. Liu's [7] erstwhile work of 1950 provides the basis for the present analysis.

The stress distributions σ_{xx} , σ_{yy} and τ_{xy} at various depths due to a constant line load profile of width $20 \mu\text{m}$ have been plotted for 0° , 45° , 90° and 135° fiber-orientation in Figures 2 through 6 respectively. $\alpha = 45^\circ$ loading scenario as a combination effect of 0° and 90° cases has been looked at. The x - and y - coordinates are normalized with respect to the half-width of line load profile ε . Strictly speaking, as in [7], the magnitude of α would depend on the friction coefficient and the specific problem at hand. However, in this investigation the α 's are chosen arbitrarily in order to characterize the behavior of stress fields, and only one case has been enumerated for brevity.

σ_{yy} distribution evolution for $\alpha = 0^\circ$ has been enumerated in figure 2. The stress fields are found to be symmetric about $x/\varepsilon = 0$ with the maximum stress magnitude occurring at $x/\varepsilon = 0$ for 0° and 90° fiber orientations. Also σ_{yy} is found to decrease more rapidly with increased depth (y/ε) for $\theta = 0^\circ$ s compared to $\theta = 90^\circ$. This can be attributed to the fact that for $\theta = 0^\circ$, the fibers are aligned along the direction of the load application therefore load transfer is not as effective as in $\theta = 90^\circ$. The "fiber" phase is the load carrying member and the "matrix" phase is the load-transferring member of any composite system. This rationale explains some of the salient stress-fields characteristics enumerated in this paper. If the load acts along the direction of the fibers, the stress-dispersion would be minimal. Conversely, if the load is acting perpendicular to the fibers (or along the matrix so to speak) the stress-dispersion would be maximal. Here stress-dispersion indicates how fast the stress magnitudes decrease or decay. For example, σ_{yy} distribution for $\theta = 0^\circ$ exhibits higher stress-dispersion as compared with σ_{yy} distribution for $\theta = 90^\circ$. Another thing to note here is the fact that σ_{yy} has been specified at the

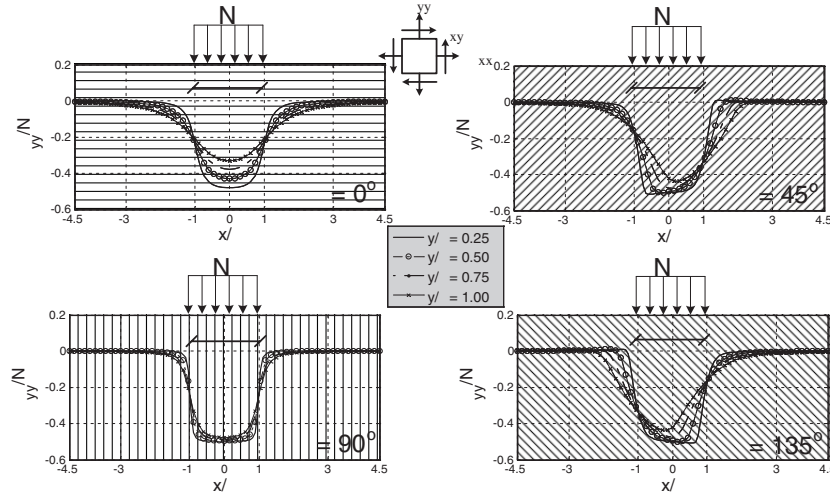


Figure 2: Variation of σ_{yy} due to line load distributions on a width 2ϵ of $20\ \mu\text{m}$ with respect to x/ϵ for Normal Load Case ($\alpha = 0^\circ$).

boundary $y = 0$. Therefore, when we look at the σ_{yy} distribution across the depth, the stress-dispersion becomes an important parameter. For $\theta = 45^\circ$ and 135° , the maximum occurs approximately along the direction of the fiber orientation and the stress-fields are no longer symmetric, however, they seem to be mirror images. Also interesting to note is that the stress fields for $\theta = 45^\circ$ and 135° are indicative of the respective fiber orientations. This observation highlights the preservation of inherent stiffness characteristics along the respective fiber direction of an UD-FRP laminate despite the quasi-homogeneous assumption postulated during problem formulation thereby revalidating the model. Similar inferences have been drawn for σ_{yy} distribution for $\alpha = 90^\circ$ and have therefore not been included here.

Figure 3 depicts σ_{yy} distribution for $\alpha = 45^\circ$ or the combination load case. Being in the domain of linear elasticity, we would expect the stresses for the combination load case to be a sum of the $\alpha = 0^\circ$ and 90° cases. Figure 4 demonstrates such an addition of stress-fields for $\theta = 0^\circ$ laminate under the application of horizontal and normal load cases to replicate the results for the combination load case. Such an addition of stress-fields would be valid for all fiber-orientations and can be demonstrated as for $\theta = 0^\circ$ laminate. Due to such a superposition, the stress-fields do not exhibit the symmetry as observed for normal and horizontal cases. We see that while both compressive and tensile stresses exist for $90^\circ/135^\circ$ laminates, only compressive stresses exist for $0^\circ/45^\circ$ laminates.

σ_{xx} distribution for combined load case has been shown in Figure 5. The additive effects demonstrated for σ_{yy} distribution in Figure 4 hold true here as well.

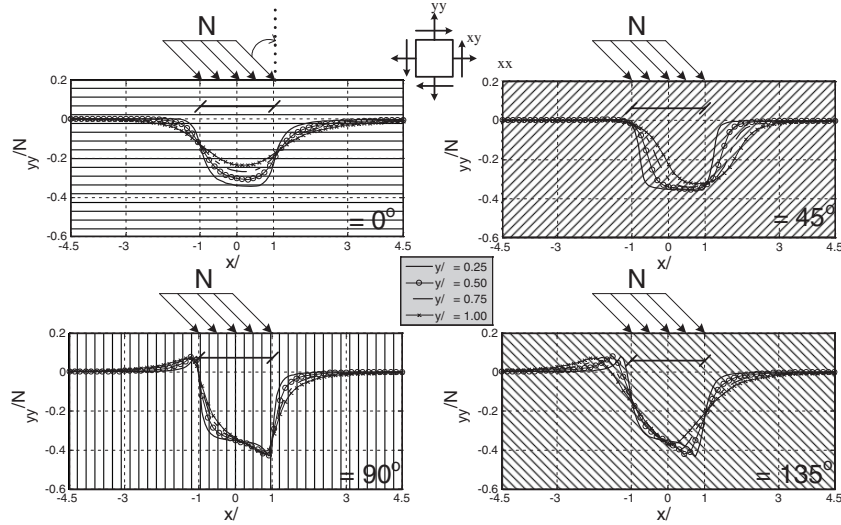


Figure 3: Variation of σ_{yy} due to line load distributions on a width 2ϵ of $20 \mu\text{m}$ with respect to x/ϵ for $\alpha = 45^\circ$.

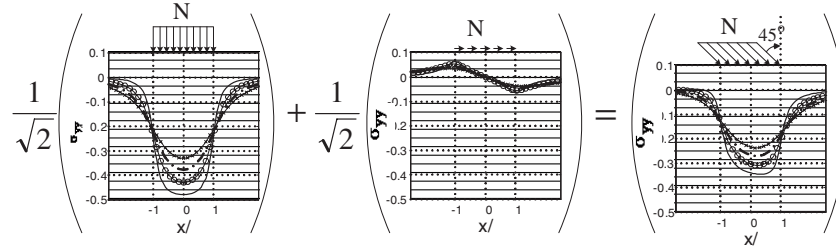


Figure 4: Stresses acting in fiber-orientation $\theta = 0^\circ$ due to combination load have been demonstrated as a sum of stresses due to normal and tangential load cases since the stress-field generation lies in the purview of linear elasticity.

Both compressive and tensile stresses are in play although the magnitude of tensile stress is very small for 45° laminate. The $45^\circ/135^\circ$ laminate stress fields cease to be mirror images due to the additive stress effects.

Finally, Figure 6 depicts the evolution of τ_{xy} distribution for the combined load case. Most FRPs have a lower shear strength values as compared with tensile or compressive strengths in parallel/perpendicular fiber directions. Therefore, the magnitude of τ_{xy} is of great importance in ascertaining the failure behavior of a given FRP laminate which will be dealt with in subsequent works.

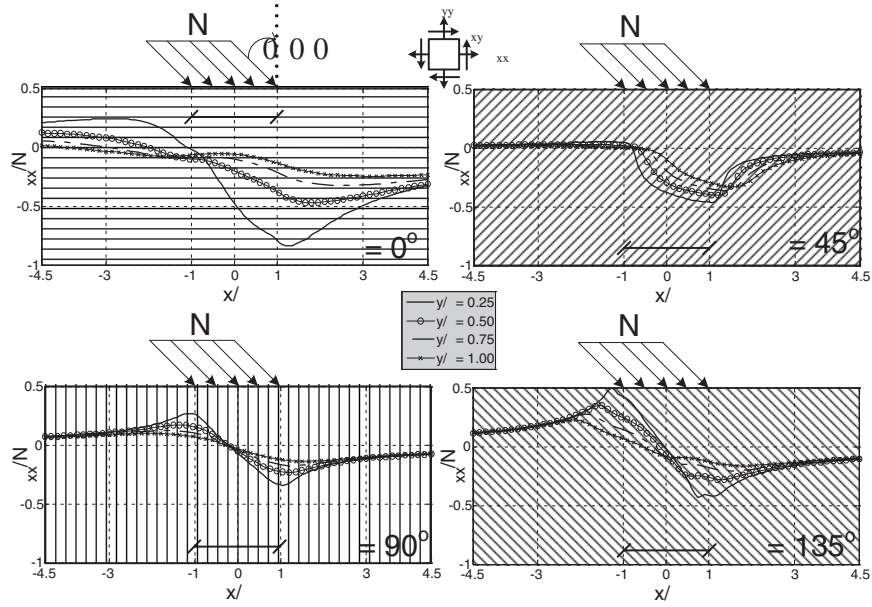


Figure 5: Variation of σ_{xx} due to line load distributions on a width 2ϵ of $20 \mu\text{m}$ with respect to x/ϵ for Combination Case ($\alpha = 45^\circ$).

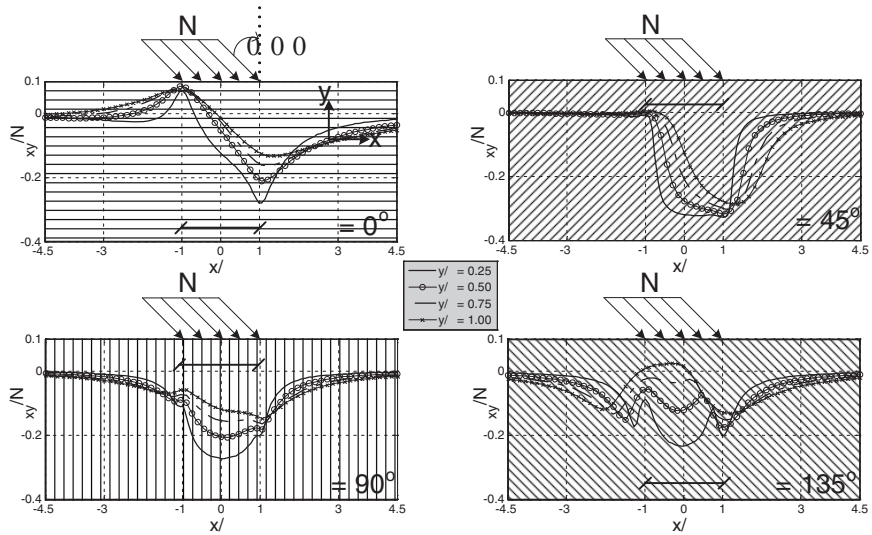


Figure 6: Variation of τ_{xy} due to line load distributions on a width 2ϵ of $20 \mu\text{m}$ with respect to x/ϵ for Combination Load Case ($\alpha = 45^\circ$).

Conclusions and Future work

The basic idea behind this investigation has been to observe and compare the corresponding stress fields in different fiber orientations due to a constant applied line load. The variation in the stress fields among the various fiber-orientations results due to the anisotropic nature of the half-space under consideration. The relation of stress-fields to failure needs to be looked into in future. And finally, the applicability of the present problem to edge-trimming process in UD-FRP laminates needs to be addressed.

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