

Analytical solution to fractional differential equation arising in thermodynamics

Areej Almuneef¹, Ahmed Eissa Hagag²

¹ Department of Mathematical Science, College of Science, Princess Nourah bint Abdulrahman University, Saudi Arabia

² Department of Basic Science, Faculty of Engineering, Sinai University, Ismailia, Egypt

Abstract

The analysis of nonlinear events related to physical phenomena is a popular issue in the modern-day. The essential purpose of this work is to discover a novel approximate solution to the fractional nonlinear Benjamin Bona Mahony Peregrine Burgers equation (BBMPB) utilizing the natural decomposition method (NDM) of fractional order. The suggested approach provides analytical solutions that are extremely near to the exact solution whereas obviating the complexities associated with many other approaches. The expected issue's uniqueness theorem and convergence analysis are explored using Banach's fixed-point theory. The reliability and accuracy of the recommended method were tested using numerical simulations. The graphs and tables reflect the results. The comparison of the suggested scheme's solution with the exact solutions demonstrates that the scheme is efficient, methodical, and extremely exact in tackling nonlinear complicated phenomena.

OPEN ACCESS

Published: 29/12/2023

Accepted: 21/12/2023

Submitted: 22/06/2023

DOI:
10.23967/j.rimni.2024.01.001

Keywords:

Fractional nonlinear Benjamin Bona Mahony Peregrine Burgers equation
Convergence analysis
Fractional natural decomposition method

Analytical solution to fractional differential equation arising in thermodynamics

Areej Almuneef^{1,a}, Ahmed Eissa Hagag^{2,b,*}

¹Department of Mathematical Science, College of Science, Princess Nourah bint Abdulrahman University, P.O. Box 105862, Riyadh 11656, Saudi Arabia

²Department of Basic Science, Faculty of Engineering, Sinai University, Ismailia, Egypt

^{1,a} aaalmuneef@pnu.edu.sa

^{2,b,*} E-mail:ahmed.shehata@su.edu.eg

The analysis of nonlinear events related to physical phenomena is a popular issue in the modern day. The essential purpose of this work is to discover a novel approximate solution to the fractional nonlinear Benjamin Bona Mahony Peregrine Burgers equation (BBMPB) utilizing the natural decomposition method (NDM) of fractional order. The suggested approach provides analytical solutions that are extremely near to the exact solution obviating the complexities associated with many other approaches. The expected issue's uniqueness theorem and convergence analysis are explored using Banach's fixed-point theory. The reliability and accuracy of the recommended method were tested using numerical simulations. The graphs and tables reflect the results. The comparison of the suggested scheme's solution with the exact solutions demonstrates that the scheme is efficient, methodical, and extremely exact in tackling nonlinear complicated phenomena.

1 Introduction

Differential equations (DEs) are becoming extremely important

in industrial applications. These are necessary and stimulating since the majority of the operations are connected with rates of change, which are clearly shown by them. In particular, DEs provide concepts for analyzing occurrences and creating ideas in medicine, finance, engineering, economics, and other related fields of research [1-2]. The analysis and examination of these kinds of equations are based on the survey of the foundations that govern the majority of physical phenomena. Furthermore, the analysis of nonlinear systems using fractional operators is crucial for studying phenomena in everyday life. While illustrating real-world issues connected with complexity, the researchers investigated its characteristics in greater depth and discovered that each notion has its own boundaries. However, several scholars discovered numerous limits and flaws in classical calculus while researching problems involving memory or hereditary characteristics. Many researchers use the core concepts and accompanying principles of FC to illustrate their points of view on many types of nonlinear phenomena [3-5]. They later suggested additional operators defined using fractional order. Accordingly, many scholars are drawn to the notion of fractional calculus while examining various models [6-8].

The research on nonlinear analysis in relation to the everyday demands of living beings drew the attention of all scholars due to its importance in modernization. Finding the solution for the relevant system is as important as modeling with mathematical tools. In this way, there are various techniques accessible in the literature [9-11]. Furthermore, each algorithm has its own set of requirements as well as its own set of restrictions. On the other hand, scholars are developing new techniques by overcoming constraints such as large computations, low precision, complex

procedures, calculating time, and so on. There are several strategies available in the literature, many of which are quite accurate. The Adomian decomposition technique is one of the approaches with excellent accuracy and dependability [12-13]. Researchers are always exploring and attempting to suggest new techniques by altering, fostering, combining, or upgrading current ones. In this way, the researchers proposed a new method by introducing natural transform (NT) to the ADM, which is called the natural decomposition method (NDM) in the classical order [14-16]. Then this method was generalized and presented in fractional order [17-20].

In many areas of mathematics and science, pseudo-parabolic equations are found, and the highest-order term in these equations has a one-time derivative. They have been utilized to study clay consolidation, thermodynamics, shear in second-order fluids, fluid flow in fissured rock, and propagation of long waves with tiny amplitudes, among other things [21-24]. The generalized Benjamin-Bona-Mahony-Burgers (BBMB) equation is a significant particular instance of pseudo-parabolic-type equations that can be written as follows:

$$\omega_t - \omega_{xxt} - \alpha \omega_{xx} + \gamma \omega_x + g(\omega)_x = 0, \quad (1)$$

where γ denotes any genuine constant value, α denotes a positive constant, $\omega(x, t)$ is the horizontal fluid velocity, and $g(\omega)$ is a nonlinear C^2 smooth function. Peregrine [25] and Benjamin et al. [26] suggested an alternative regularised long-wave equation if $\alpha = 0$, $\gamma = 1$, and $g(\omega)_x = \omega \omega_x$ in Eq. (1), which is known as the Kortewegde Vries equation

$$\omega_t + \omega_{xxt} + \omega_x + \omega \omega_x = 0. \quad (2)$$

If $g(\omega)_x = \theta \omega \omega_x + \beta \omega_{xxx}$ in Eq. (1), then the generic form of the BBMPB equation is thus obtained as follows

$$\omega_t - \omega_{xxt} - \alpha \omega_{xx} + \gamma \omega_x + \theta \omega \omega_x + \beta \omega_{xxx} = 0. \quad (3)$$

If we put $\alpha = \beta = 0$ in Eq. (3), then this is the Benjamin-Bona-Mahony (BBM) equation in its general form

$$\omega_t - \omega_{xxt} + \gamma \omega_x + \theta \omega \omega_x = 0, \quad (4)$$

where $\theta \neq 0$ and γ are arbitrary constants. Eq. (4) contains various forms of BBM equations as seen in the research [27-30].

In this article, we use NDM to solve the fractional nonlinear Benjamin Bona Mahony Peregrine Burgers equation. In addition, the behavior of the results is described in terms of fractional order. The remainder of the work is organized as follows: in the next section, we explain the fundamental concepts of FC and NDM of fractional order, which are then used to obtain the needed results. Section 3 shows the basic solution method of the proposed technique using the Caputo fractional operator. Section 4 proves the proposed algorithm's uniqueness theorem and convergence analysis. In Section 5, we obtain the solution to the fractional nonlinear BBMPB problem using the fundamental NDM. In addition, we give the numerical results and graphs for the found solution in the same section. Finally, we draw conclusions about the stored findings in terms of the considered technique and model.

2 Preliminaries to FC

2.1 Definition I

The fractional integral operator of Riemann-Liouville of a

function $\omega(\theta) \in C_\zeta, \zeta \geq -1$ is defined as [3]

$$J_a^\zeta \omega(\theta) = \begin{cases} \int_a^\theta \omega(\theta) (d\theta)^\zeta = \frac{1}{\Gamma(\zeta)} \int_a^\theta (\theta - v)^{\zeta-1} \omega(v) dv, & \theta > a, \\ \omega(\theta), & \theta = a. \end{cases} \quad (5)$$

2.2 Definition II

The Caputo fractional differential operator of order $\zeta > 0$ is defined as [3]

$$D_t^\zeta \omega(\theta) = \begin{cases} \frac{d^k}{d\theta^k} \omega(\theta), & \zeta = k \in N, \\ \frac{1}{\Gamma(k-\zeta)} \int_0^\theta (\theta - v)^{k-\zeta-1} \omega^{(k)}(v) dv, & k-1 < \zeta \leq k \in N. \end{cases} \quad (6)$$

2.3 Definition III

The Mittag-Leffler of a one-parameter function $E_\zeta(\theta)$ with $\zeta > 0$ is given as [31]

$$E_\zeta(\theta) = \sum_{i=0}^{\infty} \frac{\theta^i}{\Gamma(i\zeta + 1)} \quad (7)$$

2.4 Definition IV

The natural transform (NT) of $\omega(\theta)$, which is defined as

$$N^+ \{ \omega(\theta) \} = \int_{-\infty}^{\infty} e^{-s\theta} \omega(z\theta) d\theta, \quad s, z \in (-\infty, \infty). \quad (8)$$

2.5 Definition V

The effect of the natural transform on the Caputo operator is given as [32]

$$N^+ \{ D_t^\zeta \omega(\theta) \} = \frac{z^\zeta}{s^\zeta} N^+ \{ \omega(\theta) \} - \sum_{i=0}^{n-1} \frac{z^{i-\zeta}}{s^{i+1-\zeta}} [D^i \omega(\theta)]_{\theta=0}, \quad n-1 < \zeta \leq n. \quad (9)$$

3 Construction of fractional NDM

We consider a general form of fractional nonlinear partial differential equation to demonstrate the underlying theory and solution technique of the suggested approach as

$$D_t^\zeta \omega(x, t) + R\omega(x, t) + F\omega(x, t) = h(x, t), \quad n-1 < \zeta \leq (10)$$

with the initial condition

$$\omega(x, 0) = v(x), \quad (11)$$

where $D_t^\zeta = \frac{\partial^\zeta}{\partial t^\zeta}$ denotes the Caputo operator of $\omega(x, t)$, R denotes the linear function, F denotes the non-linear function and $h(x, t)$ signifies the source term. Using the NT on Eq. (10),

we get

$$N^+ [D_t^\zeta \omega(x, t)] + N^+ [R\omega(x, t)] + N^+ [F\omega(x, t)] = N^+ [h(x, t)].$$

Applying definition 5, we get

$$N^+ [\omega(x, t)] = \frac{Z^\zeta}{S^\zeta} \sum_{i=0}^{n-1} \frac{Z^{i-\zeta}}{S^{i+1-\zeta}} [D^i \omega(t)]_{t=0} + \frac{Z^\zeta}{S^\zeta} N^+ [h(x, t)] + \frac{Z^\zeta}{S^\zeta} N^+ [R\omega(x, t) + F\omega(x, t)].$$

Utilize the inverse NT on the above equation to obtain

$$\omega(x, t) = H(x, t) + N^- \left\{ \frac{Z^\zeta}{S^\zeta} N^+ [h(x, t) - R\omega(x, t) - F\omega(x, t)] \right\}.$$

$H(x, t)$ identified using nonhomogeneous terms and the provided guess condition. The infinite series solution is given as

$$\omega(x, t) = \sum_{n=0}^{\infty} \omega_n(x, t), \quad F\omega(x, t) = \sum_{n=0}^{\infty} A_n, \quad (14)$$

where A_n signifies the nonlinear component of $F\omega(x, t)$, and we have

$$\sum_{n=0}^{\infty} \omega_n(x, t) = H(x, t) + N^- \left\{ \frac{Z^\zeta}{S^\zeta} N^+ [h(x, t) - R\omega(x, t) - F\omega(x, t)] \right\}.$$

Lastly, the analytical solutions are provided in the form of

$$\omega(x, t) = \sum_{n=0}^{\infty} \omega_n(x, t). \quad (16)$$

4 Convergence analysis of NDM

The uniqueness and existence theorems are the instruments that lead one to infer that there is only one solution that satisfies a specific initial condition for a given problem.

4.1 Theorem 1

The solution provided with the aid of NDM for the BBMPM equation is unique wherever $\chi \in (0, 1)$, where

$$\chi = \{e^3 + \alpha\varphi^2 - \gamma\varphi - \theta(X + Y) - \beta\varphi^3\} \Xi \quad (17)$$

Proof

The analytical solution determined for the BBMPM equation is given as

$$\omega(x, t) = \sum_{j=0}^{\infty} \omega_j(x, t) q^j, \quad (18)$$

where

$$\omega_{m+1}(x, t) = N^- \left\{ \frac{\phi(x)}{S} \right\} + N^- \left\{ \left(\frac{Z^\zeta}{S^\zeta} \right) N^+ [(\omega_m)_{xxt} + \alpha(\omega_m)_{xx} - \gamma(\omega_m)_x - \theta \sum_{i=0}^m \omega_i(\omega_{m-i})_x - \beta(\omega_m)_{3x}] \right\}.$$

Let ω and ω be the two solutions for the BBMPM equation such that $|\omega| \leq X$ and $|\omega| \leq Y$, then usage of the equation above, we obtain

$$|\omega - \omega| = |N^- \left\{ \left(\frac{Z^\zeta}{S^\zeta} \right) N^+ [(\omega_{xxt} - \omega_{xxt}) + \alpha(\omega_{xx} - \omega_{xx}) - \gamma(\omega_x - \omega_x) - \theta(\omega^2 - \omega^2)] \right\}|$$

Transform by using the convolution principle for NT, we obtain

$$\begin{aligned} |\omega - \omega| &= \int_0^\phi \left\{ |\omega_{2xt} - \omega_{2xt}| + \alpha |\omega_{2x} - \omega_{2x}| - \gamma |\omega_x - \omega_x| - \theta |\omega^2 - \omega^2| \right\} dt \\ &\leq \int_0^\phi \left\{ \frac{\partial^3}{\partial x^2 \partial t} |\omega - \omega| + \frac{\alpha \partial^2}{\partial x^2} |\omega - \omega| - \frac{\gamma \partial}{\partial x} |\omega - \omega| - \theta |\omega - \omega| \right\} dt \\ &\leq \left\{ \frac{\partial^3}{\partial x^2 \partial t} |\omega - \omega| + \frac{\alpha \partial^2}{\partial x^2} |\omega - \omega| - \frac{\gamma \partial}{\partial x} |\omega - \omega| - \theta |\omega - \omega| \right\} \phi \\ &\leq \{e^3 |\omega - \omega| + \alpha\varphi^2 |\omega - \omega| - \gamma\varphi |\omega - \omega| - \theta(X + Y) |\omega - \omega| - \beta\varphi^3 |\omega - \omega|\} \end{aligned}$$

where $\varphi^n = \frac{\partial^n}{\partial x^n}$, $n = 1, 2, 3$ and $e^3 = \frac{\partial^3}{\partial x^2 \partial t}$. To minimise the previous equation as follows, we can use the integral mean value [33]

$$\begin{aligned} |\omega - \omega| &\leq \{e^3 |\omega - \omega| + \alpha\varphi^2 |\omega - \omega| - \gamma\varphi |\omega - \omega| - \theta(X + Y) |\omega - \omega| - \beta\varphi^3 |\omega - \omega|\} \phi \\ &\leq \chi |\omega - \omega|, \end{aligned}$$

$\therefore (1 - \chi) |\omega - \omega| \leq 0$. Since $0 < \chi < 1$, therefore $|\omega - \omega| = 0$, which gives $\omega = \omega$, where $\omega = \frac{\phi^{\zeta+1}}{\Gamma(\zeta+2)}$. Hence, the analytical solution is unique.

4.2 Theorem 2

Presume that

$$\|F(\omega) - F(s)\| \leq \chi \|\omega - s\|, \forall \omega, s \in B, \chi \in (0, 1),$$

where B is a Banach space with $F: B \rightarrow B$. The preceding theorem and the fixed-point principle of Banach [34] was used to infer that F has a fixed point. Furthermore, the analytical solution acquired utilizing the suggested procedure converges with a random election for $\omega_0, s_0 \in B$ to a fixed point of F and

$$\|\omega_\mu - \omega_\sigma\| \leq \frac{\chi^\phi}{1 - \chi} \|\omega_1 - \omega_0\|. \quad (24)$$

Proof

Assume that B a Banach space $(C[J], \|\cdot\|)$ of all continuous functions. We can agree that $\{\omega_\mu\}$ is a Cauchy sequence in the Banach space as

$$\begin{aligned} \|\omega_\mu - \omega_\sigma\| &= \max_{\phi \in J} |\omega_\mu - \omega_\sigma| \\ &= \max_{\phi \in J} |N^- \left\{ \left(\frac{Z^\zeta}{S^\zeta} \right) N^+ \left[\left(\frac{\partial^3 \omega_{\mu-1}}{\partial x^2 \partial t} - \frac{\partial^3 \omega_{\sigma-1}}{\partial x^2 \partial t} \right) + \alpha \left(\frac{\partial^2 \omega_{\mu-1}}{\partial x^2} - \frac{\partial^2 \omega_{\sigma-1}}{\partial x^2} \right) - \theta (\omega_{\mu-1}^2 - \omega_{\sigma-1}^2) - \beta \left(\frac{\partial^3 \omega_{\mu-1}}{\partial x^3} - \frac{\partial^3 \omega_{\sigma-1}}{\partial x^3} \right) \right] \right\}| \\ &\leq \max_{\phi \in J} \left[N^- \left\{ \left(\frac{Z^\zeta}{S^\zeta} \right) N^+ \left[\left| \frac{\partial^3 \omega_{\mu-1}}{\partial x^2 \partial t} - \frac{\partial^3 \omega_{\sigma-1}}{\partial x^2 \partial t} \right| + \alpha \left| \frac{\partial^2 \omega_{\mu-1}}{\partial x^2} - \frac{\partial^2 \omega_{\sigma-1}}{\partial x^2} \right| - \theta |\omega_{\mu-1} - \omega_{\sigma-1}| + |\omega_{\mu-1} + \omega_{\sigma-1}| - \beta \left| \frac{\partial^3 \omega_{\mu-1}}{\partial x^3} - \frac{\partial^3 \omega_{\sigma-1}}{\partial x^3} \right| \right] \right\} \right] \end{aligned}$$

Transform by using the convolution principle for NT, we obtain

$$\| \omega_\mu - \omega_\sigma \| \leq \max_{\phi \in J} [\{ \epsilon^3 | \omega_{\mu-1} - \omega_{\sigma-1} | + \alpha \phi^2 | \omega_{\mu-1} - \omega_{\sigma-1} | - \gamma \phi | \omega_{\mu-1} - \omega_{\sigma-1} | - \theta (X+Y) | \omega_{\mu-1} - \omega_{\sigma-1} | - \beta \phi^3 | \omega_{\mu-1} - \omega_{\sigma-1} | \} \int_0^\phi (\phi - \tau)^\zeta d\tau]$$

To minimise the previous equation as follows, we can use the integral mean value [33]

$$\| \omega_\mu - \omega_\sigma \| \leq \max_{\phi \in J} [\{ \epsilon^3 | \omega_{\mu-1} - \omega_{\sigma-1} | + \alpha \phi^2 | \omega_{\mu-1} - \omega_{\sigma-1} | - \gamma \phi | \omega_{\mu-1} - \omega_{\sigma-1} | - \theta (X+Y) | \omega_{\mu-1} - \omega_{\sigma-1} | - \beta \phi^3 | \omega_{\mu-1} - \omega_{\sigma-1} | \} \chi \| \omega_{\mu-1} - \omega_{\sigma-1} \|]$$

Subtracting μ by $\sigma + 1$, we obtain

$$\| \omega_{\sigma+1} - \omega_\sigma \| \leq \chi \| \omega_\sigma - \omega_{\sigma-1} \| \leq \chi^2 \| \omega_{\sigma-1} - \omega_{\sigma-2} \| \leq \dots \leq \chi^\sigma \| \omega_1 - \omega_0 \| \quad (29)$$

By using triangular inequality, we obtain

$$\begin{aligned} \| \omega_\mu - \omega_\sigma \| &= \| \omega_{\sigma+1} + \omega_{\sigma+2} + \dots + \omega_\mu - \omega_{\sigma+1} - \omega_{\sigma+2} - \dots - \omega_\sigma \| \\ &= \| \omega_{\sigma+1} + \omega_{\sigma+2} + \dots + \omega_\mu - \omega_{\mu-1} - \dots - \omega_{\sigma+2} - \omega_{\sigma+1} - \omega_\sigma \| \\ &\leq \{ \chi^\sigma + \chi^{\sigma+1} + \dots + \chi^{\mu-1} \} \| \omega_1 - \omega_0 \| \\ &\leq \chi^\sigma \{ 1 + \chi + \dots + \chi^{\mu-\sigma-1} \} \| \omega_1 - \omega_0 \| \\ &\leq \chi^\sigma \left\{ \frac{1 - \chi^{\mu-\sigma-1}}{1 - \chi} \right\} \| \omega_1 - \omega_0 \| \end{aligned}$$

As $\chi \in (0, 1)$, so $1 - \chi^{\mu-\sigma-1} < 1$, then we get

$$\| \omega_\mu - \omega_\sigma \| \leq \frac{\chi^\sigma}{1 - \chi} \| \omega_1 - \omega_0 \| \quad (31)$$

Since $\| \omega_1 - \omega_0 \| < \infty$, we find that $\| \omega_\mu - \omega_\sigma \| \rightarrow 0$ when μ and $\sigma \rightarrow \infty$. This shows that the sequence $\{ \omega_\mu \}$ generated by NDM is a convergent Cauchy sequence.

5 Solution for BBMPB equation

To offer the solution to the relevant problem, we will use the fractional natural decomposition approach. We will provide four examples to demonstrate the dependability of the proposed method. In this part, we will look at the new fractional Benjamin Bona Mahony Peregrine Burgers equation, which is stated as follows

$$\frac{\partial^\zeta \omega(x, t)}{\partial t^\zeta} = \frac{\partial^3 \omega(x, t)}{\partial x^2 \partial t} + \alpha \frac{\partial^2 \omega(x, t)}{\partial x^2} - \gamma \frac{\partial \omega(x, t)}{\partial x} - \theta \omega(x, t) - \beta \frac{\partial^3 \omega(x, t)}{\partial x^3}$$

in the operator form, with initial condition

$$\omega(x, 0) = \phi(x) \quad (33)$$

Using NT on Eq.(32), one may obtain

$$N^+ [D_t^\zeta \omega(x, t)] = N^+ \left[\frac{\partial^3 \omega}{\partial x^2 \partial t} + \alpha \frac{\partial^2 \omega}{\partial x^2} - \gamma \frac{\partial \omega}{\partial x} - \theta \omega - \beta \frac{\partial^3 \omega}{\partial x^3} \right]$$

By using the natural transformation, we find that

$$\frac{z^{-\zeta}}{z^\zeta} N^+ [\omega(x, t)] - n - 1 \sum_{s=1}^n \frac{z^{-i-\zeta}}{s^{i+1-\zeta}} [D^i \omega]_{t=0} = N^+ \left[\frac{\partial^3 \omega}{\partial x^2 \partial t} + \alpha \frac{\partial^2 \omega}{\partial x^2} - \gamma \frac{\partial \omega}{\partial x} - \theta \omega - \beta \frac{\partial^3 \omega}{\partial x^3} \right]$$

By Eqs. (34) and (35), we find that

$$N^+ [\omega(x, t)] = \frac{\phi(x)}{s} + \frac{z^{-\zeta}}{s^\zeta} N^+ \left[\frac{\partial^3 \omega}{\partial x^2 \partial t} + \alpha \frac{\partial^2 \omega}{\partial x^2} - \gamma \frac{\partial \omega}{\partial x} - \theta \omega - \beta \frac{\partial^3 \omega}{\partial x^3} \right]$$

By using the inverse NT to above Eq.

$$\omega(x, t) = \phi(x) + N^+ \left\{ \frac{z^{-\zeta}}{s^\zeta} N^+ \left[\frac{\partial^3 \omega}{\partial x^2 \partial t} + \alpha \frac{\partial^2 \omega}{\partial x^2} - \gamma \frac{\partial \omega}{\partial x} - \theta \omega - \beta \frac{\partial^3 \omega}{\partial x^3} \right] \right\}$$

Suppose that the solution to the infinite series of $\omega(x, t) = \sum_{n=0}^{\infty} \omega_n(x, t)$. Keep in mind $\omega_{xx} = \sum_{n=0}^{\infty} A_n$ is the Adomian polynomial and signify the nonlinear terms. Eq. (37) may be rewritten using this term as

$$\sum_{n=0}^{\infty} \omega_n(x, t) = \phi(x) + N^+ \left\{ \frac{z^{-\zeta}}{s^\zeta} N^+ \left[\sum_{n=0}^{\infty} \omega_{nxt} + \alpha \sum_{n=0}^{\infty} \omega_{xx} - \gamma \sum_{n=0}^{\infty} \omega_x - \theta \sum_{n=0}^{\infty} A_n - \beta \sum_{n=0}^{\infty} A_n \right] \right\}$$

5.1 Application 1

The initial condition for Eq. (32) take the following form [35]

$$\omega(x, 0) = - \left(\frac{\beta + \gamma}{\theta} \right) + \frac{\beta + \gamma}{\theta} \tanh \left(\frac{-\beta - \gamma}{2\delta} x \right) \quad (39)$$

If $\alpha = \beta = \gamma = \theta = 1$, we get

$$\omega_1 = \frac{3t^\zeta (\cosh(2x) - 3) \operatorname{sech}^4(x)}{\Gamma(\zeta + 1)} \quad (40)$$

$$\begin{aligned} \omega_2 = \frac{3t^{2\zeta} \operatorname{sech}^7(x)}{4\Gamma(2\zeta + 1)} &\{ 3[546 \sinh(x) - 93 \sinh(3x) + \sinh(5x)] + 64 \cosh(x) \} \\ &+ \frac{6\zeta t^{2\zeta-1}}{\Gamma(2\zeta + 1)} \{ -34 \cosh(2x) + \cosh(4x) + 45 \} \operatorname{sech}^6(x) \end{aligned}$$

The prior analytical solution leads to the following exact solution [35]

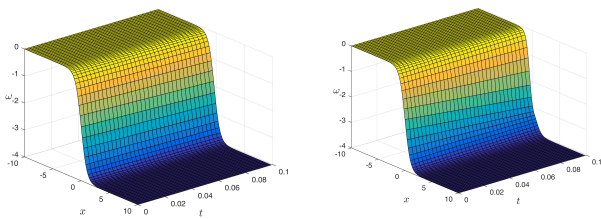
$$\omega(x, t) = - \left(\frac{\beta + \gamma}{\theta} \right) + \frac{\beta + \gamma}{\theta} \tanh \left(\frac{-\beta - \gamma}{2\delta} (x + \beta t) \right) \quad (42)$$

For the specific instance when $\zeta = 1$, the approximative findings and table 1 demonstrate that the exact solution of Eq. (32) has a generic type that is equivalent to the aforementioned analytical solutions. In order to comprehend the geometric behavior of our approximation to Eq. (32), shown in Figure 1, the exact solution in two and three dimensions is compared to the 2nd iteration of NTM. When $\zeta = 1$, $\zeta = 0.95$, $\zeta = 0.90$ and $\zeta = 0.80$, the NDM solution and the exact solution were also compared.

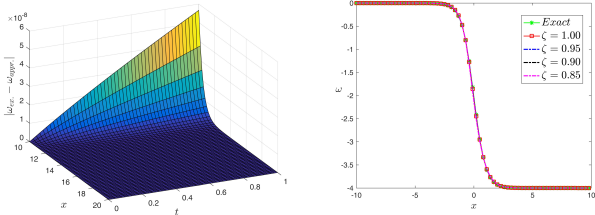
Table 1: Comparison of case 1 for the exact solution with NDM solution and absolute errors using numerical calculations at $t = 1$ and $\alpha = \beta = \gamma = \theta = 1$.

x	ω_{Ex}	$\omega_{NDM}(\zeta = 1)$	Absolute error	$\omega_{NDM}(\zeta = 0.95)$	$\omega_{NDM}(\zeta = 0.90)$
1	0.9901	-3.5640	4.26060	-3.556574	-3.549577
2	0.9507	48679	8E-01	735	917
3	-3.9998	-3.9906	9.18393	-3.990469	-3.990315
4	18408	34473	5E-03	718	480
5	-3.9999	-3.9998	1.68702	-3.999824	-3.999822
6	0.9999	0.9999	1.68702	0.9999	0.9999
7	0.9999	0.9999	1.68702	0.9999	0.9999
8	0.9999	0.9999	1.68702	0.9999	0.9999
9	0.9999	0.9999	1.68702	0.9999	0.9999
10	0.9999	0.9999	1.68702	0.9999	0.9999

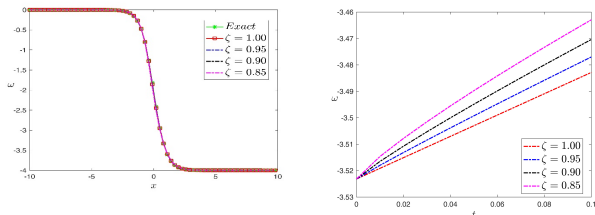
	-3.9999	-3.9999	3.09006	-3.999996	-3.999996
8	99939	96849	1E-06	793	741
	-3.9999	-3.9999	5.65965	-3.999999	-3.999999
10	99998	99942	1E-08	941	940
	-3.9999	-3.9999	1.03660	-3.999999	-3.999999
12	99999	99998	1E-09	998	998
	-3.9999	-3.9999	1.89859	-3.999999	-3.999999
14	99999	99999	2E-11	999	999
	-3.9999	-3.9999	3.47721	-3.999999	-3.999999
16	99999	99999	8E-13	999	999
	-3.9999	-3.9999	6.21724	-3.999999	-3.999999
18	-4	99999	8E-15	999	999
20	-4	-4	0	-4	-4



(a) Exact solution of ω(x,t). (b) NDM solution of ω(x,t).



(c) Absolute error at t = 0.01. (d) Comparison of results at t = 0.01.



(e) Comparison between exact and various values of ζ. (f) Comparison between various values of ζ at x = 1.

Figure 1: Periodic wave analytical (NDM) solutions ω(x,t) of Eq. (32) with initial condition (39) and α = β = γ = θ = 1.

5.2 Application 2

The initial condition for Eq. (32) take the following form [35]

$$\omega(x, 0) = -\left(\frac{\beta + \gamma}{\theta}\right) - \frac{\beta + \gamma}{\theta} \coth\left(\frac{-\beta - \gamma}{2\delta} x\right), \quad (43)$$

If α = β = γ = θ = 1, we get

$$\omega_1 = -\frac{3t^\zeta (\cosh(2x) + 3) \operatorname{csch}^4(x)}{\Gamma(\zeta + 1)}, \quad (44)$$

$$\omega_2 = \frac{3t^{2\zeta} \operatorname{csch}^7(x)}{4\Gamma(2\zeta + 1)} \{64\sinh(x) + 32\sinh(3x) + 1638\cosh(x) + 279\cosh(3x)\} + \frac{3t^{2\zeta-1}}{\Gamma(2\zeta)} (34\cosh(2x) + \cosh(4x) + 45) \operatorname{csch}^6(x).$$

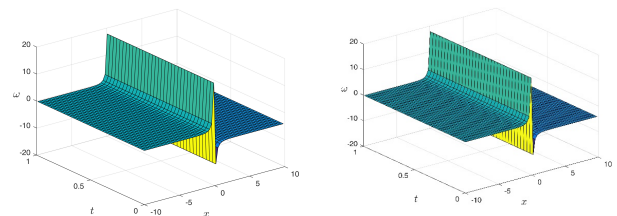
The prior analytical solution leads to the following exact solution [35]

$$\omega(x, t) = -\left(\frac{\beta + \gamma}{\theta}\right) - \frac{\beta + \gamma}{\theta} \coth\left(\frac{-\beta - \gamma}{2\delta}(x + \beta t)\right). \quad (46)$$

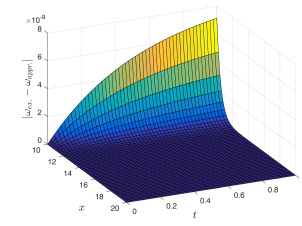
For the specific instance when ζ = 1, the approximative findings and table 2 demonstrate that the exact solution of Eq. (32) has a generic type that is equivalent to the aforementioned analytical solutions. In order to comprehend the geometric behavior of our approximation to Eq. (32), shown in Figure 2, the exact solution in two and three dimensions is compared to the 2nd iteration of NTM. When ζ = 1, ζ = 0.95, ζ = 0.90 and ζ = 0.80, the NTM solution and the exact solution were also compared.

Table 2: Comparison of case 2 for the exact solution with NDM solution and absolute errors using numerical calculations at t = 1 and α = β = γ = θ = 1.

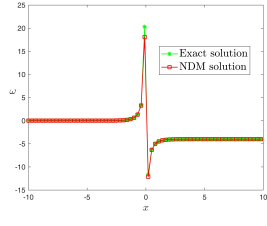
x	ω _{Ex}	ω _{NDM} (ζ = 1)	Absolute error	ω _{NDM} (ζ = 0.95)	ω _{NDM} (ζ = 0.90)
2	-4.0099	-4.6001	5.90172	-4.610901	-4.621002
	39646	12243	5E-01	692	357
4	-4.0001	-4.0094	9.23883	-4.009586	-4.009741
	81607	20444	6E-03	308	583
6	-4.0000	-4.0001	1.68721	-4.000175	-4.000177
	03326	72047	0E-04	075	909
8	-4.0000	-4.0000	3.09006	-4.000003	-4.000003
	00060	03150	8E-06	206	258
10	-4.0000	-4.0000	5.65965	-4.000000	-4.000000
	00001	00057	1E-08	058	059
12	-4.0000	-4.0000	1.03660	-4.000000	-4.000000
	00000	00001	1E-09	001	001
14	-4.0000	-4.0000	1.89859	-4.000000	-4.000000
	00000	00000	8E-11	000	000
16	-4.0000	-4.0000	3.47460	-4.000000	-4.000000
	00000	00000	7E-13	000	000
18	-4.0000	-4.0000	6.01094	-4.000000	-4.000000
	00000	00000	4E-15	000	000
20	-4.0000	-4.0000	1.01960	-4.000000	-4.000000
	00000	00000	5E-16	000	000



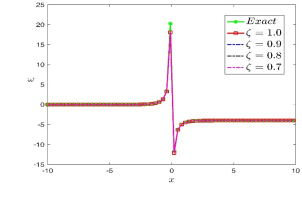
(a) Exact solution of ω(x,t). (b) NDM solution of ω(x,t).



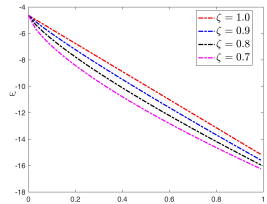
(c) Absolute error at t = 0.01.



(d) Comparison of results at t = 0.01.



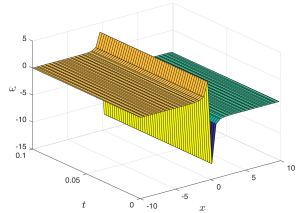
(e) Comparison between exact and various values of ζ.



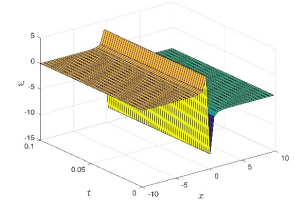
(f) Comparison between various values of ζ at x = 1.

Figure 2: Periodic wave analytical (NDM) solutions ω(x,t) of Eq. (32) with initial condition (42) and α = β = γ = θ = 1.

4	-4.0001 81607	-4.0094 20444	9.238836 E-03	-4.009586 308	-4.009741 583
6	-4.0000 03326	-4.0001 72047	1.687210 E-04	-4.000175 075	-4.000177 909
8	-4.0000 00060	-4.0000 03150	3.090068 E-06	-4.000003 206	-4.000003 258
10	-4.0000 00001	-4.0000 00057	5.659651 E-08	-4.000000 058	-4.000000 059
12	-4.0000 00000	-4.0000 00001	1.036601 E-09	-4.000000 001	-4.000000 001
14	-4.0000 00000	-4.0000 00000	1.898619 E-11	-4.000000 000	-4.000000 000
16	-4.0000 00000	-4.0000 00000	3.476832 E-13	-4.000000 000	-4.000000 000
18	-4.0000 00000	-4.0000 00000	6.403864 E-15	-4.000000 000	-4.000000 000
20	-4.0000 00000	-4.0000 00000	7.907517 E-17	-4.000000 000	-4.000000 000



(a) Exact solution of ω(x,t).



(b) NDM solution of ω(x,t).

5.3 Application 3

The initial condition for Eq. (32) take the following form [35]

$$\omega(x, 0) = -\left(\frac{\beta + \gamma}{\theta}\right) + \frac{\beta + \gamma}{2\theta} \tanh\left(\frac{-\beta - \gamma}{4\delta} x\right) + \frac{\beta + \gamma}{2\theta} \coth\left(\frac{-\beta - \gamma}{4\delta} x\right), \quad (47)$$

If α = β = γ = θ = 1, we get

$$\omega_1 = -\frac{3t^\alpha (\cosh(2x) + 3) \operatorname{csch}^4(x)}{\Gamma(\alpha + 1)}, \quad (48)$$

$$\omega_2 = -\frac{3t^{2\zeta} \operatorname{csch}^7(x)}{4\Gamma(2\zeta + 1)} \{64 \sinh(x) + 32 \sinh(3x) + 1638 \cosh(x) + \dots$$

$$- \frac{3t^{2\zeta - 1}}{\Gamma(2\zeta)} (34 \cosh(2x) + \cosh(4x) + 45) \operatorname{csch}^9(x), \dots$$

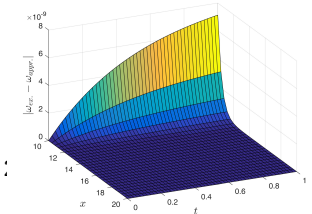
The prior analytical solution leads to the following exact solution [35]

$$\omega(x, t) = -\left(\frac{\beta + \gamma}{\theta}\right) + \frac{\beta + \gamma}{2\theta} \tanh\left(\frac{-\beta - \gamma}{4\delta} (x + \beta t)\right) + \frac{\beta + \gamma}{2\theta} \coth\left(\frac{-\beta - \gamma}{4\delta} (x + \beta t)\right)$$

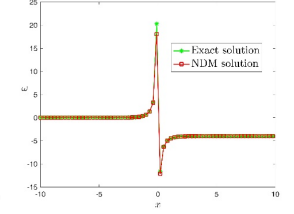
For the specific instance when ζ = 1, the approximative findings and table 3 demonstrate that the exact solution of Eq. (32) has a generic type that is equivalent to the aforementioned analytical solutions. In order to comprehend the geometric behavior of our approximation to Eq. (32), shown in Figure 3, the exact solution in two and three dimensions is compared to the 2nd iteration of NTM. When ζ = 1, ζ = 0.95, ζ = 0.90 and ζ = 0.80, the NTM solution and the exact solution were also compared.

Table 3: Comparison of case 3 for the exact solution with NDM solution and absolute errors using numerical calculations at t = 1 and α = β = γ = θ = 1.

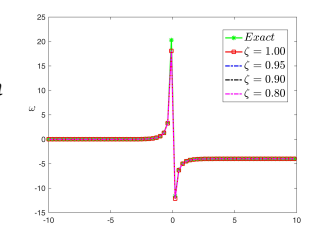
x	ω _{EX}	ω _{NDM} (ζ = 1)	Absolute error	ω _{NDM} (ζ = 0.95)	ω _{NDM} (ζ = 0.90)
2	-4.0099 39646	-4.6001 12243	5.901725 E-01	-4.610901 692	-4.621002 357



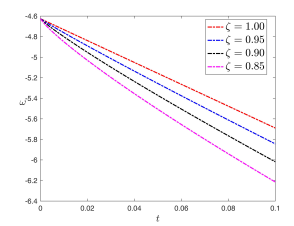
(c) Absolute error at t = 0.01.



(d) Comparison of results at t = 0.01.



(e) Comparison between exact and various values of ζ.



(f) Comparison between various values of ζ at x = 1.

Figure 3: Periodic wave analytical (NDM) solutions ω(x,t) of Eq. (32) with initial condition (46) and α = β = γ = θ = 1.

5.4 Application 4

The initial condition for Eq. (32) take the following form [35]

$$\omega(x, 0) = \frac{-10\beta\mu - 10\gamma\mu + \delta}{20\theta\mu} - \frac{(6\delta\mu) \tanh^2(\mu x)}{5\theta} - \frac{(12\delta\mu) \tanh(\mu x)}{5\theta}$$

where $\mu = \frac{5\beta + 5\gamma\sqrt{25\beta^2 + 50\beta\gamma + 25\gamma^2 - 24\delta^2}}{24\delta}$. If $\alpha = \beta = \gamma = \theta = 1$, we get

$$\omega_1 = \frac{t^\zeta}{10800\Gamma(\zeta + 1)} \left\{ -3(959\sqrt{19} + 4180)\sinh\left(\frac{1}{6}(\sqrt{19} + 5)x\right) + -95(79\sqrt{19} + 344) \right\} \left(\tanh\left(\frac{1}{12}(\sqrt{19} + 5)x\right) + \dots \right)$$

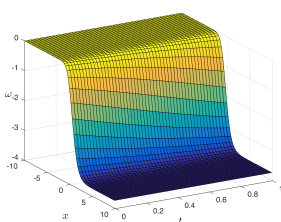
The prior analytical solution leads to the following exact solution [35]

$$\omega(x, t) = \frac{-10\beta\mu - 10\gamma\mu + \delta}{20\theta\mu} - \frac{(12\delta\mu)}{5\theta} \tanh(\mu(x - Vt)) - \frac{(6\delta\mu)}{5\theta}$$

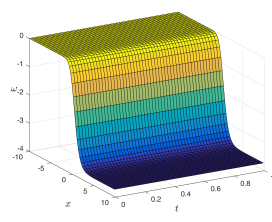
For the specific instance when $\zeta = 1$, the approximative findings and table 4 demonstrate that the exact solution of Eq. (32) has a generic type that is equivalent to the aforementioned analytical solutions. In order to comprehend the geometric behavior of our approximation to Eq. (32), shown in Figure 4, the exact solution in two and three dimensions is compared to the 2nd iteration of NTM. When $\zeta = 1, \zeta = 0.95, \zeta = 0.90$ and $\zeta = 0.80$, the NTM solution and the exact solution were also compared.

Table 4: Comparison of case 4 for the exact solution with NDM solution and absolute errors using numerical calculations at $t = 1$ and $\alpha = \beta = \gamma = \theta = 1$.

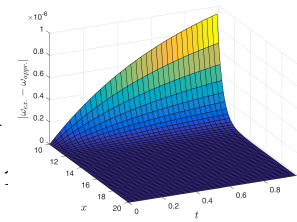
x	ω_{Ex}	$\omega_{NDM}(\zeta = 1)$	Absolute error	$\omega_{NDM}(\zeta = 0.95)$	$\omega_{NDM}(\zeta = 0.90)$
2	-3.6600	-3.2127	4.47357	-3.208190	-3.203946
	80833	23611	2E-01	131	062
4	-3.7398	-3.7015	3.82821	-3.700965	-3.700438
	12014	29867	4E-02	987	105
6	-3.7433	-3.7416	1.73519	-3.741632	-3.741608
	93916	58721	4E-03	941	807
8	-3.7435	-3.7434	7.67363	-3.743474	-3.743473
	52259	75523	0E-04	382	314
10	-3.7435	-3.7435	3.38987	-3.743555	-3.743555
	59254	55864	1E-06	813	766
12	-3.7435	-3.7435	1.49742	-3.743559	-3.743559
	59563	59413	4E-07	411	409
14	-3.7435	-3.7435	6.61462	-3.743559	-3.743559
	59576	59570	7E-09	570	569
16	-3.7435	-3.7435	2.92190	-3.743559	-3.743559
	59577	59577	2E-10	577	577
18	-3.7435	-3.7435	1.29074	-3.743559	-3.743559
	59577	59577	5E-11	577	577
20	-3.7435	-3.7435	5.70210	-3.743559	-3.743559
	59577	59577	5E-13	577	577



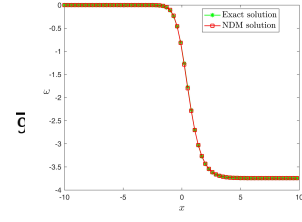
(a) Exact solution of $\omega(x, t)$.



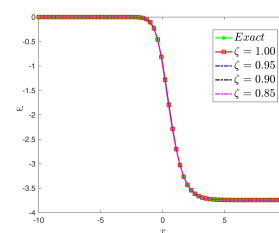
(b) NDM solution of $\omega(x, t)$.



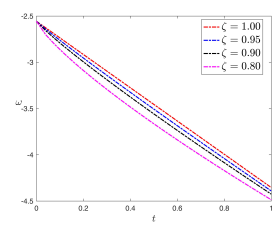
(c) Absolute error at $t = 0.01$.



(d) Comparison of results at $t = 0.01$.



(e) Comparison between exact and various values of ζ .



(f) Comparison between various values of ζ at $x = 1$.

Figure 4: Periodic wave analytical (NDM) solutions $\omega(x, t)$ of Eq. (32) with initial condition (50) and $\alpha = \beta = \gamma = \theta = 1$.

6 Conclusion

Studying and exploring nonlinear physical models using new techniques always help us advance in science and technology. In the current framework, we used NDM to evaluate the BBMPB equation with fractional order. Banach's fixed-point theory is used to investigate the anticipated issue's uniqueness theorem and convergence analysis. The anticipated method's dependability and applicability are demonstrated by presenting four cases. The behaviors for the obtained findings are provided in 2D, 3D graphs, and tables for featured fractional order. These graphs aid to conclude the stimulating behaviors of the analogical models. Furthermore, while solving nonlinear issues, NDM does not require any conversion, perturbation, or consideration of extra polynomials or parameters. The examination of these kinds of occurrences can provide new ideas for investigating more real-world events. It can also generate ideas for employing an accurate method to evaluate nonlinear models related to science and technology. This work elucidates the proposed model, which is notably dependent on time instant and its history and can be convincingly illustrated utilizing fractional notions.

Conflict of Interest: The author declare that there is no conflict of interest.

Data Availability: No data were used to support this study.

Acknowledgments: This research project was funded by the Deanship of Scientific Research, Princess Nourah bint Abdulrahman University, through the Program of Research Project Funding After Publication, grant No (43- PRFA-P-43).

References

[1] R. Hilfer, Applications of Fractional Calculus in Physics, Academic press, Orlando, 1999.
 [2] A. Kilbas, H. Srivastava and J. Trujillo, Theory and Applications of Fractional Differential Equations, Elsevier, 2006.

- [3] I. Podlubny, *Fractional Differential Equations* (Academic Press, New York, 1999).
- [4] R. Amin, et al, Solution of Variable-Order Nonlinear Fractional Differential Equations Using Haar Wavelet Collocation Technique, *Fractals* 31 (2023) 2340022.
- [5] M. Manjula, et al, Existence, uniqueness and approximation of nonlocal fractional differential equation of sobolev type with impulses, *AIMS Mathematics* 8 (2023) 4645-4665.
- [6] Z. Alqahtani and A. E. Hagag, A fractional numerical study on a plant disease model with replanting and preventive treatment. *Métodos numéricos para cálculo y diseño en ingeniería: Revista internacional*, 39(3) (2023) 1-21.
- [7] S. Panda, V. Vijayakumar, and K. S. Nisar, Applying periodic and anti-periodic boundary conditions in existence results of fractional differential equations via nonlinear contractive mappings, *Boundary Value Problems* 2023 (2023) 91.
- [8] M. Sivashankar, et al, Some properties and stability of Helmholtz model involved with nonlinear fractional difference equations and its relevance with quadcopter, *Chaos, Solitons & Fractals* 168 (2023) 113161.
- [9] C. Cattani, Haar wavelet-based technique for sharp jumps classification, *Math. Comput. Model.* 39 (2004) 255.
- [10] Z. Alqahtani and A. E. Hagag, A new semi-analytical solution of compound KdV-Burgers equation of fractional order, *Métodos numéricos para cálculo y diseño en ingeniería: Revista internacional*, 39(4) (2023) 1-16.
- [11] A. Khan, et al, Exact Controllability of Hilfer Fractional Differential System with Non-instantaneous Impluseses and State Dependent Delay, *Qualitative Theory of Dynamical Systems*, 22 (2023) 62.
- [12] S. A. El-Wakil, A. Elhanbaly, M. A. Abdou. Adomian decomposition method for solving fractional nonlinear differential equations, *Applied Mathematics and Computation* 182 (1) (2006) 313-324.
- [13] V. Daftardar-Gejji, H. Jafari. Adomian decomposition: a tool for solving a system of fractional differential equations, *Journal of Mathematical Analysis and Applications* 301 (2) (2005) 508-518.
- [14] M. Rawashdeh, and S. Maitama, Solving coupled system of nonlinear PDE's using the natural decomposition method, *International Journal of Pure and Applied Mathematics* 92 (2014) 757-776.
- [15] M. Rawashdeh, and S. Maitama, Finding exact solutions of nonlinear PDEs using the natural decomposition method, *Mathematical Methods in the Applied Sciences* 40 (2017) 223-236.
- [16] R. Khandelwal, P. Kumawat, and Y. Khandelwal, A study of natural transform based on decomposition method for solving nonlinear ordinary differential equation, *Int. J. Stat. Appl. Math.* 3 (2018) 664-669.
- [17] M. S. Rawashdeh, The fractional natural decomposition method: Theories and applications, *Math. Methods Appl. Sci.* 40 (2017) 2362.
- [18] R. Shah, et al., Natural transform decomposition method for solving fractional-order partial differential equations with proportional delay, *Mathematics* 7 (2019) 532.
- [19] M. Rawashdeh, The fractional natural decomposition method: theories and applications, *Mathematical Methods in the Applied Sciences* 40 (2017) 2362-2376.
- [20] M. Elbadri, et al., A new solution of time-fractional coupled KdV equation by using natural decomposition method, *In Abstract and Applied Analysis*, 2020 (2020).
- [21] M. O. Korpusov, and A. G. Sveshnikov, Blow-up of solutions of strongly nonlinear equations of pseudoparabolic type, *Journal of Mathematical Sciences* 148 (2008) 1-142.
- [22] S. A. Dubey, Numerical solution for nonlocal Sobolev-type differential equations, *In Electron. J. Differ. Equ. Conf* 19 (2010) 75-83.
- [23] E. Kaikina, P. Naumkin, and I. Shishmarev, The Cauchy problem for an equation of Sobolev type with power nonlinearity, *Izvestiya: Mathematics* 69 (2005) 59.
- [24] G. Karch, Asymptotic behaviour of solutions to some pseudoparabolic equations, *Mathematical Methods in the Applied Sciences* 20 (1997) 271-289.
- [25] D. H. Peregrine, Calculations of the development of an undular bore, *Journal of Fluid Mechanics* 25 (1966) 321-330.
- [26] T. Benjamin, J. Bona, and J. Mahony, Model equations for long waves in nonlinear dispersive systems, *Philosophical Transactions of the Royal Society of London. Series A, Mathematical and Physical Sciences* 272 (1972) 47-78.
- [27] R. Camassa, and D. Holm, An integrable shallow water equation with peaked solitons, *Physical review letters* 71 (1993) 1661.
- [28] D. Rollins, Painlevé analysis and Lie group symmetries of the regularized long-wave equation, *Journal of mathematical physics* 32 (1991) 3331-3332.
- [29] Y. Kodama, On solitary-wave interaction, *Physics Letters A* 123 (1987) 276-282.
- [30] W. Hereman et al., Exact solitary wave solutions of nonlinear evolution and wave equations using a direct algebraic method, *Journal of Physics A: Mathematical and General* 19 (1986) 607.
- [31] F. Mainardi, On the initial value problem for the fractional diffusion equation, in: *Wave and Stability in Continuous Media*, S. Rionero, T. Ruggeeri (Eds.), World Scientific, Singapore, 246, 1994.
- [32] D. Loonker and P. K. Banerji, Solution of fractional ordinary differential equations by natural transform, *International Journal of Mathematical Engineering and Science* 12(2) (2013) 1-7.
- [33] P. Singh and D. Sharma, Convergence and error analysis of series solution of nonlinear partial differential equation, *Nonlinear Engineering* 7 (2018) 303-308.
- [34] E. Kreyszig, *Further Applications: Banach Fixed Point Theorems*, Erwin Kreyszig, *Introductory Functional Analysis with Applications*. New York, NY: Wiley Classic Libraries; (1989) 299-321.
- [35] Ö. F. Gözükızıl, and Ş. Akçağıl, The tanh-coth method for some nonlinear pseudoparabolic equations with exact solutions, *Advances in Difference Equations* 2013(1) (2013) 1-18.35