

MHD FLOW OF CARREAU NANOFUID EXPLORED USING CNT OVER A NONLINEAR STRETCHED SHEET

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ABSTRACT

In the present investigation is to magnetohydrodynamics (MHD) radiative flow of an incompressible steady flow of Carreau nanofluid explored with carbon nanotubes. The boundary layer flow and heat transfer to a Carreau nanofluid model over a non-linear stretching surface is introduced. The Carreau model, adequate for many non-Newtonian fluids is used to characterize the behavior of the fluids having shear thinning properties and fluids with shear thickening properties for numerical values of the power law exponent n . The modeled boundary layer conservation equations are converted to non-linear coupled ordinary differential equations by a suitable transformation. R language with `bvp` solver was adopted to obtain numerical solutions of the resulting equations by using the Runge-Kutta method along with shooting technique. This analysis reveals many important physical aspects of flow and heat transfer. Computations are performed for different values of the stretching parameter (m), the Weissenberg number (We) and the Prandtl number (Pr). The obtained results show that the velocity of shear thinning fluid is depressed by the Weissenberg number while contrasting behavior for the shear thickening fluid is observed. A comparison with previously published data in limiting cases is performed and they are in excellent agreement.

Keywords: Shear thinning and Shear thickening, Radiation, Runge-Kutta method, R language.

1. INTRODUCTION

It is renowned fact that the magnetohydrodynamic (MHD), which is the science of motion of electrically conducting fluids, is one of the thrust areas of modern research. The elementary examples of electrically conducting fluids include plasmas, liquid metals (mercury or liquid sodium) and electrolytes. The basic theme of MHD is that if an electrically conducting fluid is placed in magnetic field then motion of fluid may create a force called electromotive force. The electromotive force has the ability to induce current. Ever since then, this field has a broad spectrum of science and engineering, specifically in geophysics, fusion reactors, dispersion of metals, modern metallurgy and MHD generators etc. Moreover, MHD flows are of immense concern in problems related with physiological fluids. Pavlov (1974) was the pioneer who discussed the influence of magnetic field on MHD flow past a stretching surface. Another eminent contribution was given by Andersson (1992) who examined the MHD flow of a viscous fluid. Nanofluids can be described as a colloidal suspension of nanoparticles (1–100 nm) scattered uniformly in a base fluid. The nanoparticles are made from materials that are chemically stable like metals, oxides ceramics, carbides and non metals such as graphite and carbon nanotubes. Nanotubes are affiliates of the fullerene structural family. Their name arises from their long, resonating edifice with the ramparts formed by one-atom-thick sheets of carbon, called graphene. These sheets are rolled at specific and discrete (“chiral”) angles, and the combination of the rolling angle and radius decides the nanotube properties; for

example, whether the individual nanotube shell is a metal or semiconductor. Nanotubes are categorized as single-walled nanotubes (SWNTs) and multi-walled nanotubes (MWNTs). The commonly used base fluids are ethylene, water, oils and lubricants. Choi (1995) who was working with a group of Argonne National Laboratory USA, introduced the term Nanofluid. Nanofluids enhance the thermophysical properties of the base fluids used. Nanofluids have a wide range of practical applications, they can be used as smart materials to work as heat valve to control the flow of heat. Advanced electronic gadgets experience thermal management challenges from reduction of accessible surface area for heat and high level of heat generation. These challenges can be avoided by either increasing the capacity of heat transfer or by finding an optimum geometry of cooling appliances. Due to their high thermal conductivity and increased heat transfer coefficient, nanofluids can be used for liquid coolant of computer processors. In presence of magnetic field, nanofluids have several engineering, industrial and biomedical applications such as cooling of nuclear reactors by liquid sodium and induction flow meter, magnetogravimetric separations, smart fluids for vibration damping, aerodynamic sensors magnetic drug targeting, nanocryosurgery and nanodrug delivery.

Till date, as a consequence of enormous industrial applications the notion of non-Newtonian fluids have dragged much insight when equated to Newtonian liquids. Moreover, plenty of experimental and notional studies have been executed to scrutinize the mechanism of non-Newtonian

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transport owing to their extensive applications in numerous genetic and developed progressions, materials and motorized engineering. Numerous materials, for instance biomedical flows, splashes, bio-fluids in genetic material and polymers are identified as non-Newtonian fluids. The non-Newtonian Carreau fluid model is the combination of Newtonian and power-law models. This fluid model is capable of describing both shear thinning and shear thickening phenomena. In spite of all such concentrations, innumerable investigators are quiet engaged to scrutinize the analysis of non-Newtonian liquids under diverse prospective. Rheological equations from molecular network theories was presented by Carreau (1972). Hydromagnetic natural convection from an isothermal inclined surface adjacent to a thermally stratified porous medium has been studied by Chamkha (1997). "Unsteady flow and heat transfer on a semi-infinite flat plate with an aligned magnetic field (Takhar et al., 1999)". "Unsteady three-dimensional MHD-boundary-layer flow due to the impulsive motion of a stretching surface (Takhar et al., 2001)". Magneto-hydrodynamic three dimensional flow and heat transfer over a stretching surface in a viscoelastic fluid was explored by Ahmad and Nazar (2011). Boundary layer flow and heat transfer to Carreau fluid over a nonlinear stretching sheet was presented by Khan and Hashim (2015). Numerical simulation of water based magnetite nanoparticles between two parallel disks were simulated by Haq et al. (2016). Impact of nonlinear thermal radiation and gyrotactic microorganisms on the Magneto-Burgers nanofluid were studied by Khan et al. (2017). A microscopic study of MHD fractional inertial flow through Forchheimer medium were shown by Anwar and Rasheed (2017). Carbon nanotubes significance in Darcy-Forchheimer flow were established by Hayat et al. (2018). Lu et al. (2018) investigated the effect of a numerical treatment of MHD radiative flow of Micropolar nanofluid with homogeneous-heterogeneous reactions past a nonlinear stretched surface. "Unsteady squeezing carbon nanotubes based nano-liquid flow with Cattaneo-Christov heat flux and homogeneous-heterogeneous reactions (Lu et al., 2019)". Numerical study of Carreau nanofluid flow past vertical plate with the Cattaneo-Christov heat flux model was shown by Vasu and Ray (2019).

The objective of present investigation is to discuss the boundary layer flow and heat transfer of Carreau nanofluid model over a non-linear stretching surface. We obtained the boundary layer equations for the velocity field of the Carreau nanofluid explored with CNT. This model which, in comparison with the Newtonian case involves three additional dimensionless parameters, allows the characterization of a wide variety of shear-thinning fluids. The numerical method namely Runge-Kutta shooting method is used to solve the non-linear governing problem. The obtained results for both the fluid velocity and its temperature are presented for different values of parameters entering in the problem. The effects of engineering parameters, like the Weissenberg number (We), the power law index n and the Prandtl number (Pr) on the local skin friction and local Nusselt number are presented numerically in tabular form.

2. MATHEMATICAL FORMULATION

Consider the problem of an incompressible two dimensional flow of a generalized non Newtonian fluid, namely Carreau nanofluid, due to a stretching sheet which coincides with the plane $y=0$. It is assumed that the flow being confined to $y > 0$. A uniform magnetic field B_0 is applied parallel to the y -axis. The induced magnetic field is neglected for small magnetic reynold number. Heat and mass transfer characteristics are taken into account in the presence of thermal radiation and thermophoresis effects. The uniform temperature of the surface T_w is larger than ambient fluid temperature T_∞ and it is moving with nonlinear velocity U_w . The species concentration at the surface C_w and ambient concentration C_∞ are constants. The following governing equations are introduced (Khan and Hashim, 2015)

$$\frac{\partial p}{\partial x} + \frac{\partial q}{\partial y} = 0 \quad (1)$$

$$\begin{aligned} p \frac{\partial p}{\partial x} + q \frac{\partial p}{\partial y} = & \nu_{nf} \frac{\partial^2 p}{\partial x^2} \left(1 + \tau^2 \left(\frac{\partial p}{\partial y} \right)^2 \right)^{\left(\frac{n-1}{2} \right)} \\ & + \nu_{nf} (n-1) \tau^2 \left(\frac{\partial p}{\partial y} \right)^2 \frac{\partial^2 p}{\partial x^2} \left(1 + \tau^2 \left(\frac{\partial p}{\partial y} \right)^2 \right)^{\left(\frac{n-3}{2} \right)} \\ & - \frac{\sigma_{nf}}{\rho_{nf}} B^2 p + g(\beta_T(T - T_\infty) + \beta_C(C - C_\infty)) \end{aligned} \quad (2)$$

$$\begin{aligned} \left(p \frac{\partial T}{\partial x} + q \frac{\partial T}{\partial y} \right) (\rho c_p)_{nf} = & k_{nf} \frac{\partial^2 T}{\partial y^2} + \mu_{nf} \left(\frac{\partial p}{\partial y} \right)^2 - \frac{\sigma_{nf}}{\rho_{nf}} B^2 p^2 \\ & + \frac{16\sigma^s T_\infty^3}{3k^s} \frac{\partial^2 T}{\partial y^2} \end{aligned} \quad (3)$$

$$p \frac{\partial C}{\partial x} + q \frac{\partial C}{\partial y} = D \frac{\partial^2 C}{\partial y^2} - \frac{\partial}{\partial y} (V_T (C - C_\infty)) \quad (4)$$

where p and q are the velocity components in x and y directions respectively, ν_{nf} kinematic nanofluid viscosity, τ material time constant, n expression power law index, ρ_{nf} density of nanofluid, k_{nf} thermal conductivity of nanofluid, σ_{nf} electrical conductivity of nanofluid, g the gravitational acceleration β_T and β_C thermal expansion coefficients of temperature and concentration, c_p specific heat, σ^s the Stefan-Boltzmann constant, k^s the mean absorption coefficient, D the diffusion coefficient and B nonlinear uniform magnetic field V_T the thermophoretic velocity defined as

$$B = B_0 x^{\left(\frac{m-1}{2} \right)}, \quad V_T = -k_1 \frac{\nu}{T_{ref}} \frac{\partial T}{\partial y} \quad (5)$$

where k_1 is the thermophoretic coefficient, T_{ref} is the reference temperature.

Corresponding boundary conditions are

$$\begin{aligned} p = p_w(x) = bx^m, q = 0, T = T_w, C = C_w \quad \text{at } y = 0 \\ p \rightarrow 0, T \rightarrow T_\infty, C \rightarrow C_\infty \quad \text{as } y \rightarrow \infty \end{aligned} \quad (6)$$

where b and m are positive real numbers relating to the stretching sheet.

Nanofluid parameters are defined as

$$\begin{aligned} \nu_{nf} &= \frac{\mu_{nf}}{\rho_{nf}}, \\ \rho_{nf} &= (1 - \phi)\rho_{bf} + \phi\rho_{CNT}, \\ \mu_{nf} &= \frac{\mu_{bf}}{(1 - \phi)^{2.5}}, \\ (\rho c_p)_{nf} &= (1 - \phi)(\rho c_p)_{bf} + \phi(\rho c_p)_{CNT}, \\ \frac{k_{nf}}{k_{bf}} &= \frac{(k_{CNT} + 2k_{bf}) - 2\phi(k_{bf} - k_{CNT})}{(k_{CNT} + 2k_{bf}) + \phi(k_{bf} - k_{CNT})}, \\ \frac{\sigma_{nf}}{\sigma_{bf}} &= 1 + \frac{3(\sigma_{CNT} - \sigma_f)\phi}{(\sigma_{CNT} + 2\sigma_{bf}) - (\sigma_{CNT} - \sigma_{bf})\phi} \end{aligned} \quad (7)$$

Now introducing the following dimensionless variables for stream function, temperature, concentration and the independent variable as:

$$\begin{aligned} \psi &= \sqrt{\frac{2b\nu_{bf}}{m+1}} x^{\left(\frac{m-1}{2} \right)} f(\eta), \quad \eta = yx^{\left(\frac{m-1}{2} \right)} \sqrt{\frac{b(m+1)}{2\nu_{bf}}}, \\ p &= \frac{\partial \psi}{\partial y}, \quad q = -\frac{\partial \psi}{\partial x} \\ \theta(\eta) &= \frac{(T - T_\infty)}{(T_w - T_\infty)}, \quad h(\eta) = \frac{(C - C_\infty)}{(C_w - C_\infty)} \end{aligned} \quad (8)$$

Using the above variables into Eqs. (1)–(4) and boundary conditions (6), governing equations are reduced as ordinary differential equations as follows

$$\frac{P_1}{P_2}(1 + nWe_x^2(f'')^2)(1 + We_x^2(f'')^2)^{\frac{n-3}{2}} f''' - \frac{P_3}{P_1} l_1 M f' + \gamma l_1 (\theta + Nh) - \beta(f'^2) + f f'' = 0 \quad (9)$$

$$(1 + \frac{4}{3P_5} R) \theta'' + \frac{P_4}{P_5} (Pr \theta' f') + \frac{P_1}{P_5} Pr Ec (f'')^2 - \frac{P_3}{P_5} M l_1 Ec Pr (f')^2 = 0 \quad (10)$$

$$h'' + Sc(f h') - \tau^s Sc(h \theta'' + \theta' h') = 0 \quad (11)$$

and the corresponding boundary conditions

$$f(0) = 0, \quad f'(0) = 1, \quad \theta(0) = 1, \quad h(0) = 1 \quad \text{at} \quad \eta = 0$$

$$f'(\eta) = 0, \quad \theta(\eta) = 0, \quad h(\eta) = 0 \quad \text{as} \quad \eta \rightarrow \infty \quad (12)$$

where prime denotes differentiation with respect to η ,

$$\beta = \frac{2m}{m+1}, \quad l_1 = \frac{2}{m+1}, \quad We_x = \left(\frac{(m+1)\tau^2 x^{3m-1}}{2\nu_f} b^3 \right)^{0.5},$$

$$\gamma = \frac{Gr_x}{Re_x^2}, \quad Gr_x = \frac{\beta_T(T_w - T_\infty)x^3}{\nu_f^2}, \quad Re_x = \frac{u_w x}{\nu_f},$$

$$Pr = \frac{\mu c_p}{k}, \quad Ec = \frac{u_w^2}{(c_p)_f(T_w - T_\infty)}, \quad Sc = \frac{\nu_f}{D},$$

$$M = \frac{\sigma_f B_0^2}{b \rho_f}, \quad N = \frac{\beta_C(C_w - C_\infty)}{\beta_T(T_w - T_\infty)},$$

$$R = \frac{4\sigma^s T_\infty^3}{k^s k}, \quad \tau^s = -\frac{k_1(T_w - T_\infty)}{T_{ref}} \quad (13)$$

Some nanofluid parameters are

$$P_1 = \frac{1}{(1-\phi)^{2.5}}, \quad P_2 = \left[(1-\phi) + \phi \frac{\rho_{CNT}}{\rho_{bf}} \right],$$

$$P_3 = 1 + \frac{3(\sigma_{CNT} - \sigma_{bf})\phi}{(\sigma_{CNT} + 2\sigma_{bf}) - (\sigma_{CNT} - \sigma_{bf})\phi}$$

$$P_4 = \left[(1-\phi) + \phi \frac{(\rho c_p)_{CNT}}{(\rho c_p)_{bf}} \right], \quad (14)$$

$$P_5 = \frac{(1-\phi) + 2\phi \left(\frac{k_{CNT}}{k_{CNT} - k_{bf}} \right) \ln \left(\frac{k_{CNT} + k_{bf}}{2k_{bf}} \right)}{(1-\phi) + 2\phi \left(\frac{k_{bf}}{k_{CNT} - k_{bf}} \right) \ln \left(\frac{k_{CNT} + k_{bf}}{2k_{bf}} \right)}$$

Engineering parameters likeskin friction coefficient Cf_x , local Nusselt Nu_x and Sherwood Sh_x numbers are defined as

$$Cf_x = \frac{\tau_w}{\rho_f p_w^2}, \quad Nu_x = \frac{xq_w}{k_f(T_w - T_\infty)}, \quad Sh_x = \frac{xj_w}{D(C_w - C_\infty)} \quad (15)$$

in which τ_w , q_w , and j_w denote the surface shear stress, the surface heat flux and the surface mass flux respectively. Here

$$\tau_w = \mu_{nf} \left(\frac{\partial p}{\partial y} \left(1 + \tau^2 \left(\frac{\partial p}{\partial y} \right)^2 \right)^{\left(\frac{n-1}{2} \right)} \right)_{y=0}, \quad (16)$$

$$q_w = -k_{nf} \left(\frac{\partial T}{\partial y} \right)_{y=0}, \quad j_w = -D \left(\frac{\partial C}{\partial y} \right)_{y=0}$$

In dimensionless coordinates one can express that

$$Re_x^{0.5} Cf_x = \left(\frac{m+1}{2} \right)^{0.5} (1 + we_x^2(f'')^2)^{\frac{n-1}{2}} (1-\phi)^{-2.5} f''(0),$$

$$Re_x^{-0.5} Nu_x = - \left(\frac{m+1}{2} \right)^{0.5} P_5 \theta'(0),$$

$$Re_x^{-0.5} Sh_x = - \left(\frac{m+1}{2} \right)^{0.5} h'(0) \quad (17)$$

3. NUMERICAL METHOD FOR SOLUTION

The nonlinear differential Eqs. (9) to (11) with boundary conditions (12) have been solved numerically using Runge–Kutta fourth order method along with shooting technique using bvp solver package in R software. Thus, Eqs. (9) to (11) have been converted into first-order equations as follows:

$$g'_1 = g_2$$

$$g'_2 = g_3$$

$$g'_3 = \frac{\beta(g_2)^2 - g_1 g_3 + \frac{P_3}{P_2} l_1 M g_2 - l_1 \gamma (g_4 + N g_5)}{\frac{P_1}{P_2} (1 + nwe_x^2(g_3)^2) (1 + we_x^2(g_3)^2)^{\frac{n-3}{2}}}$$

$$g'_4 = g_5 \quad (18)$$

$$g'_5 = \frac{-\frac{P_4}{P_5} (Pr g_5 f_1) - \frac{P_1}{P_5} Pr Ec (g_3)^2 + \frac{P_3}{P_5} M l_1 Ec Pr (g_2)^2}{(1 + \frac{4}{3P_5} R)}$$

$$g'_6 = g_7$$

$$g'_7 = -Sc(g_1 g_7) + \tau^s Sc(g_6 g'_5 + g_7 g_5)$$

$$\text{where } g_1 = f, g_2 = f', g_3 = f'', g_4 = \theta, g_5 = \theta', g_6 = h, g_7 = h'. \quad (19)$$

Boundary conditions (12) becomes,

$$g_1(0) = 0, g_2(0) = 1, g_4(0) = 1, g_6(0) = 1$$

$$g_2(\infty) = 0, g_4(\infty) = 0, g_6(\infty) = 0 \quad (20)$$

In view of the above boundary conditions, out of seven only four initial conditions are known and solution procedure is begun by three initial guesses of $g_3(0), g_5(0), g_7(0)$ with the shooting method iteration to assure the iteration convergence required to satisfy asymptotically boundary conditions at infinity. Thus, in the present study the value of $\eta = \eta_{max}$ is taken to be 10 and step-size is taken to be $\Delta \eta = 0.01$ with relative error tolerance 10^{-5} .

4. RESULTS AND DISCUSSION

The behavior of different variables for the velocity, thermal field, nanoparticles concentration, skin friction coefficient and local Nusselt and Sherwood numbers are investigated through plots.

4.1. Accuracy assessment

Accuracy of the present results are compared with results mentioned in Khan and Hashim (2015). In the absence of nanoparticles, Magnetic parameters, Weissenberg number, local buoyancy parameter, Eckert number, Schmidt number, radiation parameter, these comparisons are found to be in excellent agreement as shown in tables 1 and 2.

4.2. Dimensionless velocity

Fig. 1 demonstrates that by enhancing Weissenberg number, velocity profile decreases in case of shear thinning ($n < 1$) fluid while it shows increasing behavior for shear thickening ($n > 1$) fluid. Since Weissenberg number is the ratio of the relaxation time of the fluid and a specific process time, it enhances the thickness of fluid which in-turn causes the reduction of

velocity. Fig. 2 is plotted in order to examine the influence of magnetic field on Carreau nanofluid velocity for shear thickening fluid. It is clear from figures that the Carreau nanofluid velocity is a decreasing function of applied magnetic field. The inciting values of magnetic field parameter reflect higher values of Lorentz force. Since this force is resistive in nature, the Carreau nanofluid particles faced resistance and hence a decline in velocity curves appeared. Fig. 2 also shows that the velocity profile is increasing function of γ due to this boundary layer increases. Investigators observe that all curves move to zero at the threshold value of η is 4.7. For $\eta = 0$ the velocities have the same values since there is no slip condition between the nanofluid and the surface and it is due to the boundary condition at the wall.

4.3. Dimensionless temperature

Fig. 3 portrayed the behavior of radiation parameter(R), Prandtl number (Pr), and magnetic parameter(M). Thermal field and layer thickness are increasing functions of R. As expected, heat is produced due to radiation process in the working fluid, so the temperature field enhances. Temperature for Prandtl number (Pr) is drawn in Fig. 3. Here both temperature and layer thickness are diminished for larger Pr. Physically higher Pr lead to lower thermal diffusivity which results in the reduction of thermal field. Fig. 3 clearly indicates that the temperature and layer thickness are increased for larger magnetic parameter M. Physically higher magnetic parameter increases heat transfer coefficient which enhances the fluid temperature. Fig. 4 delineates temperature variation via Weissenberg number (We). For larger Weissenberg number (we) both temperature and layer thickness is decreased. In fact, less heat is produced through random motion of the fluid particles due to larger Weissenberg number (we) variable and consequently the temperature decreases. Investigators observe that thermal boundary layers move to zero at the threshold value of η is 5.7. For $\eta = 0$ the temperature remains the same due to the boundary condition at the wall.

4.4. Dimensionless concentration

Graphical illustration for concentration against Schmidt number (Sc), thermophoretic parameter(τ^s) are captured in Fig. 5. It is revealed that concentration and related layer thickness are decayed for larger Schmidt number (Sc). Schmidt number (Sc) is the ratio of momentum to mass diffusivities. For enhancement of Schmidt number, the mass diffusivity diminishes which is responsible in reduction of concentration. Boundary layer of concentration decreases with increasing thermophoretic parameter shown in the figure. Outcomes of generative/destructive chemical reaction γ on concentration field are delineated in the figure. There is an enhancement in fluid concentration in view of destructive chemical reaction variable. Investigators observe that boundary layers of concentration move to zero at the threshold value of η is 5.7. For $\eta = 0$ the concentration remains the same due to the boundary condition at the wall.

4.5. Dimensionless Engineering parameters

Behaviors of magnetic parameter(M), Weissenberg number(We) on skin friction coefficient, local Nusselt number, and local Sherwood number are disclosed in Table.3.

4.6. Different Base fluids and Carbon nanotubes

Figs. 6-7 shows that the momentum boundary layer of water based single-walled carbon nanotube is more when compared to kerosene and engine oil at both shear thinning and shear thickening fluid. Figs. 8-9 depict that the thermal boundary layer of water is more when compared to remaining base fluids in both shear thinning and shear thickening cases. Figs. 10-11 represent that concentration boundary layer of water is less when compared to the remaining base fluids in shear thinning but in shear thickening the concentration of boundary layer increases. Momentum boundary layer of water based multi-walled carbon nanotube is more when compared to kerosene and engine oil at both shear thinning and shear thick-

ening fluid as shown in Figs. 12-13. Figs. 14-15 depict that the thermal boundary layer of water is more as compared to remaining base fluid in both shear thinning and shear thickening cases. Figs. 16-17. represents concentration boundary layer of water is less with reference to the remaining base fluids in shear thinning but in shear thickening the concentration of boundary layer increases.

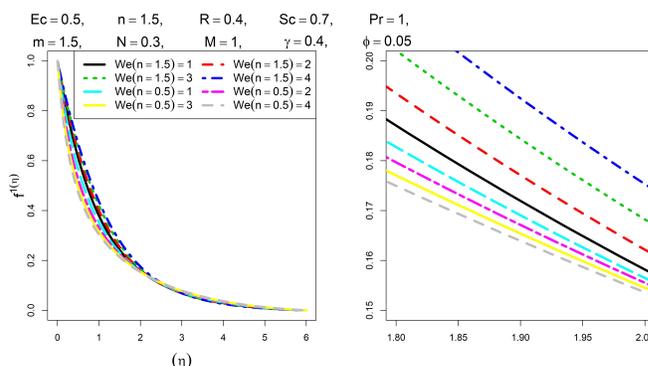


Fig. 1 Velocity profile for Shear thinning and Shear thickening fluid with Weissenberg number (We)

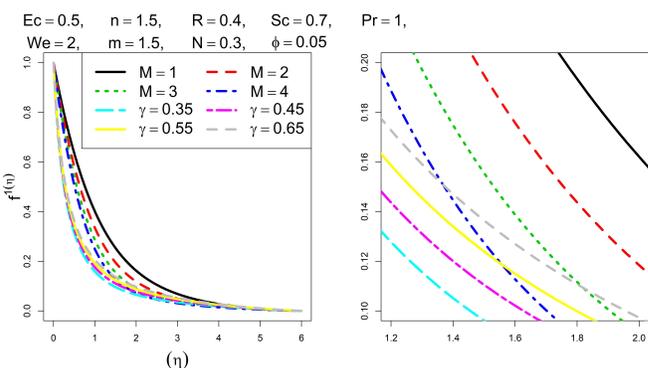


Fig. 2 Velocity profile for Shear thickening fluid with Magnetic parameter(M), and Local buoyancy parameter(γ)

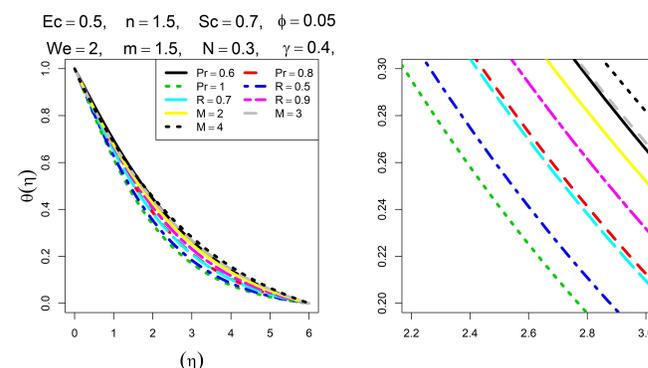


Fig. 3 Temperature profile for Shear thickening fluid with Radiation(R), Prandtl number (Pr), and magnetic parameter(M)

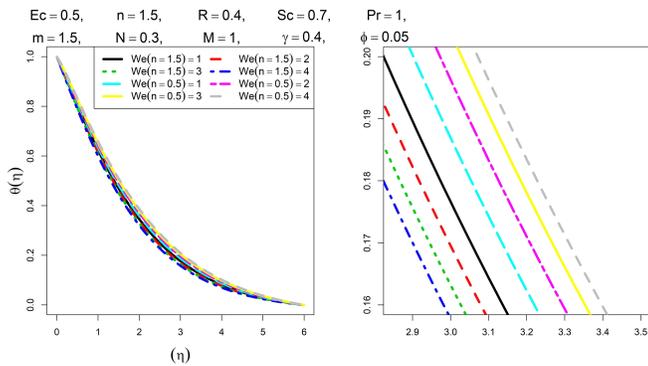


Fig. 4 Temperature profile for Shear thickening and Shear thinning fluid with Weissenberg number (We)

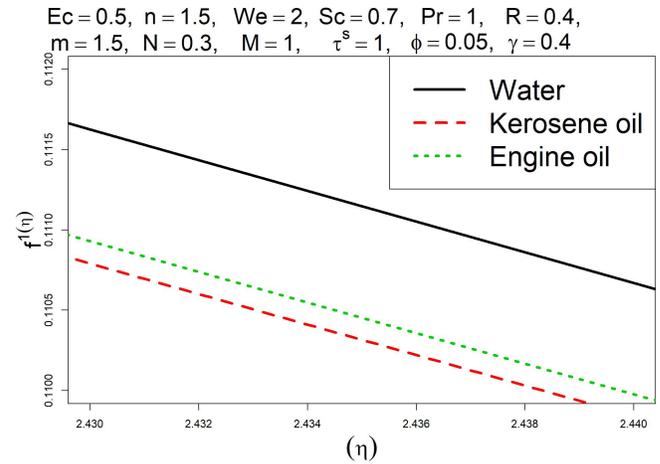


Fig. 7 Velocity profile for Shear thickening fluid with SWCNT

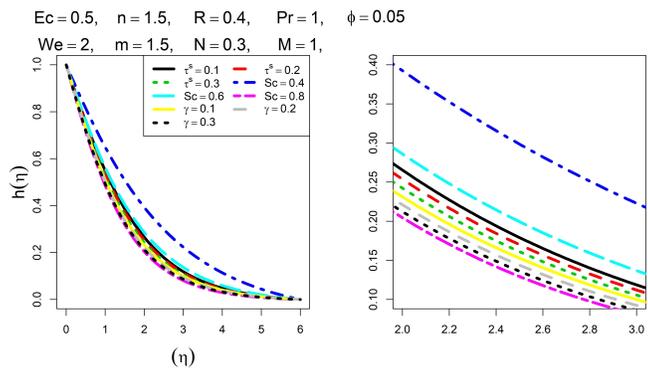


Fig. 5 Concentration profile for thermophoretic parameter (τ^s), Schmidt number (Sc), and Local buoyancy parameter (γ)

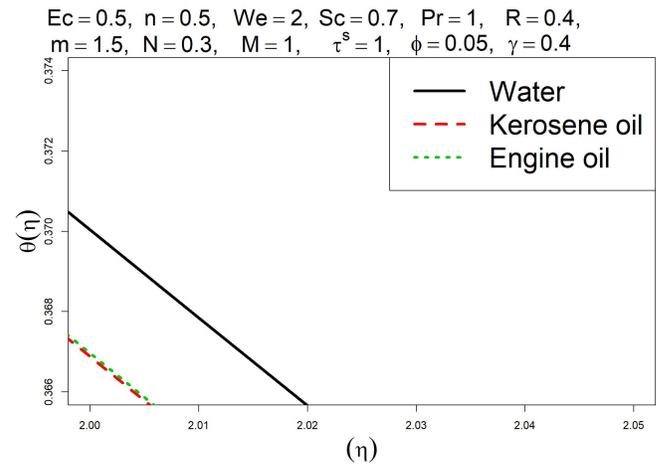


Fig. 8 Temperature profile for Shear thinning fluid with SWCNT

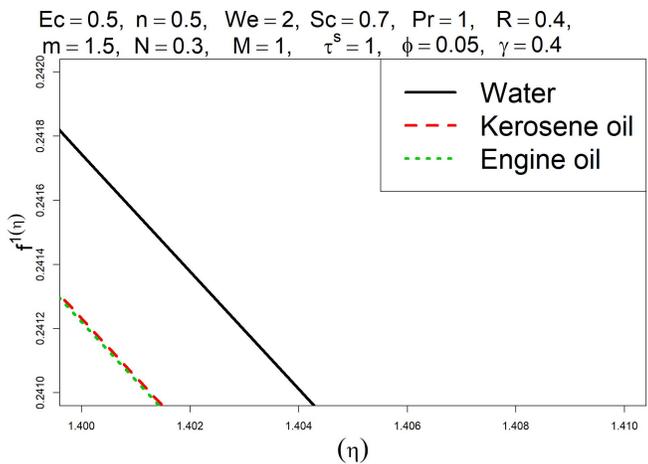


Fig. 6 Velocity profile for Shear thinning fluid with SWCNT

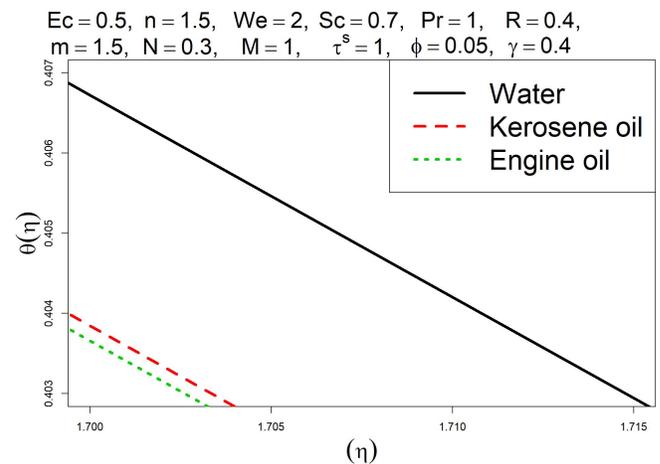


Fig. 9 Temperature profile for Shear thickening fluid with SWCNT

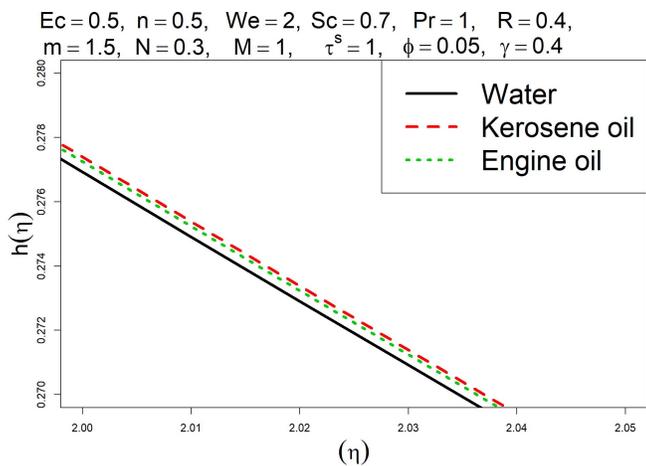


Fig. 10 Concentration profile for Shear thinning fluid with SWCNT

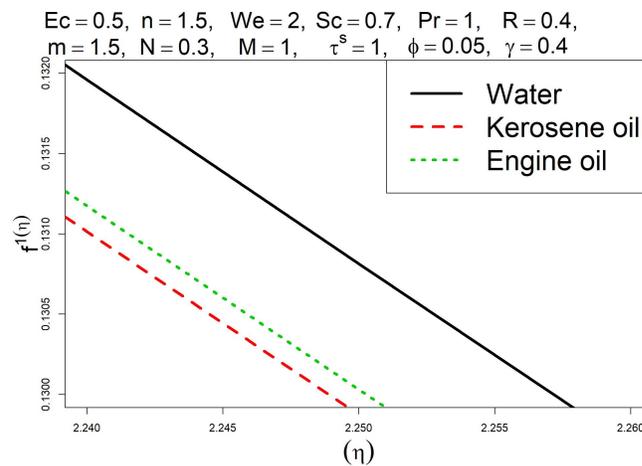


Fig. 13 Velocity profile for Shear thickening fluid with MWCNT

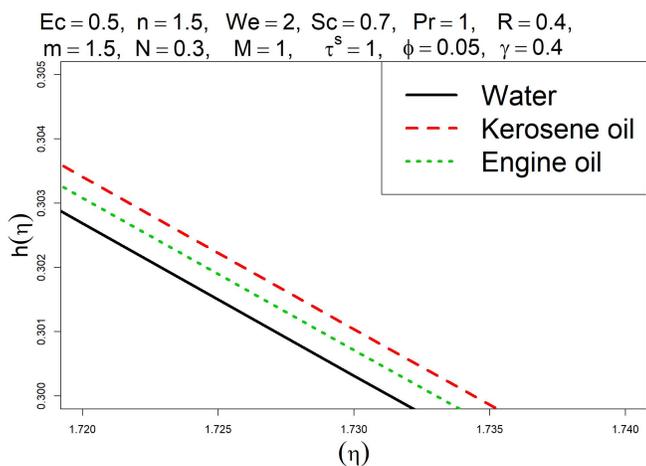


Fig. 11 Concentration profile for Shear thickening fluid with SWCNT

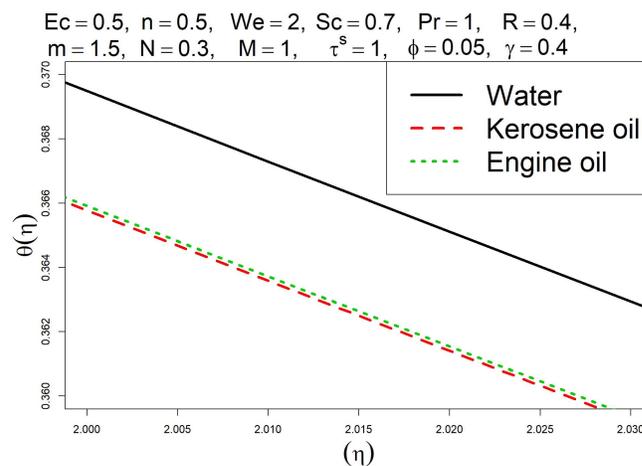


Fig. 14 Temperature profile for Shear thinning fluid with MWCNT

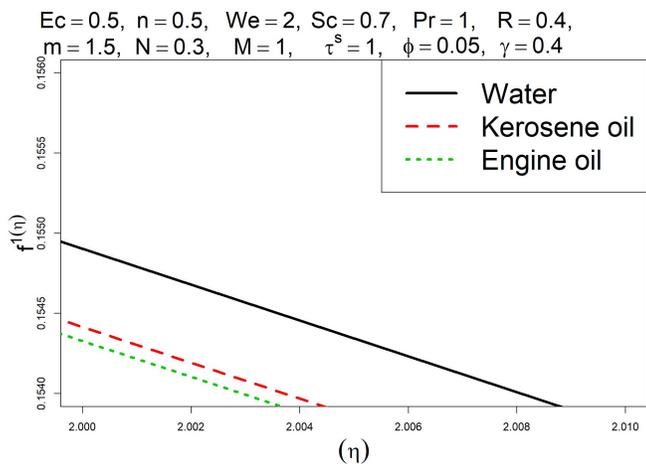


Fig. 12 Velocity profile for Shear thinning fluid with MWCNT

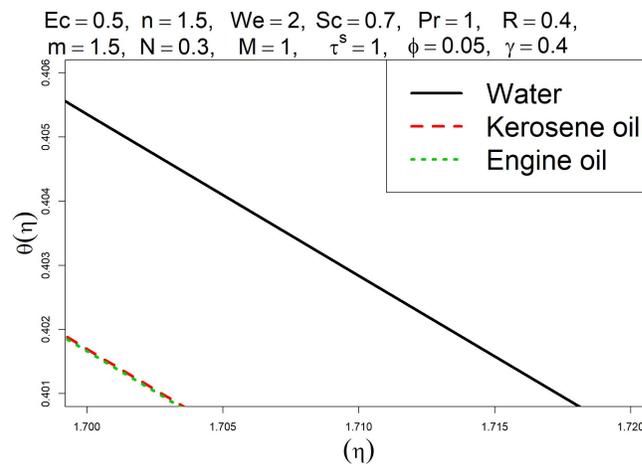


Fig. 15 Temperature profile for Shear thickening fluid with MWCNT

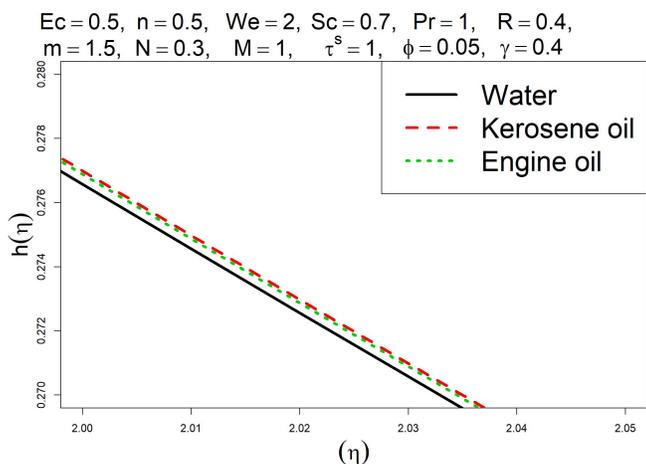


Fig. 16 Concentration profile for Shear thinning fluid with MWCNT

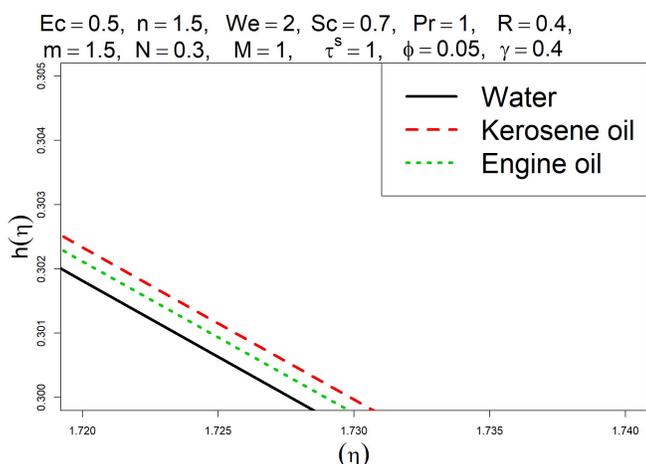


Fig. 17 Concentration profile for Shear thickening fluid with MWCNT

Table 1 A Comparison of the values of $f''(0)$ with $n=1.0$, $We=0.0$, $M=0.0$, $\gamma=0.0$, $Pr=1.0$, $Ec=0.0$, $Sc=0.0$, $R=0.0$, $\phi=0.0$ for different values of m .

m	(Khan and Hashim, 2015)	Present study
0.0	0.6275549	0.6275627
0.5	0.889544	0.8895524
1.0	1.0	1.0
1.5	1.061601	1.0616094
3.0	1.148593	1.1486015

Table 2 A Comparison of the values of $\theta'(0)$ with $n=1.0$, $We=0.0$, $M=0.0$, $\gamma=0.0$, $Pr=1.0$, $Ec=0.0$, $Sc=0.0$, $R=0.0$, $\phi=0.0$ for different values of m .

m	(Khan and Hashim, 2015)	Present study
0.2	0.610202	0.6102159
0.5	0.595201	0.5952237
1.5	0.574730	0.5747717
3.0	0.564662	0.5647180
10.0	0.554951	0.5549518

Table 3 Engineering parameters.

M	m=1,we=3.0 Pr=1.0,N=0.3,	$Re_x^{0.5} Cf_x$ n<1	$Re_x^{0.5} Cf_x$ n>1
0.0	$\gamma=0.4, Ec=0.5,$	-0.6972142	-1.078309
0.3	$Sc=0.7, \tau^s=0.3,$	-0.7764511	-1.249008
0.6	$\tau^s=0.3,$	-0.8461109	-1.409071
0.8	$\phi=0.05$	-0.8883215	-1.510718
we	m=1,M=3.0, Pr=1.0,N=0.3,	$Re_x^{0.5} Cf_x$ n<1	$Re_x^{0.5} Cf_x$ n>1
2.0	$\gamma=0.4, Ec=0.5,$	-0.9520452	-1.271478
3.0	$Sc=0.7, \tau^s=0.3,$	-0.8572339	-1.354897
3.5	$\tau^s=0.3,$	-0.8192214	-1.392846
4.0	$\phi=0.05$	-0.7862717	-1.428434
M	m=1,we=3.0 Pr=1.0,N=0.3,	$Re_x^{-0.5} Nu_x$ n<1	$Re_x^{-0.5} Nu_x$ n>1
0.0	$\gamma=0.4, Ec=0.5,$	0.3741720	0.4515941
0.3	$Sc=0.7, \tau^s=0.3,$	0.3542758	0.4706764
0.6	$\tau^s=0.3,$	0.3274603	0.4864306
0.8	$\phi=0.05$	0.3071349	0.4956823
we	m=1,M=3.0, Pr=1.0,N=0.3,	$Re_x^{-0.5} Nu_x$ n<1	$Re_x^{-0.5} Nu_x$ n>1
2.0	$\gamma=0.4, Ec=0.5,$	0.4060392	0.4836219
3.0	$Sc=0.7, \tau^s=0.3,$	0.3760004	0.4990199
3.5	$\tau^s=0.3,$	0.3630353	0.5056448
4.0	$\phi=0.05$	0.3513489	0.5116711
M	m=1,we=3.0 Pr=1.0,N=0.3,	$Re_x^{-0.5} Sh_x$ n<1	$Re_x^{-0.5} Sh_x$ n>1
0.0	$\gamma=0.4, Ec=0.5,$	0.5992370	0.6526943
0.3	$Sc=0.7, \tau^s=0.3,$	0.5699964	0.6406981
0.6	$\tau^s=0.3,$	0.5418290	0.6297297
0.8	$\phi=0.05$	0.5238310	0.6229667
we	m=1,M=3.0, Pr=1.0,N=0.3,	$Re_x^{-0.5} Sh_x$ n<1	$Re_x^{-0.5} Sh_x$ n>1
2.0	$\gamma=0.4, Ec=0.5,$	0.5897155	0.6324397
3.0	$Sc=0.7, \tau^s=0.3,$	0.5750112	0.6416226
3.5	$\tau^s=0.3,$	0.5688794	0.6455807
4.0	$\phi=0.05$	0.5634512	0.6491782

5. CONCLUSIONS

MHD flow of Carreau nanofluid explored using CNT over a nonlinear stretched sheet with thermal radiation, buoyancy and magnetic parameters has been studied. This investigation has explored that threshold value of momentum boundary layer is less than the threshold value of thermal and concentration boundary layers because of the effect of radiation and Local buoyancy parameter.

It was observed that the velocity profile diminished for $n<1$ and the contrary was true for $n>1$ with increasing value of Weissenberg number (we). Moreover opposite behavior was observed for the temperature and concentration profiles. The temperature and thermal boundary layer thickness was decreased by the increasing Prandtl number (Pr). Increasing phenomenon of Schmidt number leads to low nanoparticle concentration. Investigators observed that radiation is one of the cause to reduce the boundary layer of Concentration with Local buoyancy parameter relative to thermophoretic parameter. Investigators noticed that due to thermal conductivity and specific heat of Water, water based carbon nanotube (SWCNT and MWCNT) profiles are enhanced when compared to the remaining base fluids except at concentration profile, because of thermophoretic parameter.

This study provides information and limitations of the use of the Single walled carbon nanotube with water as base fluid. The thermal performance of the nanofluids are not much different and it may be important to observe other nanofluids which may have higher heat transfer rate. Nu-

merical Simulations were obtained by bvp solver using R-Software.

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NOMENCLATURE

x, y	Cartesian coordinates
p, q	Velocity components in x, y directions respectively
B	Magnetic field intensity
B_0	Constant magnetic field intensity
g	Gravitational acceleration
T	Temperature of the flow field
C	Species concentration
T_w	Temperature at the wall
T_∞	Temperature of the fluid outside the boundary layer
T_{ref}	Reference temperature
C_w	Concentration at the wall
C_∞	Concentration outside the boundary layer
D	Diffusion coefficient
V_T	Thermophoretic velocity
k_1	Thermophoretic coefficient
c_p	Specific heat
k_{nf}	Thermal conductivity of nanofluid
k_{bf}	Thermal conductivity of basefluid
k_{CNT}	Thermal conductivity of CNT
k^s	Mean absorption coefficient
we_x	Local Weissenberg number
Gr_x	Local Grashof number
Re_x	Local Reynolds number
Pr	Prandtl number
Ec	Eckert number
Sc	Schmidt number
M	Magnetic parameter
N	Concentration buoyancy parameter
R	Radiation parameter
n	Power law index
m, b	positive real numbers relating to the stretching
<i>Greek Symbols</i>	
ρ	Fluid density
γ	Local buoyancy parameter
τ	Material time constant
β_T	Thermal expansion coefficients of temperature
β_C	Thermal expansion coefficients of Concentration
σ^s	Stefan-Boltzmann constant
ν_{nf}	Coefficient of kinematics viscosity of nanofluid
μ_{nf}	Coefficient of viscosity of nanofluid
μ_{bf}	Coefficient of viscosity of basefluid
ϕ	Nanoparticle volume fraction
σ_{nf}	Electrical conductivity of nanofluid
σ_{bf}	Electrical conductivity of basefluid
σ_{CNT}	Electrical conductivity of CNT
ψ	Stream function
η	Similarity variable
θ	Dimensionless fluid temperature
h	Dimensionless fluid Concentration
τ^s	Thermophoretic parameter

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